

Anisotropies in the turbulent solar wind, perspective on multipoint observations from numerical simulations

Part 1

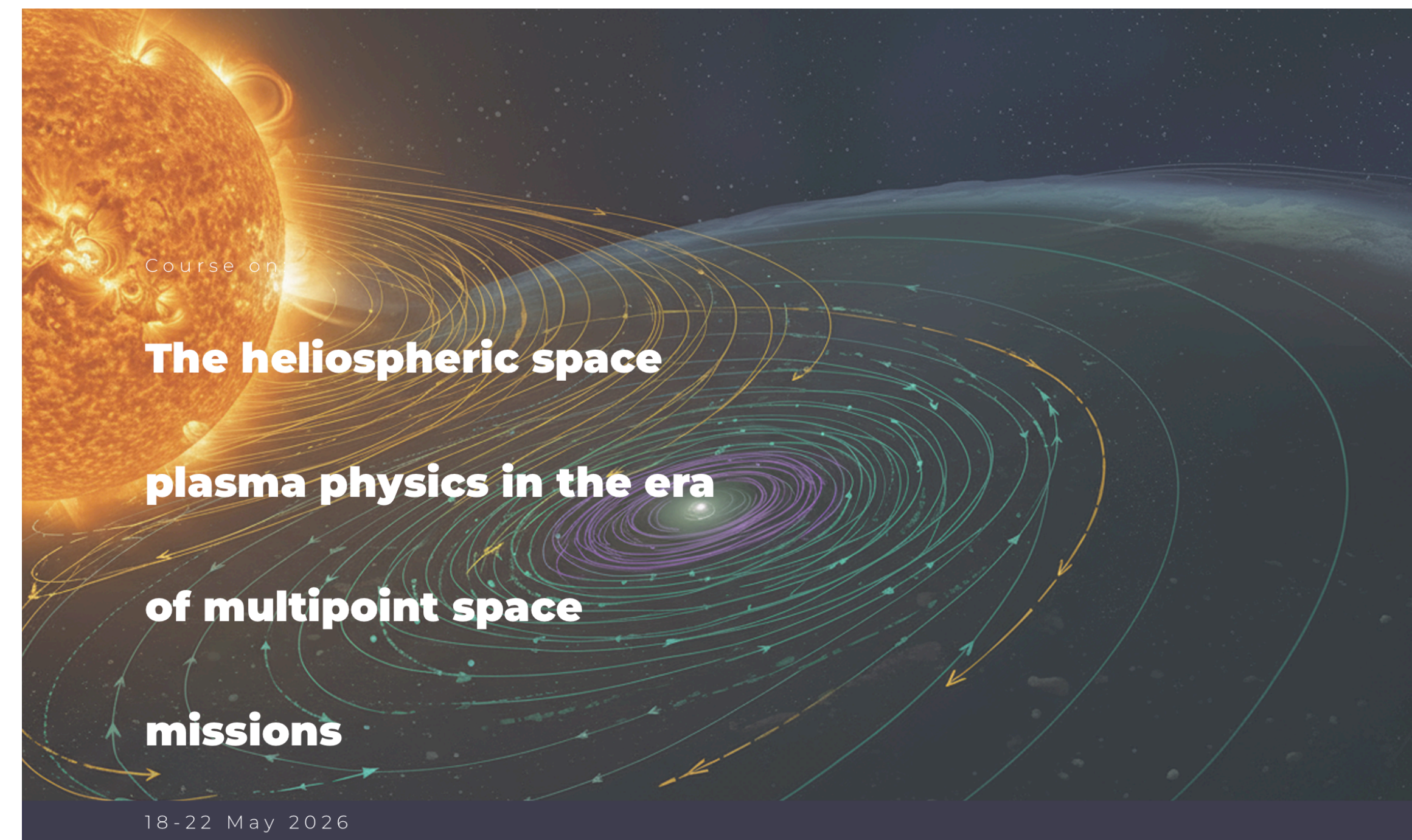
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Arcetri Space & Astrophysical Plasma (ASAP) group

<https://sites.google.com/unifi.it/asap>



Outline

- What anisotropy? (spectral and scale dependent anisotropy)
 - Why anisotropy ? (a semi-quantitative argument)
 - How anisotropy is driven ? (anisotropy of the cascade)
 - Measurements in the solar wind, some surprises:
 - Special solar wind anisotropy (Maltese Cross) at different scales
 - Scale dependent anisotropy
-
- Solar wind radial expansion: side effects
 - some explanations for the observed solar wind anisotropies
 - Why do we need multi-spacecraft ?
 - an example by revising spectral anisotropy
 - anisotropy of the cascade (the driver) revealed?

What kind of anisotropy? Spectral Anisotropy

Uneven distribution of energy in the **wavevector** space

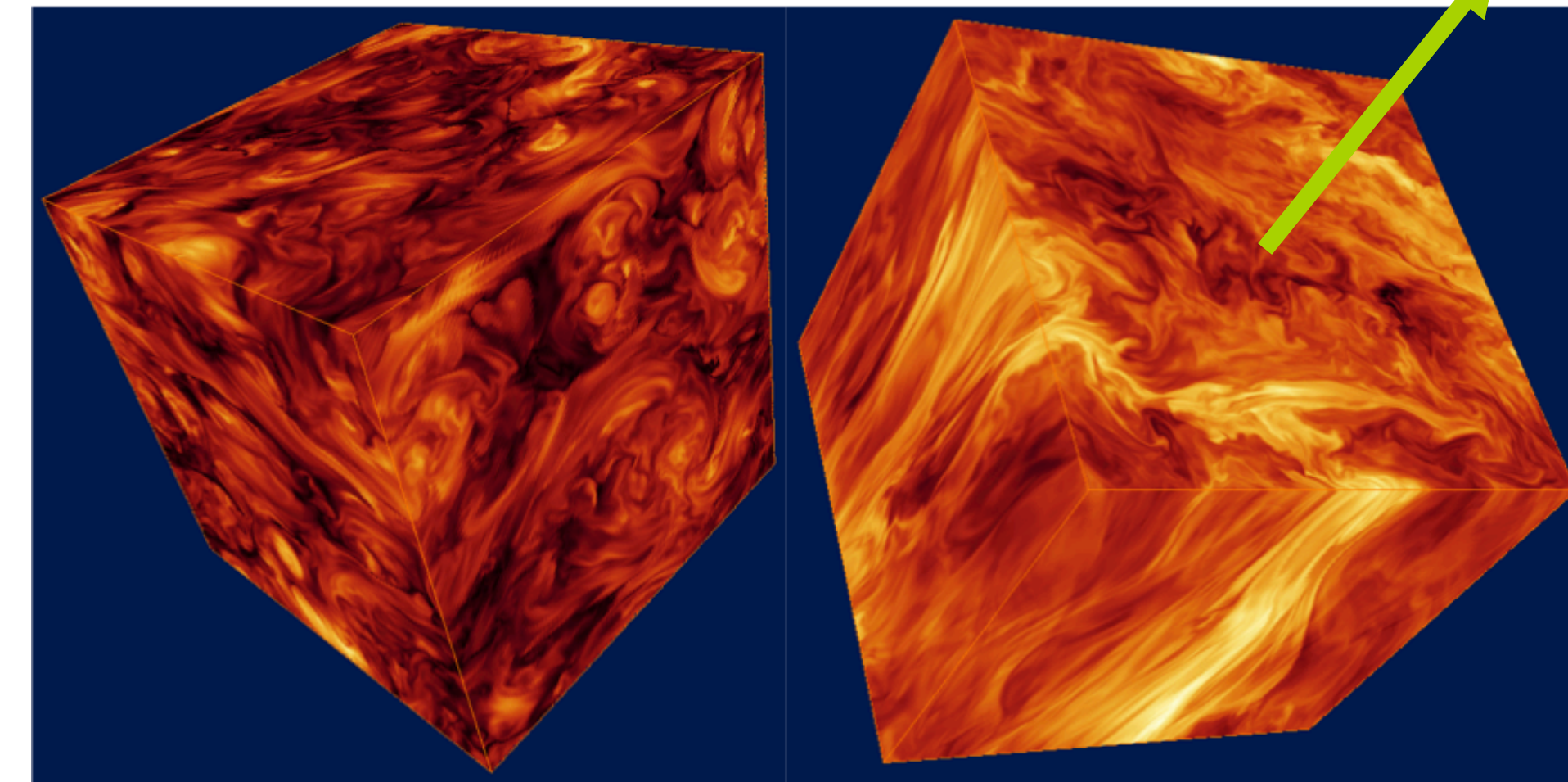
or

Uneven distribution of energy in the **increment** space

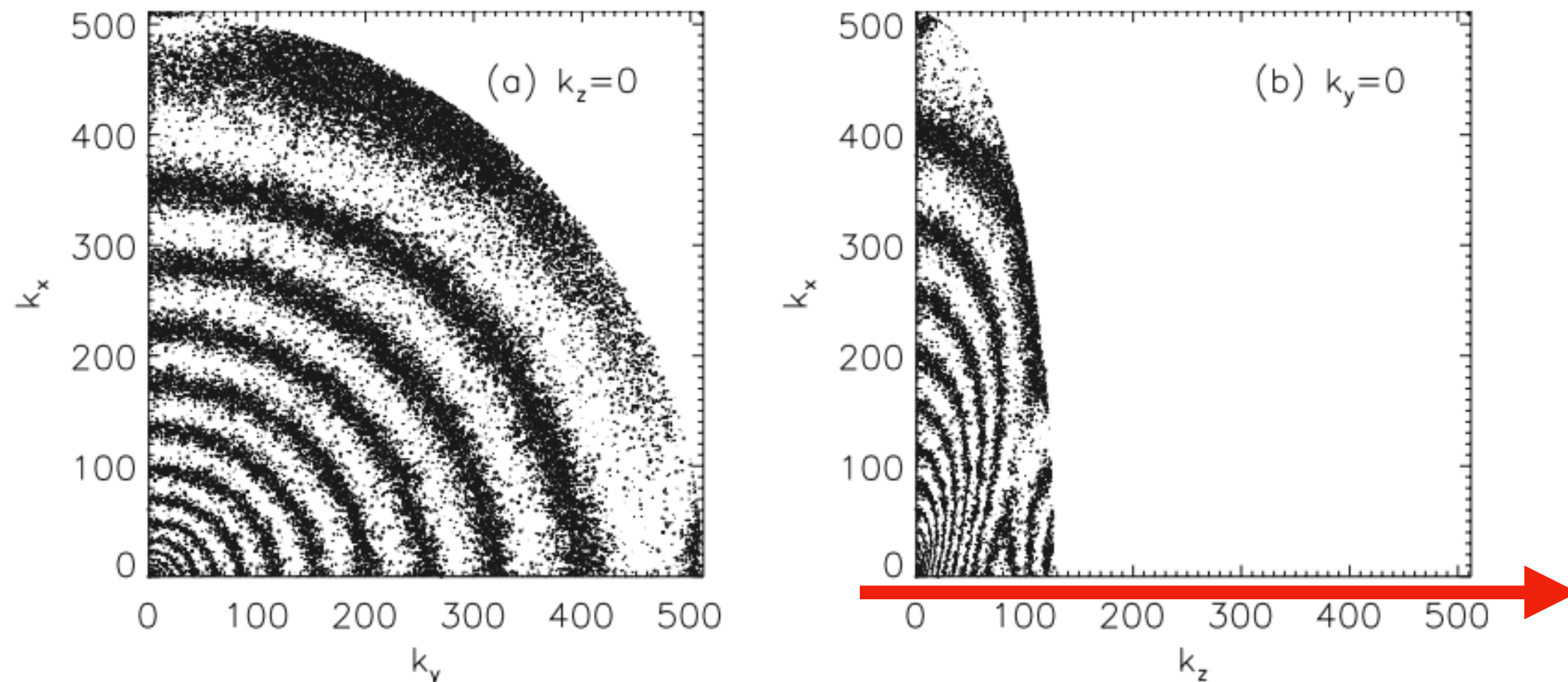
Muller & Biskamp 2005, PRE

NO mean field

WITH mean field



spectra



$$C(\ell) = \text{FT}^{-1}[E(\mathbf{k})]$$

$$C(\ell) = \langle u_i(\mathbf{x} + \ell) u_i(\mathbf{x}) \rangle$$

$$SF(\ell) = \langle |u_i(\mathbf{x} + \ell) - u_i(\mathbf{x})|^2 \rangle$$

$$SF = 2[C(0) - C(\ell)]$$

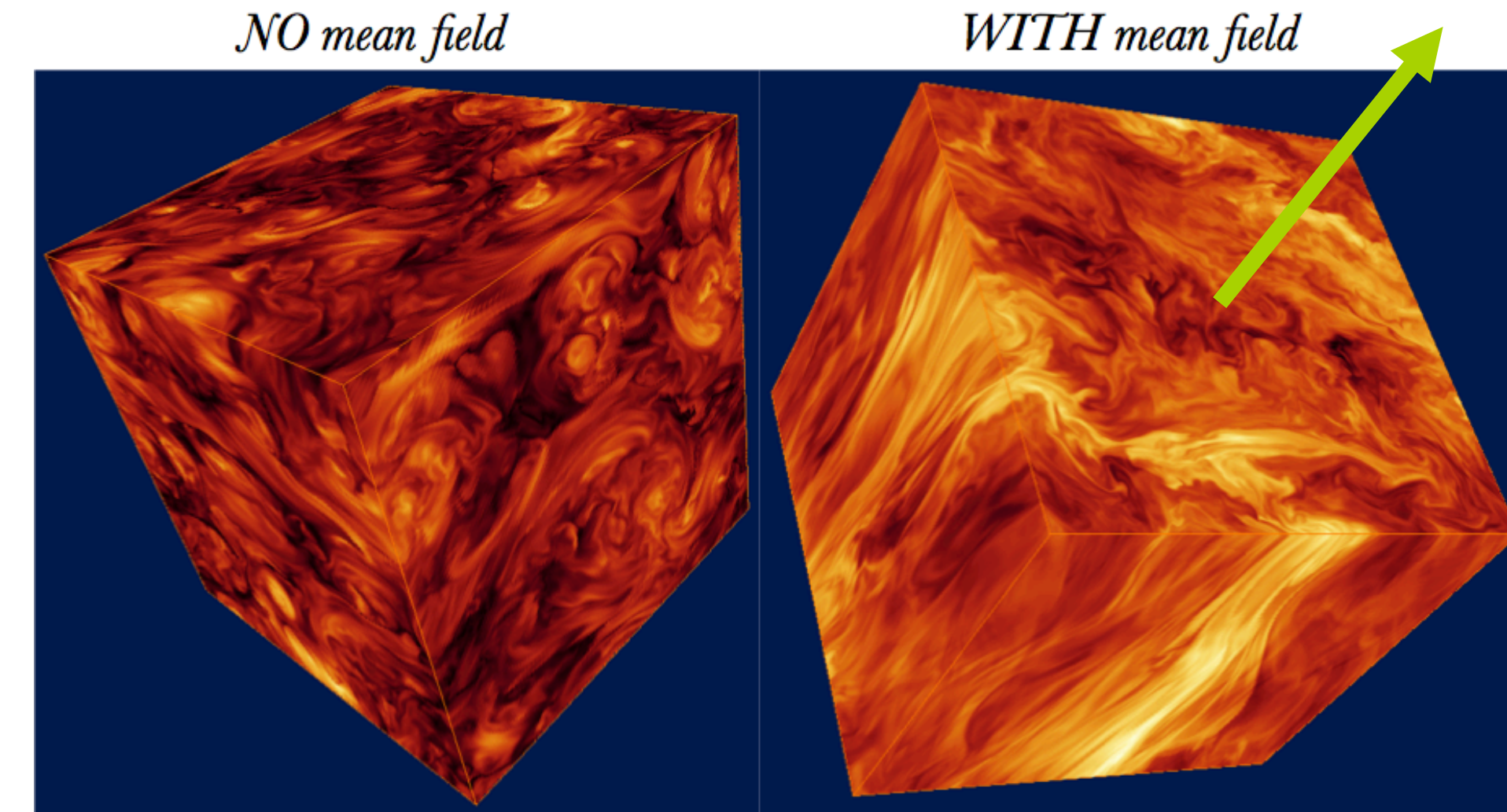
What kind of anisotropy? Spectral Anisotropy

Muller & Biskamp 2005, PRE

Uneven distribution of energy in the **wavevector** space

or

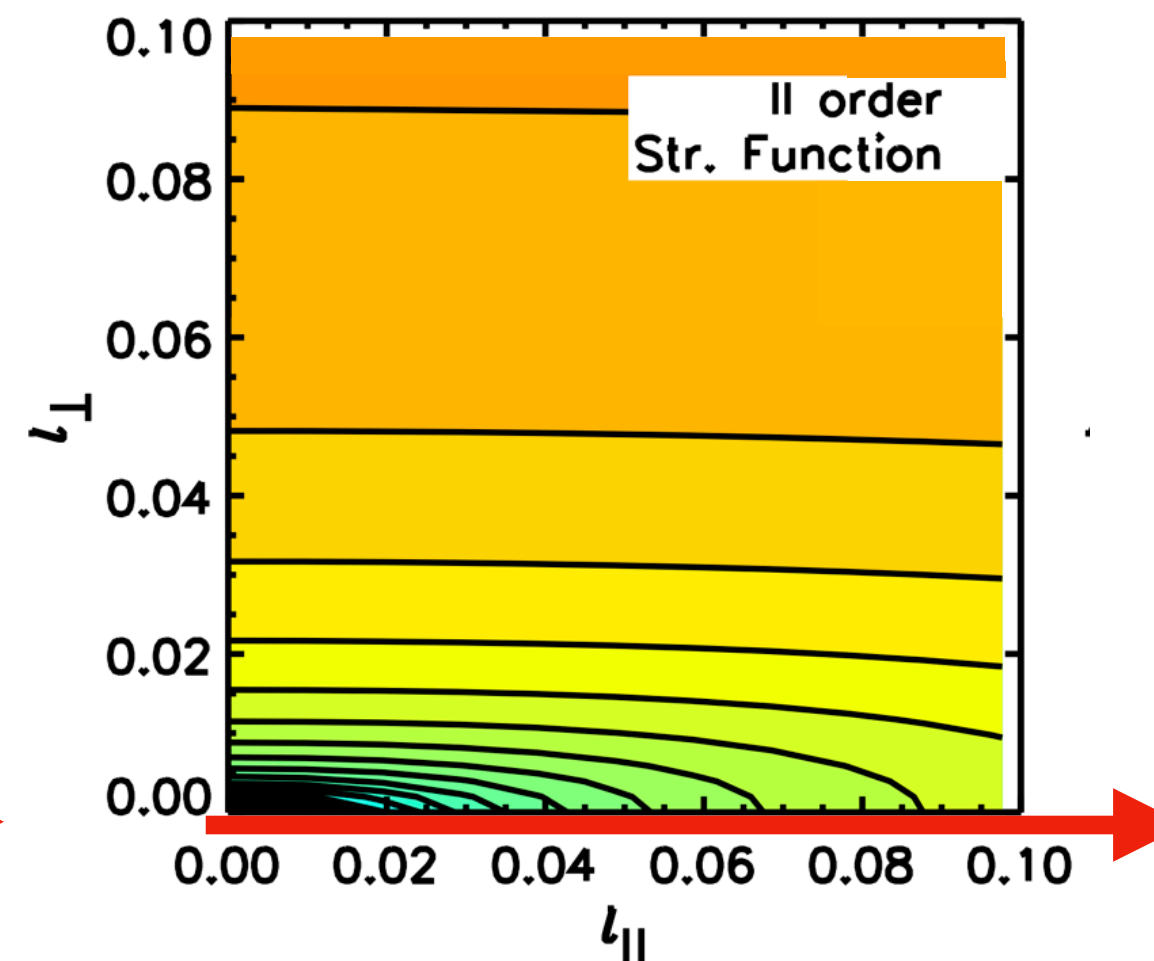
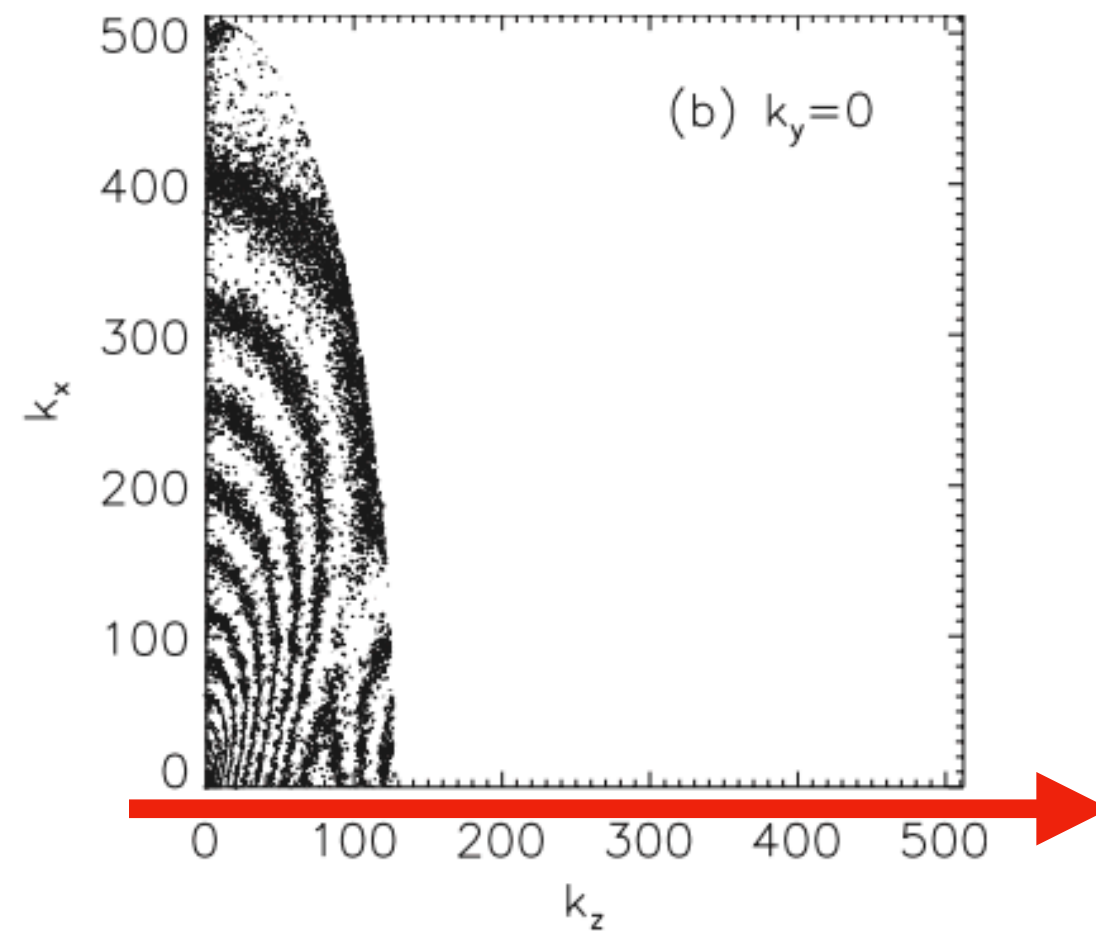
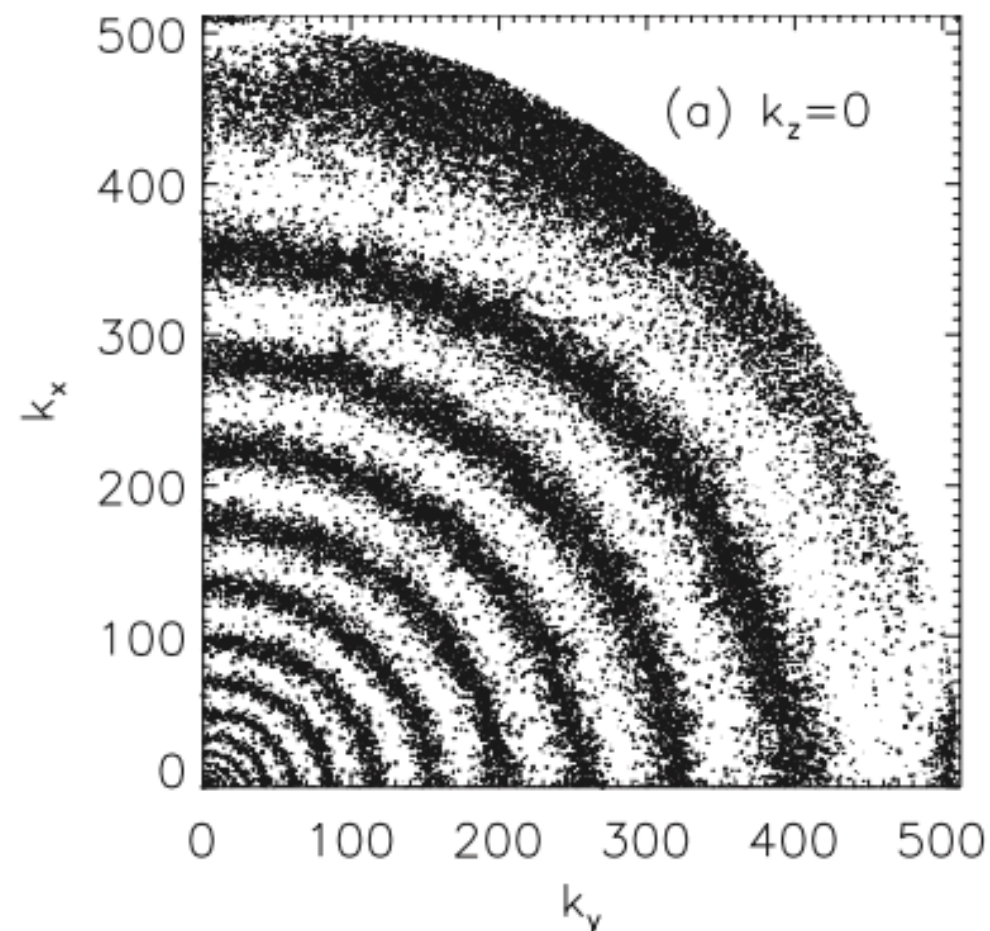
Uneven distribution of energy in the **increment** space



spectra

flipped

SF



$$C(\boldsymbol{\ell}) = \text{FT}^{-1}[E(\mathbf{k})]$$

$$C(\boldsymbol{\ell}) = \langle u_i(\mathbf{x} + \boldsymbol{\ell}) u_i(\mathbf{x}) \rangle$$

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$$SF = 2[C(0) - C(\boldsymbol{\ell})]$$

Why Spectral Anisotropy?

$$\partial_t \mathbf{z}^\pm \mp \mathbf{B}_0 \cdot \nabla \mathbf{z}^\pm + \mathbf{z}^\mp \cdot \nabla \mathbf{z}^\pm + \nabla P = 0$$

$$\mathbf{z}^\pm = \delta \mathbf{u} \mp \delta \mathbf{b} / \sqrt{4\pi\rho}$$

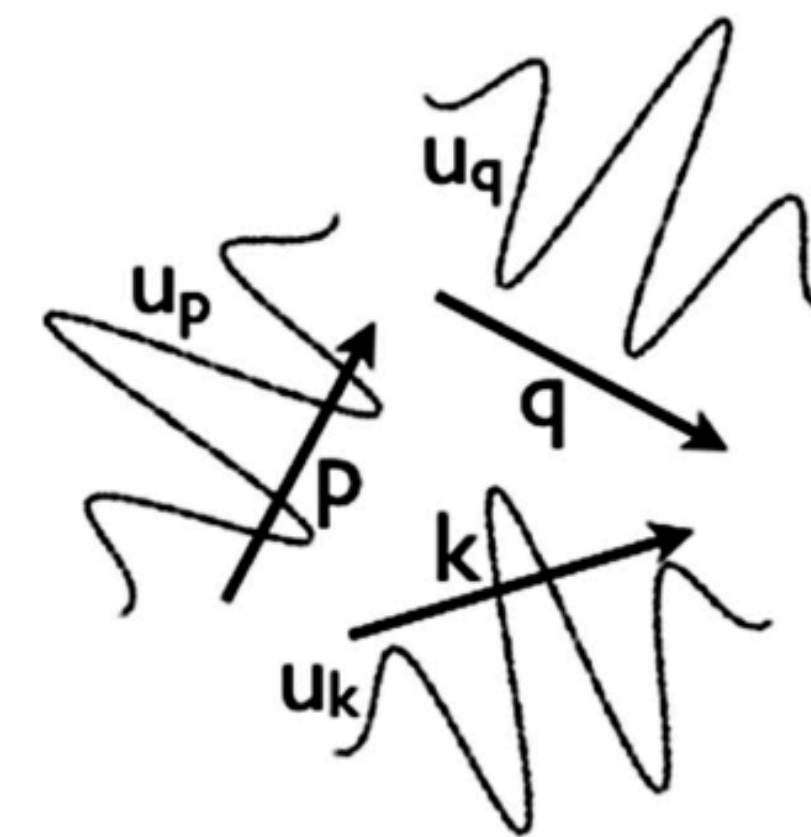
$$\partial_t \widehat{z}_i^\pm(k) \mp ik_\parallel B_0 \widehat{z}_i^\pm(k) + iP_{ijkl}(k) \widehat{z}_l^\pm \widehat{z}_j^\mp(k) = 0$$

$$P_{ijkl}(k) = k_j (\delta_{il} - k_l k_i / k^2)$$

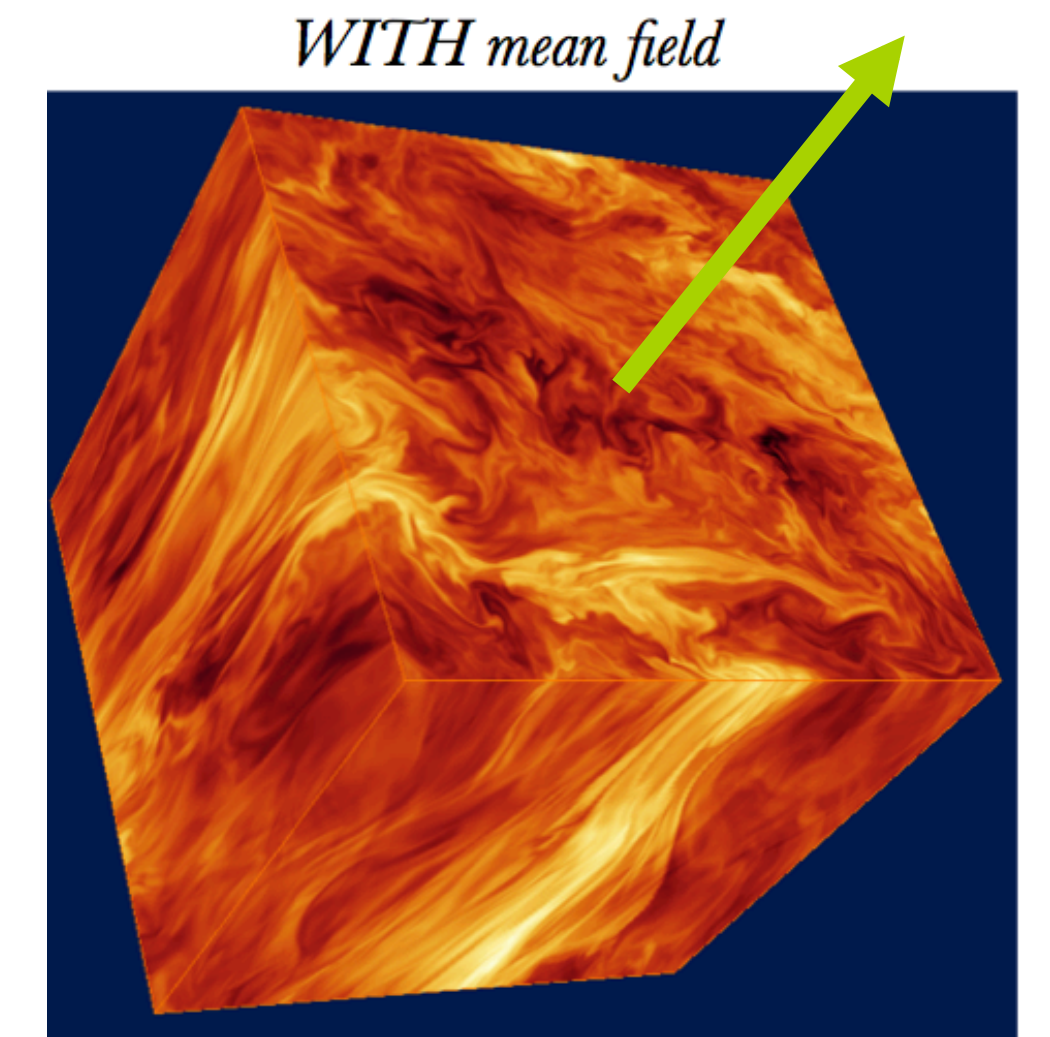
$$\widehat{\mathbf{z}}^\pm(k) = \mathbf{Z}^\pm(k) e^{\pm ik_\parallel B_0 t}$$

$$\partial_t \mathbf{Z}_i^\pm(k) + iP_{ijkl}(k) \int_{k=p+q} d^3 p d^3 q \mathbf{Z}_l^\pm(p) \mathbf{Z}_j^\mp(q) e^{\mp i\omega t} = 0$$

$$\omega = 2q_\parallel B_0 \sim 1/t_A$$



hydro



Muller & Biskamp 2005, PRE

Why Spectral Anisotropy?

$$\partial_t \mathbf{z}^\pm \mp \mathbf{B}_0 \cdot \nabla \mathbf{z}^\pm + \mathbf{z}^\mp \cdot \nabla \mathbf{z}^\pm + \nabla P = 0$$

$$\mathbf{z}^\pm = \delta \mathbf{u} \mp \delta \mathbf{b} / \sqrt{4\pi\rho}$$

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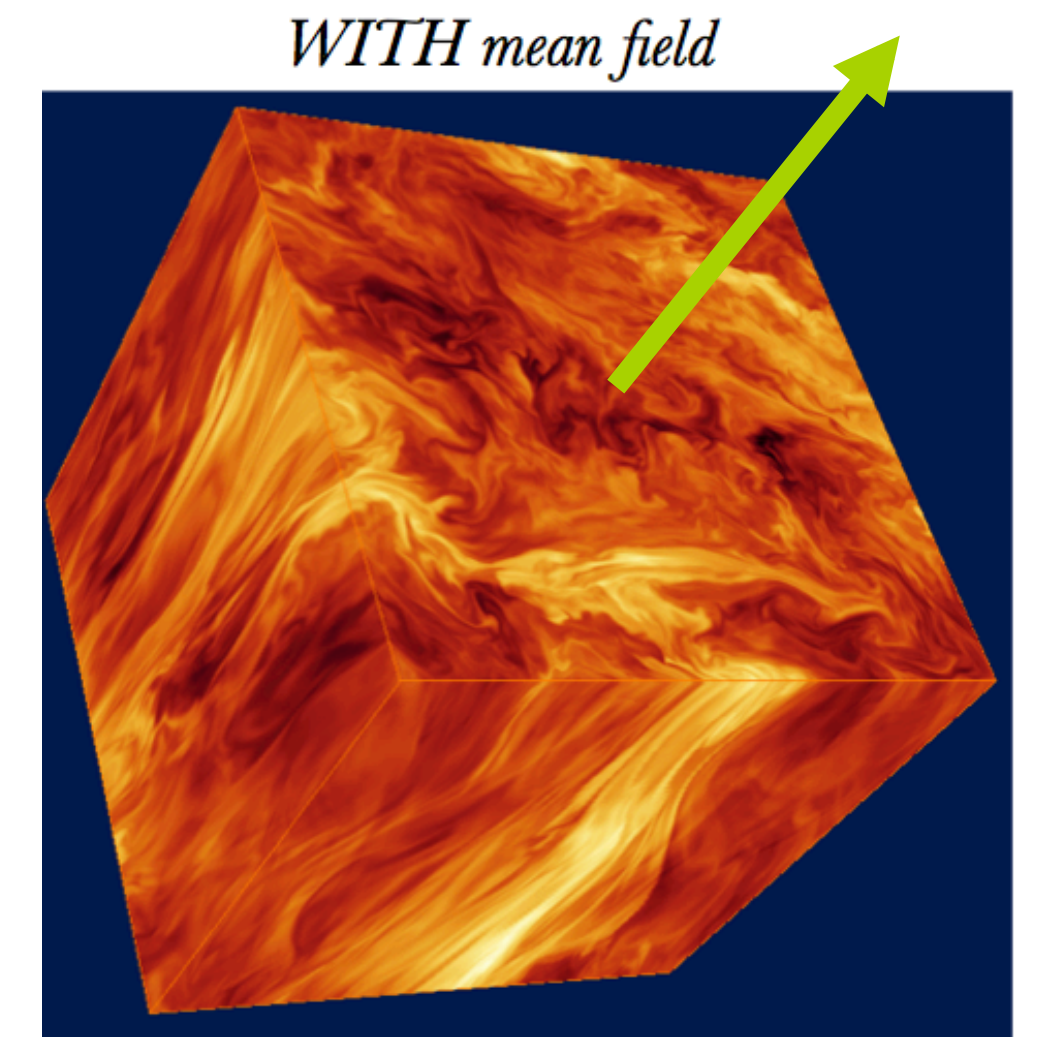
$$\widehat{z}^\pm(k) = \mathbf{Z}^\pm(k) e^{\pm ik_\parallel B_0 t}$$

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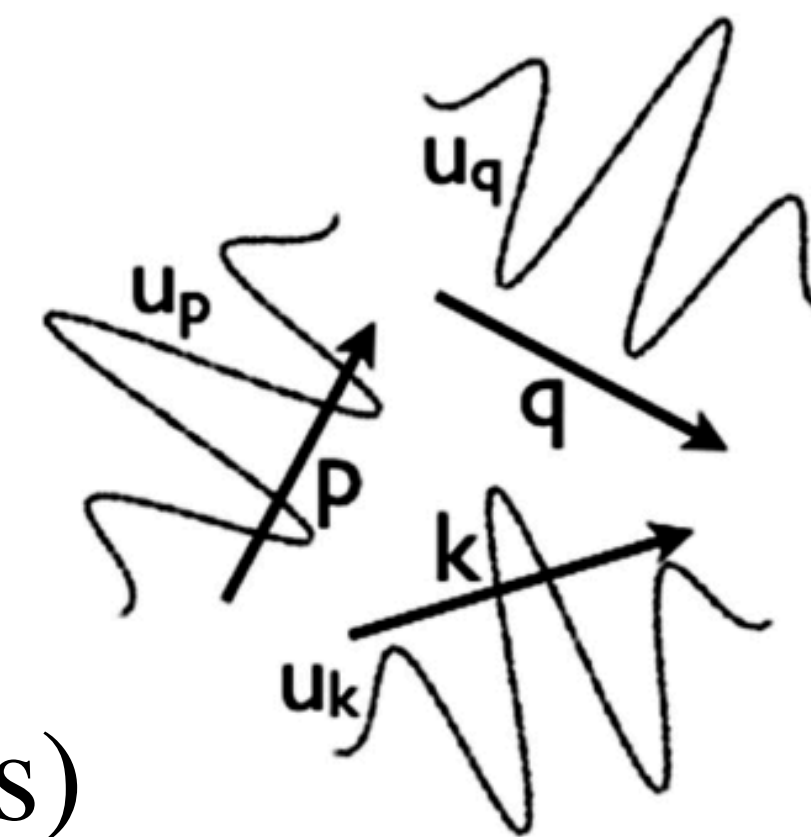
$$\omega = 2q_\parallel B_0 \sim 1/t_A$$

If $q \parallel$ to $B_0 \Rightarrow$ no contribution to NL cascade

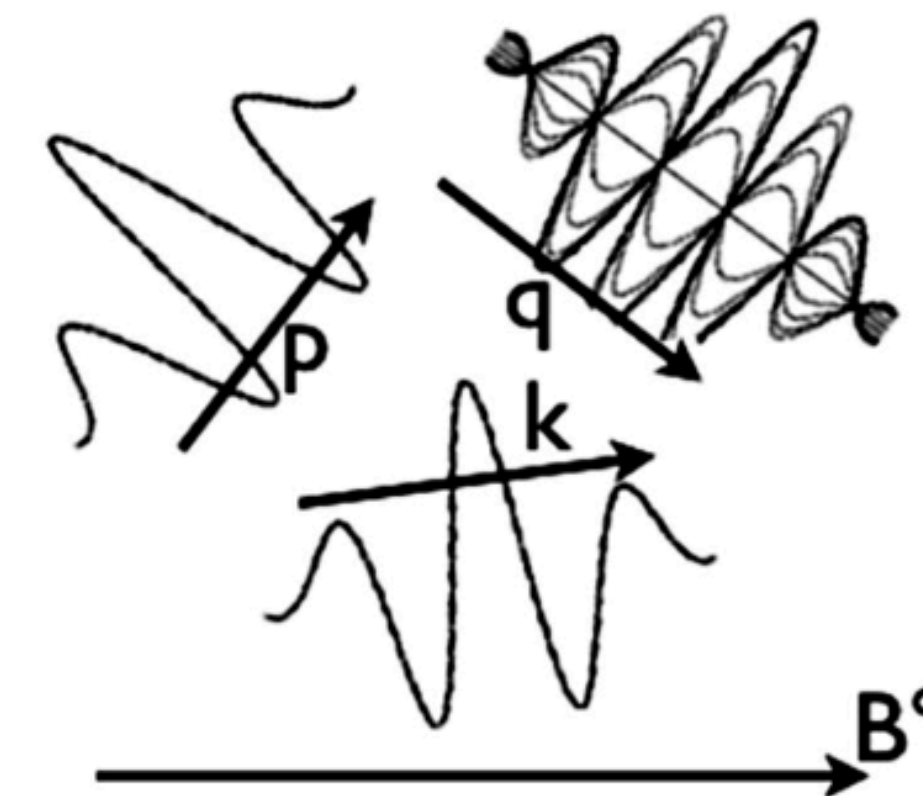
If $q \perp$ to $B_0 \Rightarrow$ no oscillatory part (Hydro-like interactions)



Muller & Biskamp 2005, PRE

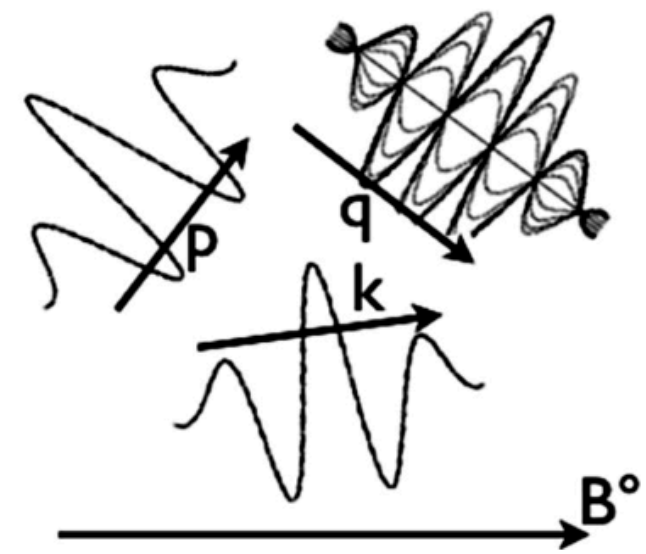
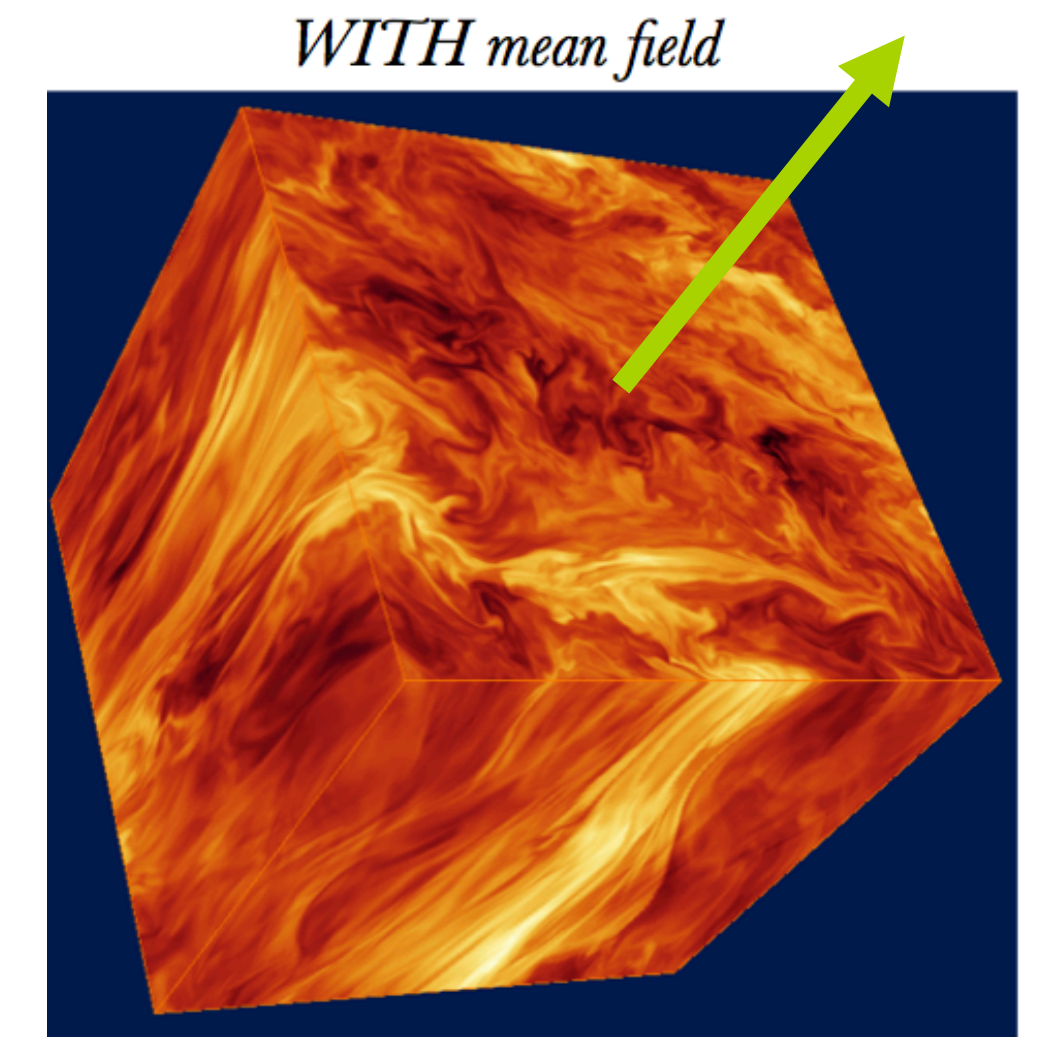


hydro



MHD with B^0

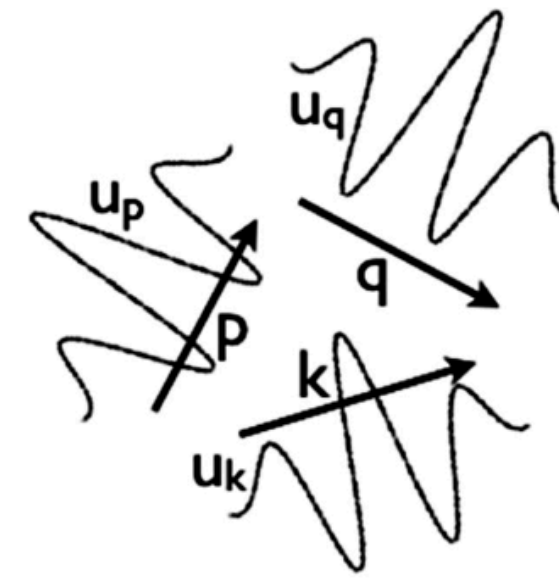
Why Spectral Anisotropy?



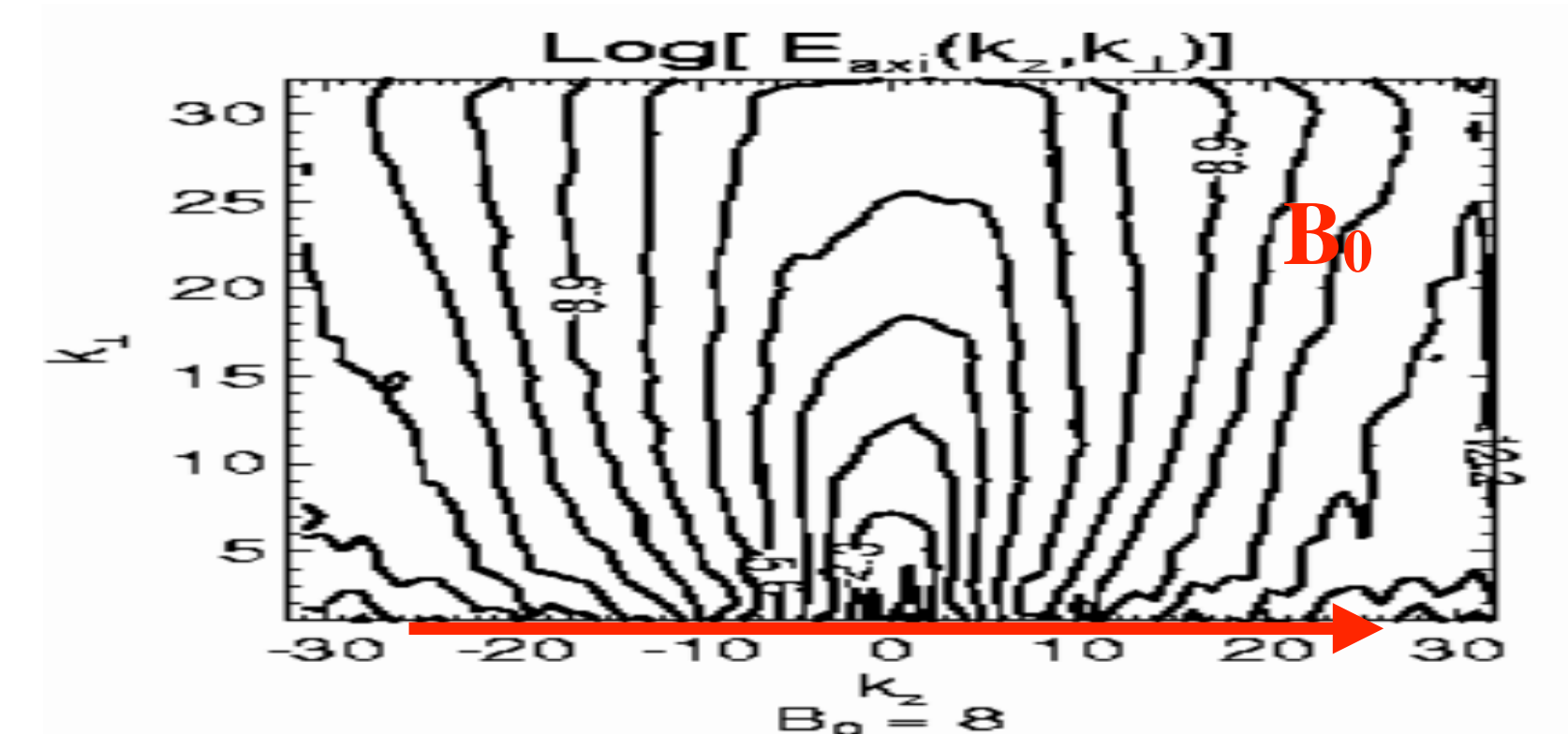
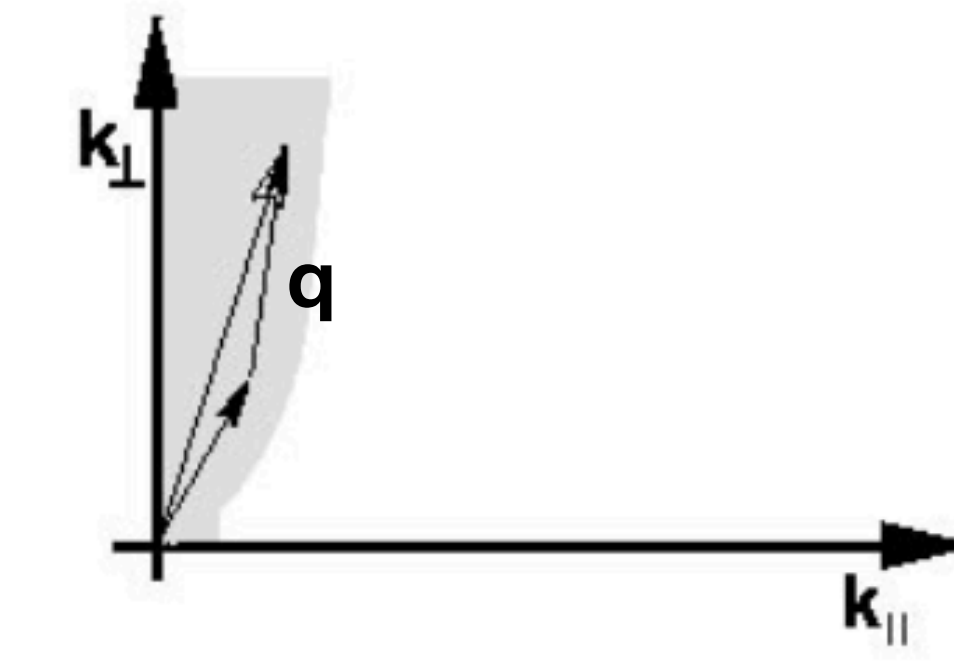
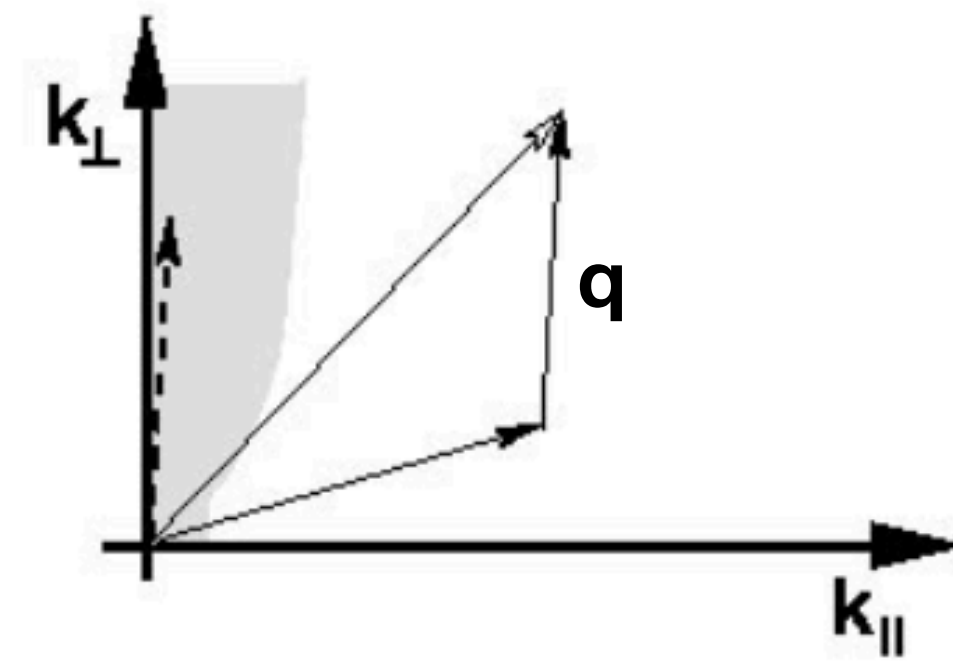
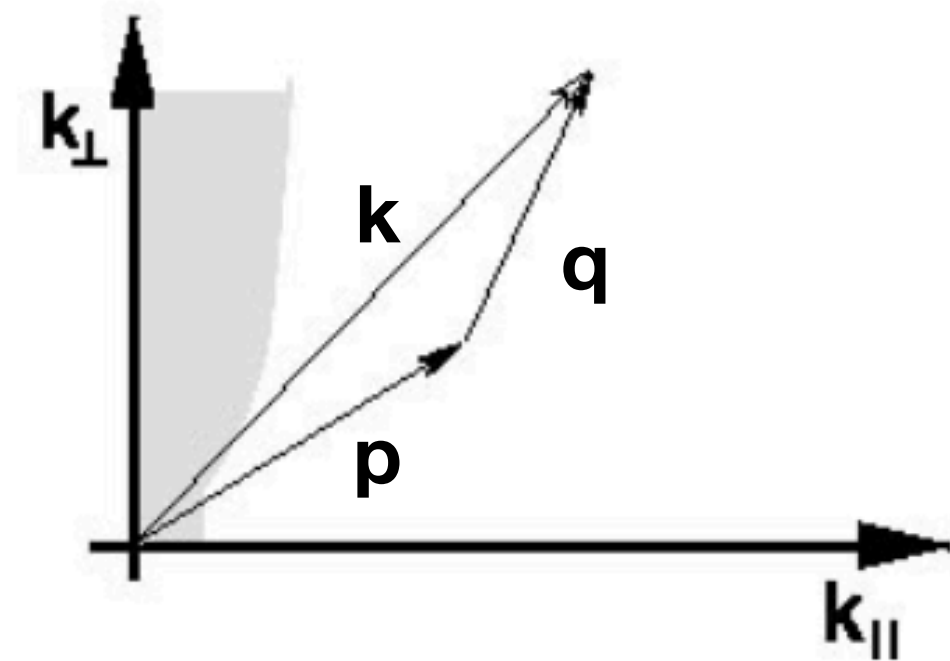
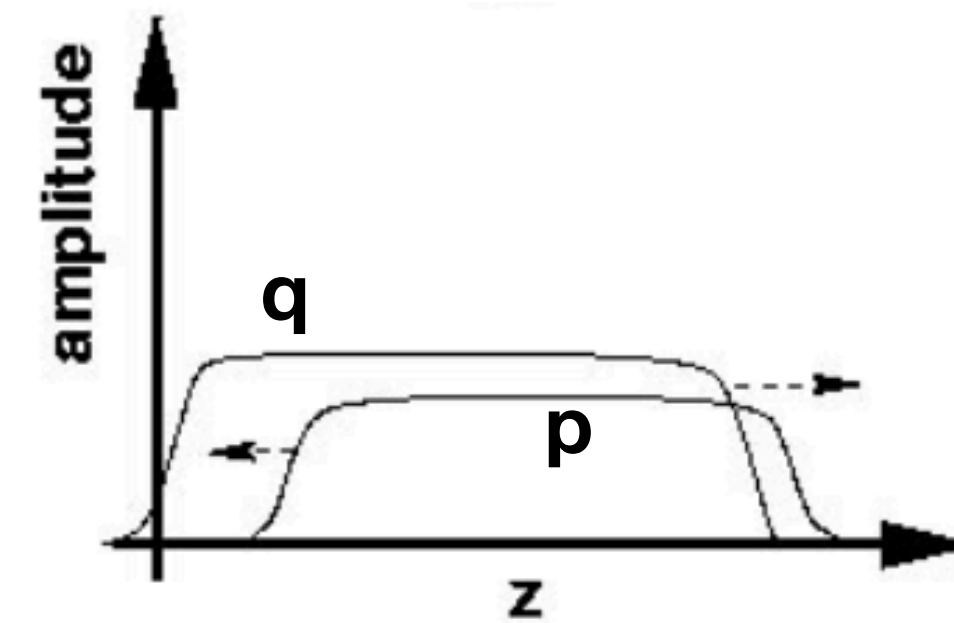
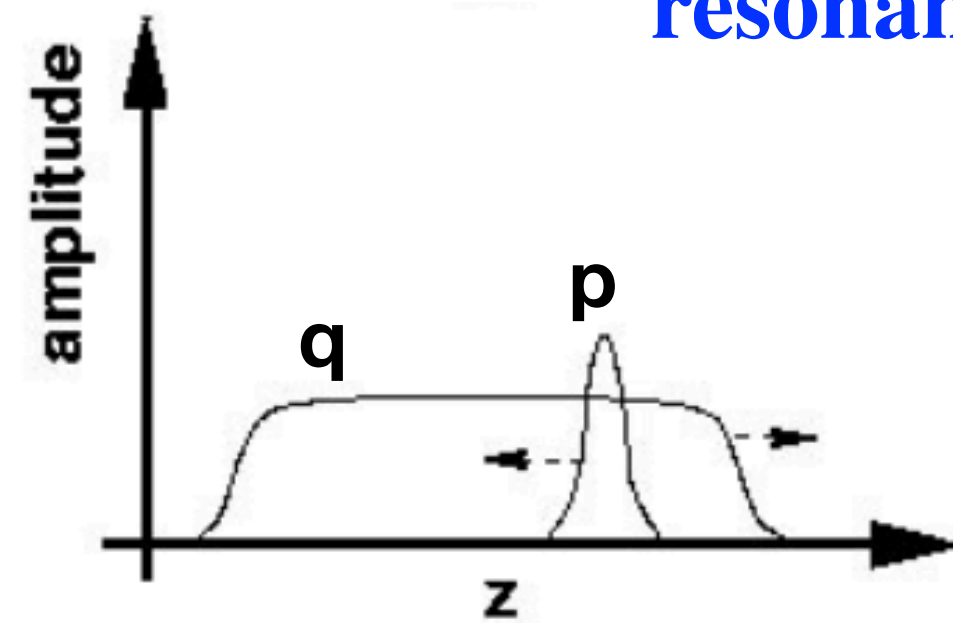
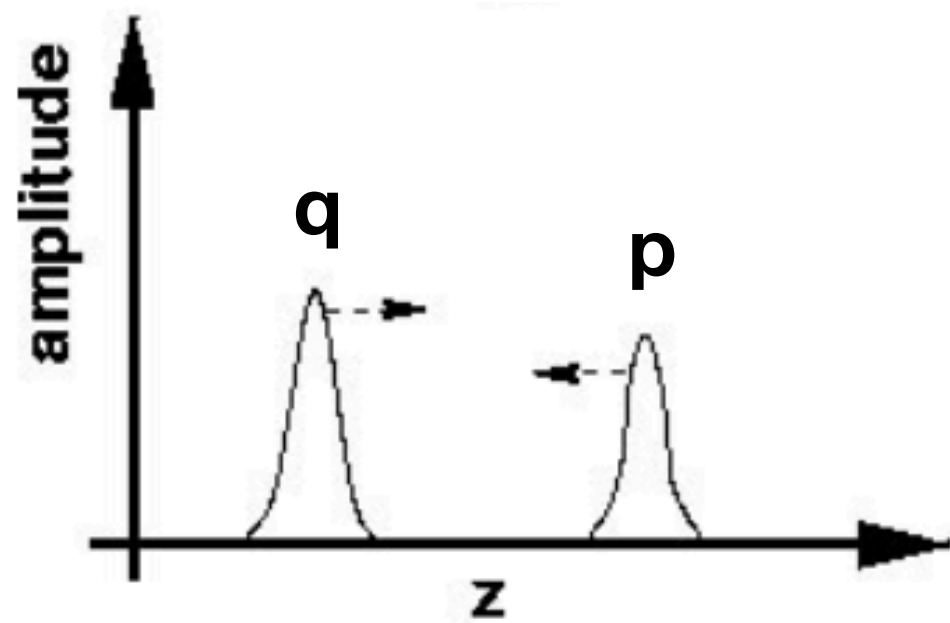
Wave+Wave: negligible

Hydro+Wave : quasi resonant

Hydro+Hydro: resonant



spectrum develops preferentially along k_{\perp}

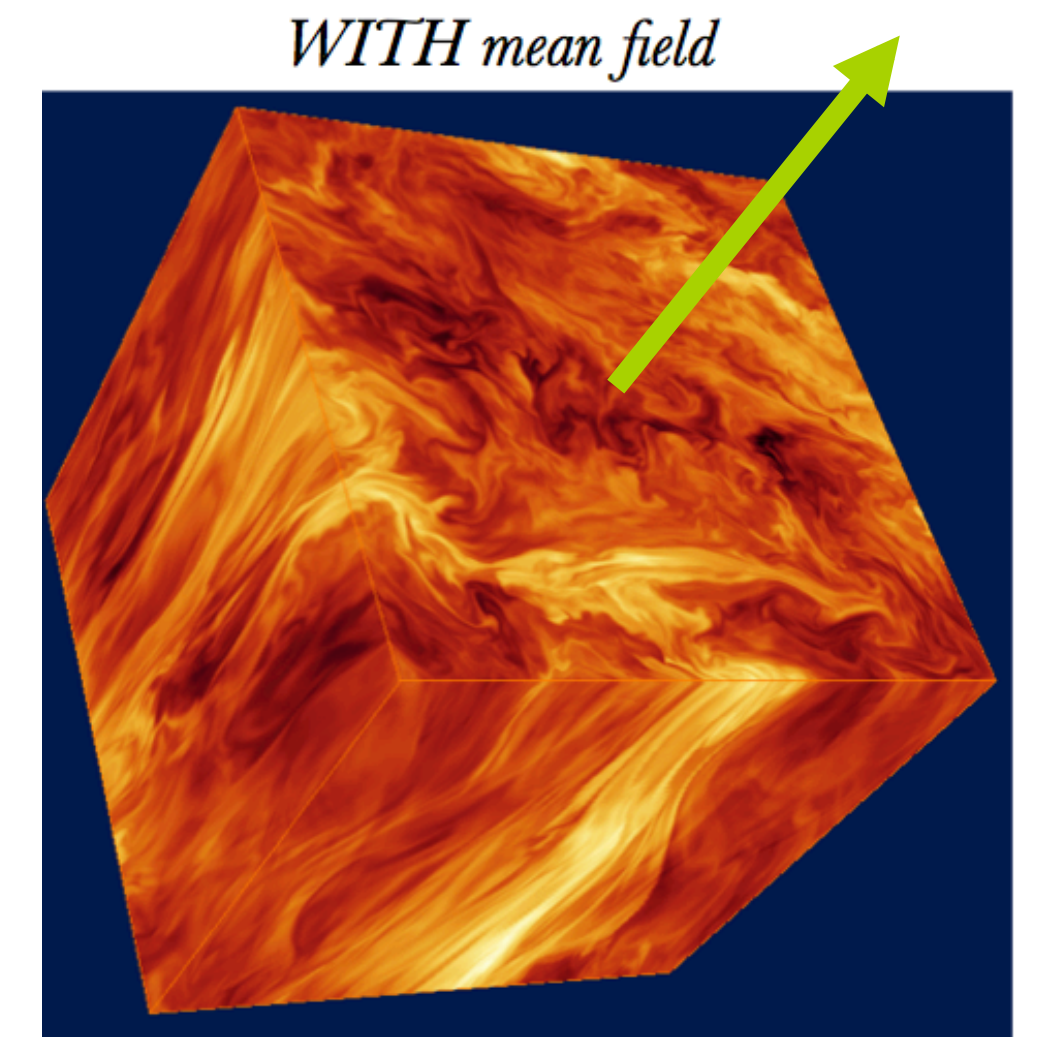


Spectral Anisotropy

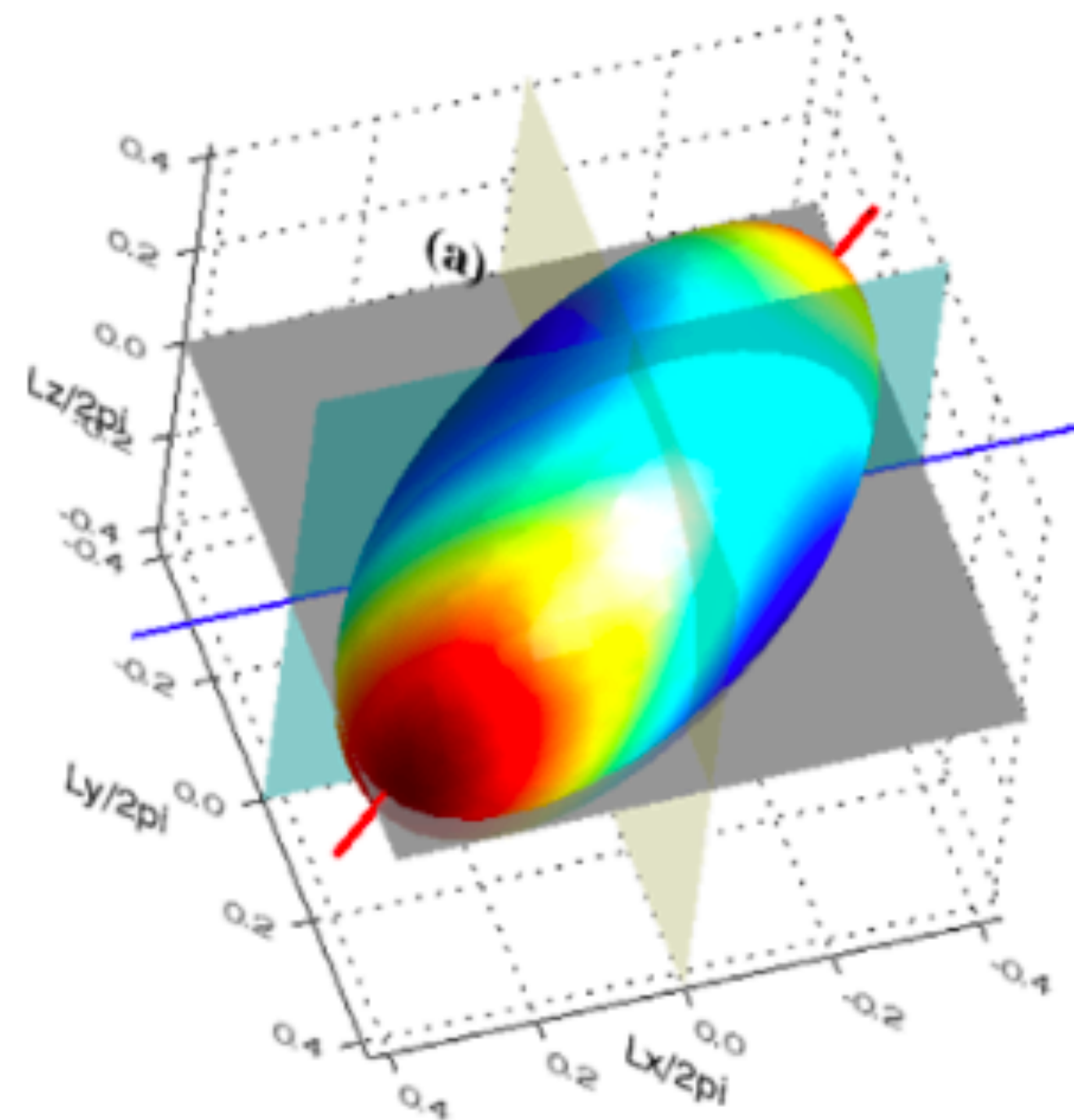
Spectral Anisotropy

$$C(\ell) = \text{FT}^{-1}[E(\mathbf{k})]$$

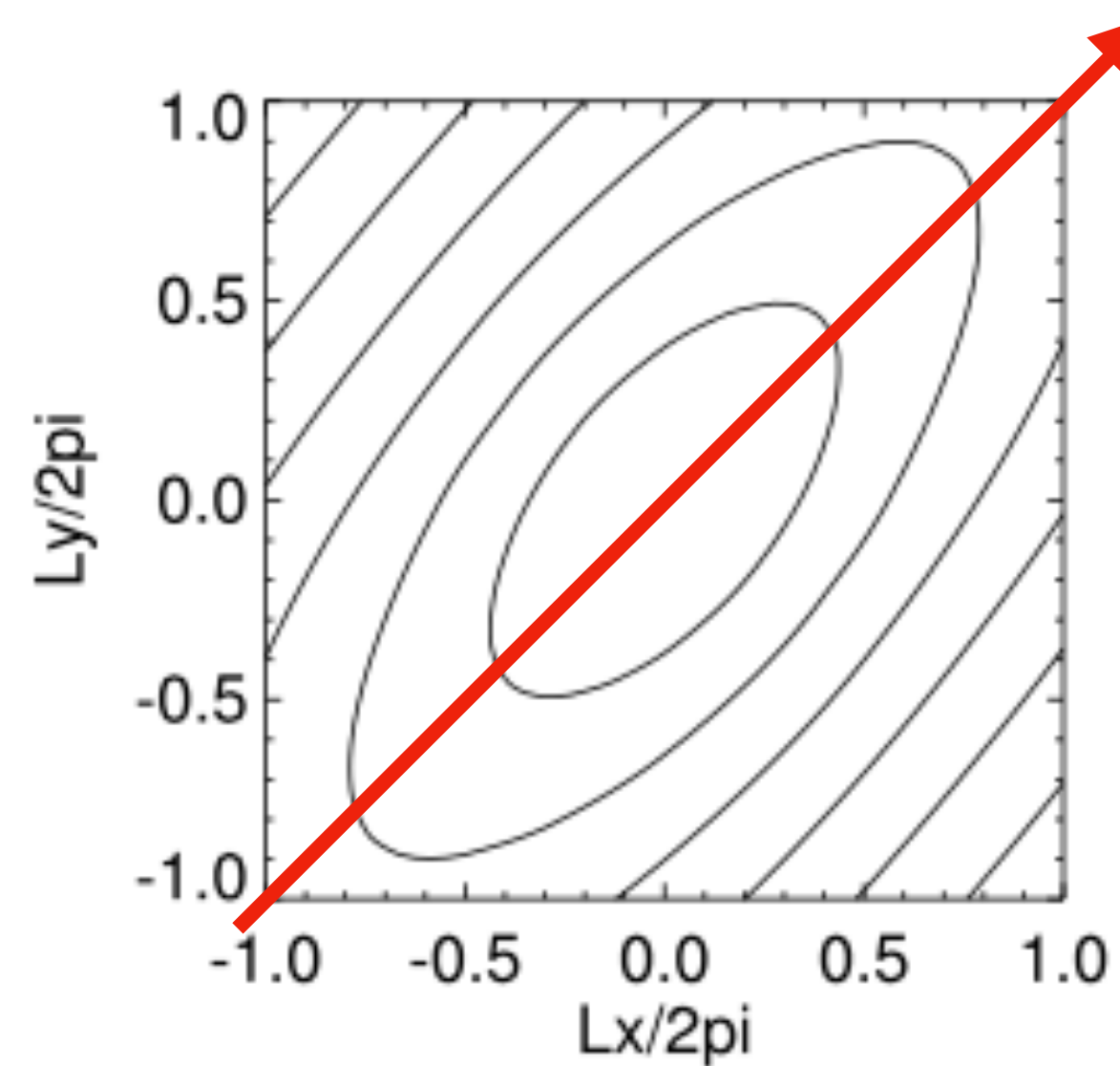
Driver?



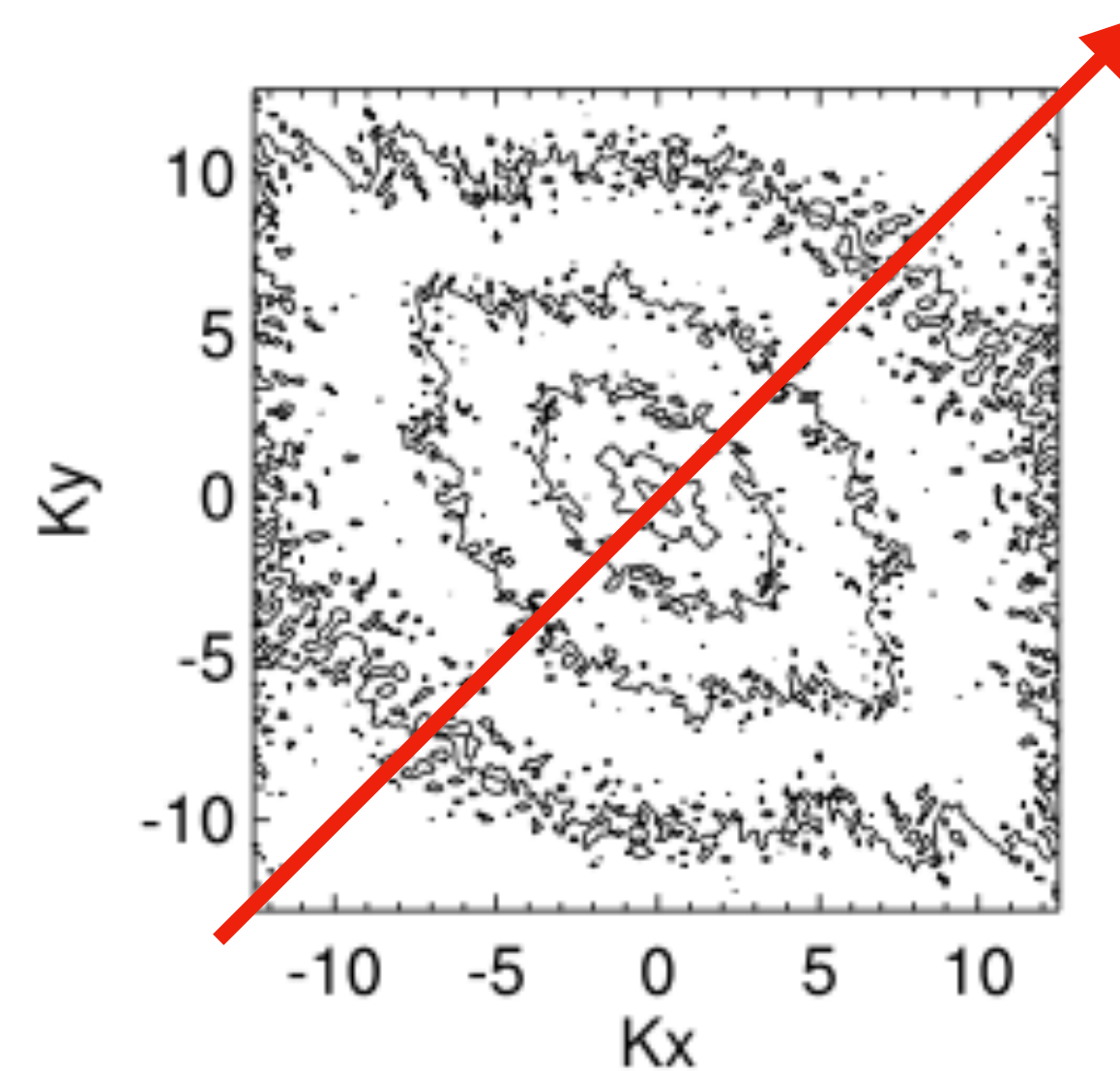
$C(\ell)$ 3D



$C(\ell_x, \ell_y)$



$E(k_x, k_y)$

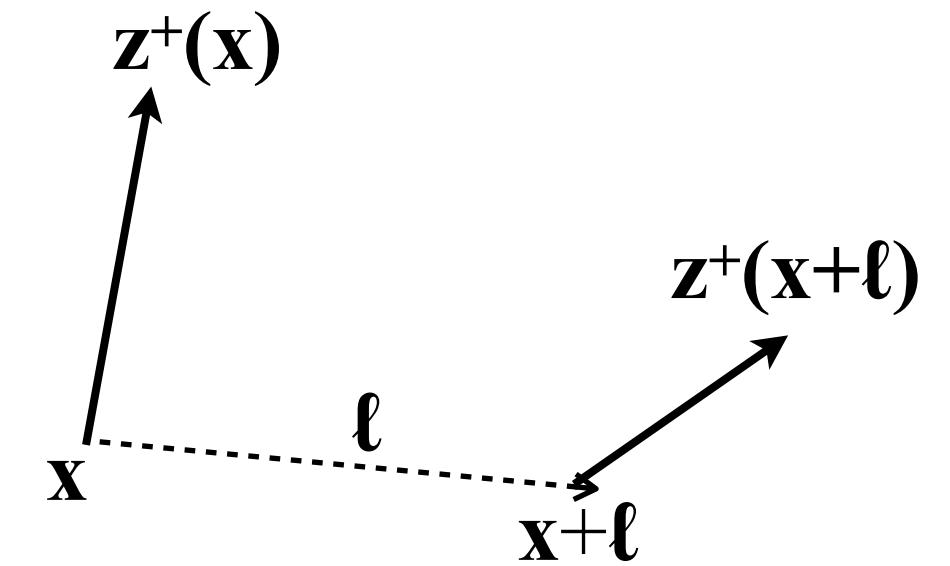


Isotropic cascade

$$\partial_t \mathbf{z}^\pm \mp \mathbf{B}_0 \cdot \nabla \mathbf{z}^\pm + \mathbf{z}^\mp \cdot \nabla \mathbf{z}^\pm + \nabla P = 0$$

$$\mathbf{z}^\pm = \delta \mathbf{u} \mp \delta \mathbf{b} / \sqrt{4\pi\rho}$$

- evaluate the eq. at \mathbf{x} and $\mathbf{x}+\boldsymbol{\ell}$,
- multiply by $\Delta \mathbf{z}^\pm$,
- average in the volume and obtain equations for S^\pm ,
- sum them so you get total energy $S=S^++S^-$



$$\Delta \mathbf{z}^\pm(\mathbf{x}, \boldsymbol{\ell}) = \mathbf{z}^\pm(\mathbf{x} + \boldsymbol{\ell}) - \mathbf{z}^\pm(\mathbf{x})$$

$$\partial_t S + \nabla_\ell \cdot \mathbf{Y} = -4\epsilon + 2\nu \nabla_\ell^2 S$$

Conservation equation for II-order Structure function S (mean field is not here!)

$$S = 1/2 [\langle |\Delta \mathbf{z}^-|^2 \rangle + \langle |\Delta \mathbf{z}^+|^2 \rangle]$$

$$\epsilon = \nu/2 [\langle \sum_i (\partial_{\mathbf{x}} z_i^+)^2 \rangle + \langle \sum_i (\partial_{\mathbf{x}} z_i^-)^2 \rangle]$$

$$\mathbf{Y} = 1/2 [\langle \Delta \mathbf{z}^- |\Delta \mathbf{z}^+|^2 \rangle + \langle \Delta \mathbf{z}^+ |\Delta \mathbf{z}^-|^2 \rangle]$$

II-order Structure Function

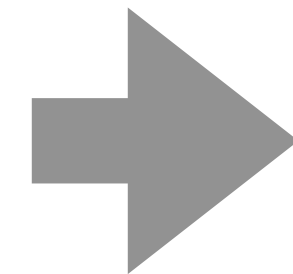
Dissipation rate

III-order Structure Function

Isotropic cascade

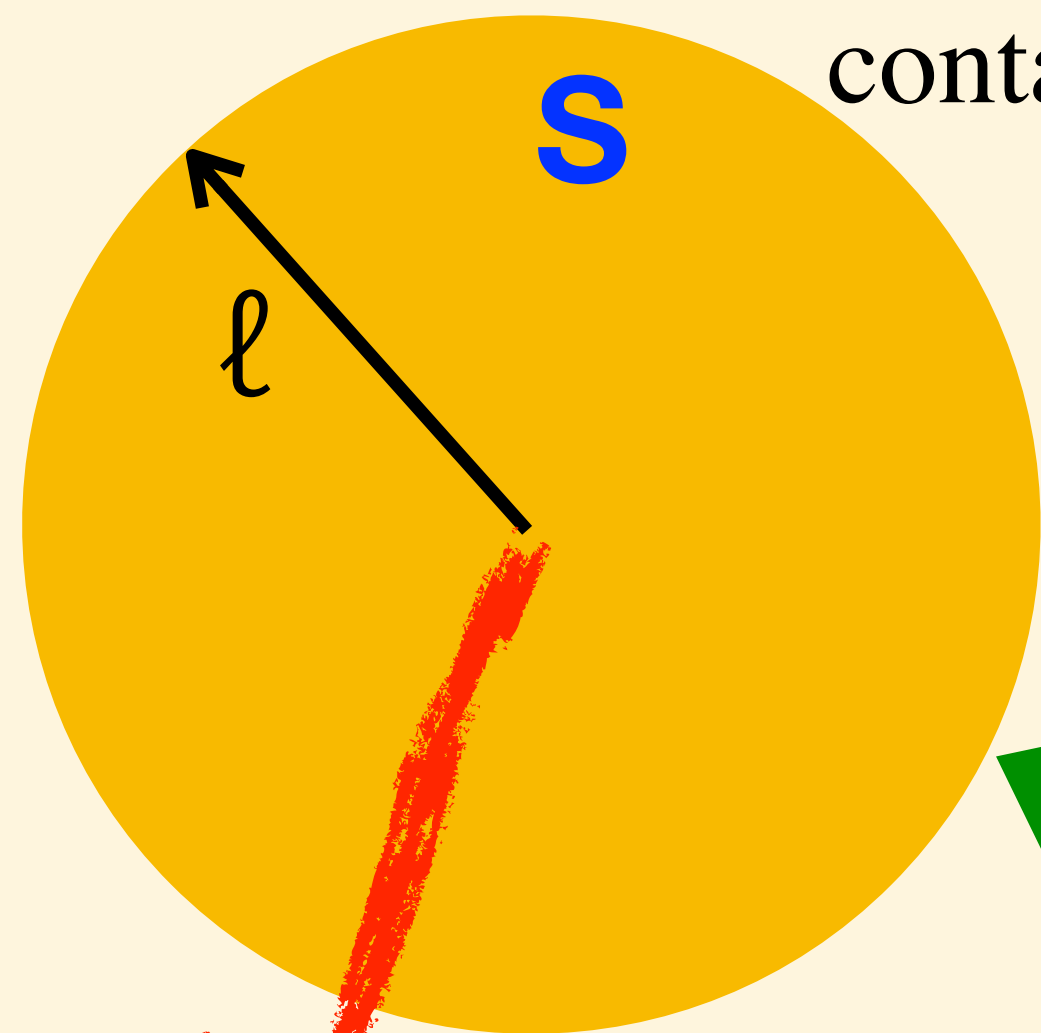
Isotropic => can study its angular average (or a cut)

$$\partial_t S + \nabla_\ell \cdot \mathbf{Y} = -4\epsilon + 2\nu \nabla_\ell^2 S$$



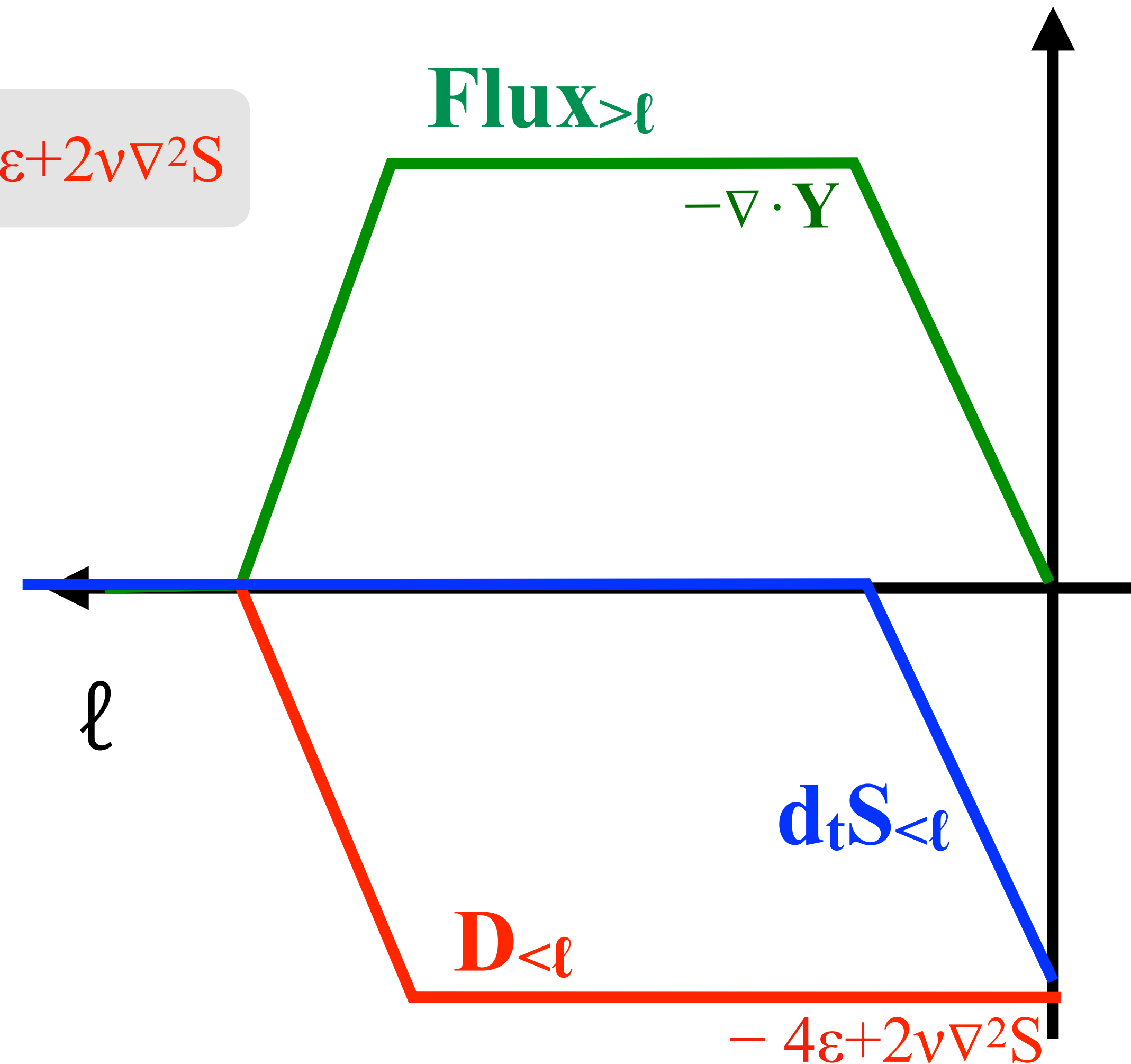
$$dS/dt = -\nabla \cdot \mathbf{Y} - 4\epsilon + 2\nu \nabla^2 S$$

dS/dt variation of energy contained in scales up to ℓ



Flux : energy flux from scales larger than ℓ

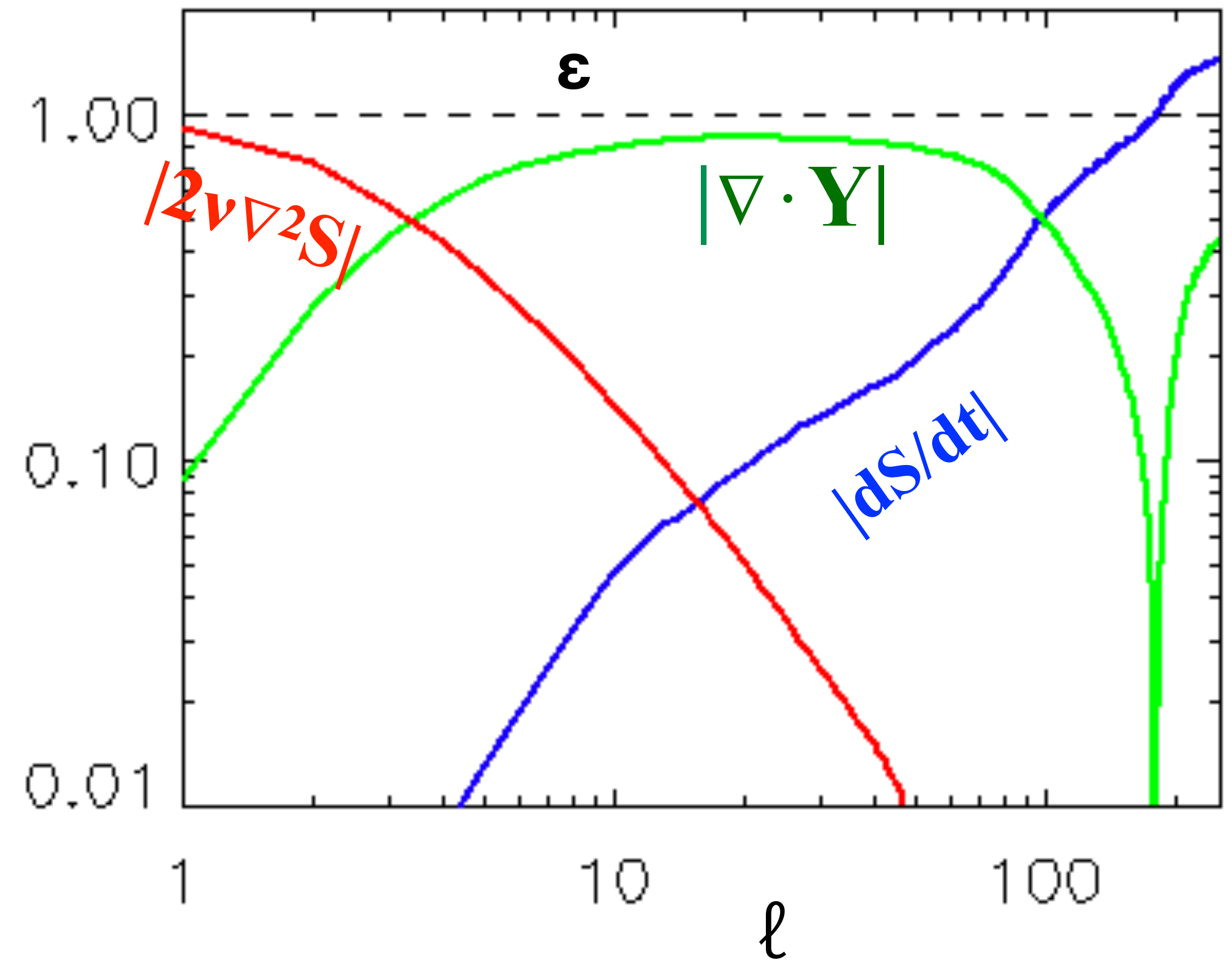
Diss : dissipation of energy contained in scales up to ℓ



Isotropic cascade

More popular representation

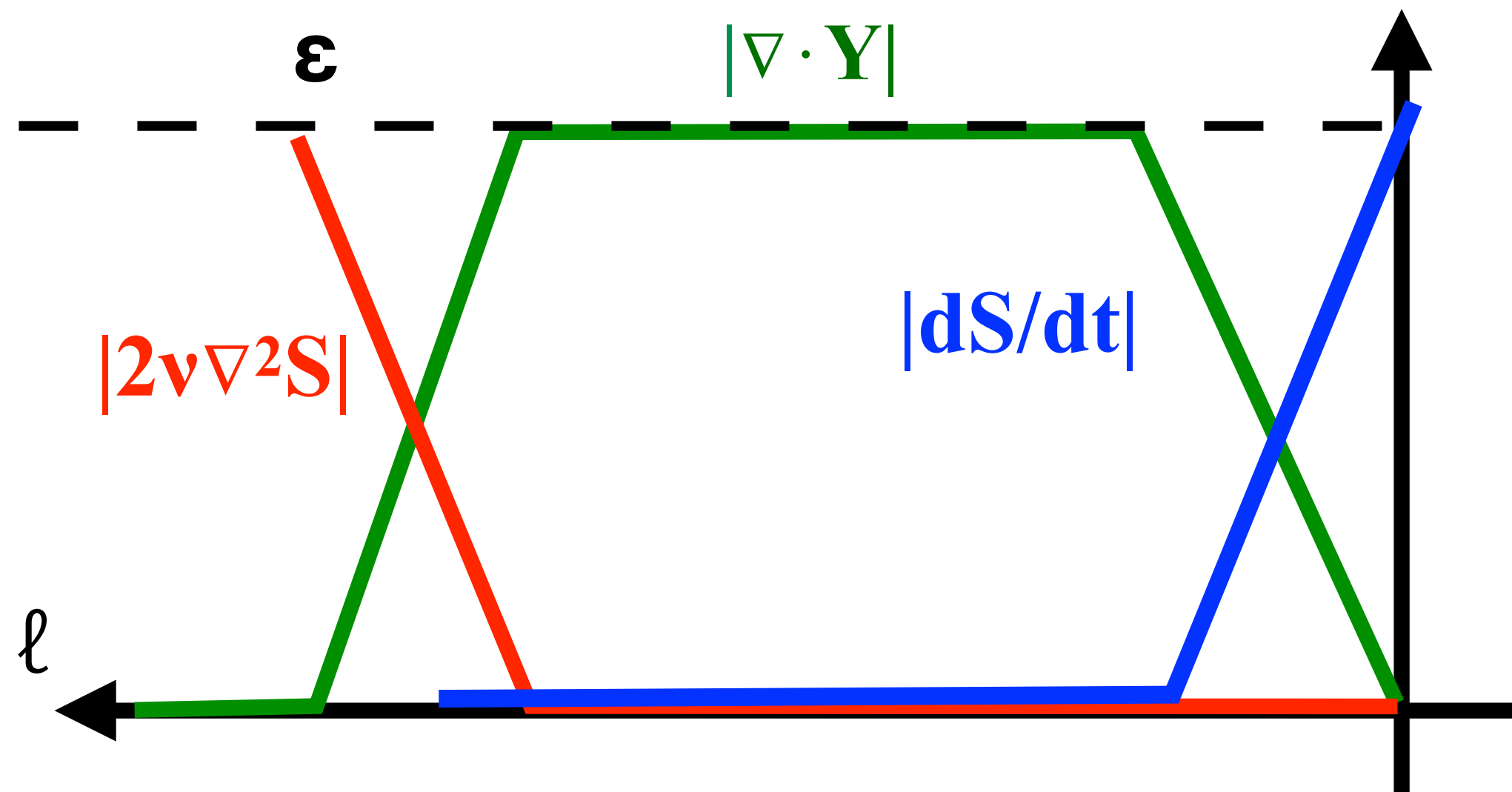
$$dS/dt + \nabla \cdot \mathbf{Y} - 2\nu \nabla^2 S = -4\varepsilon$$



Isotropic cascade

More popular representation

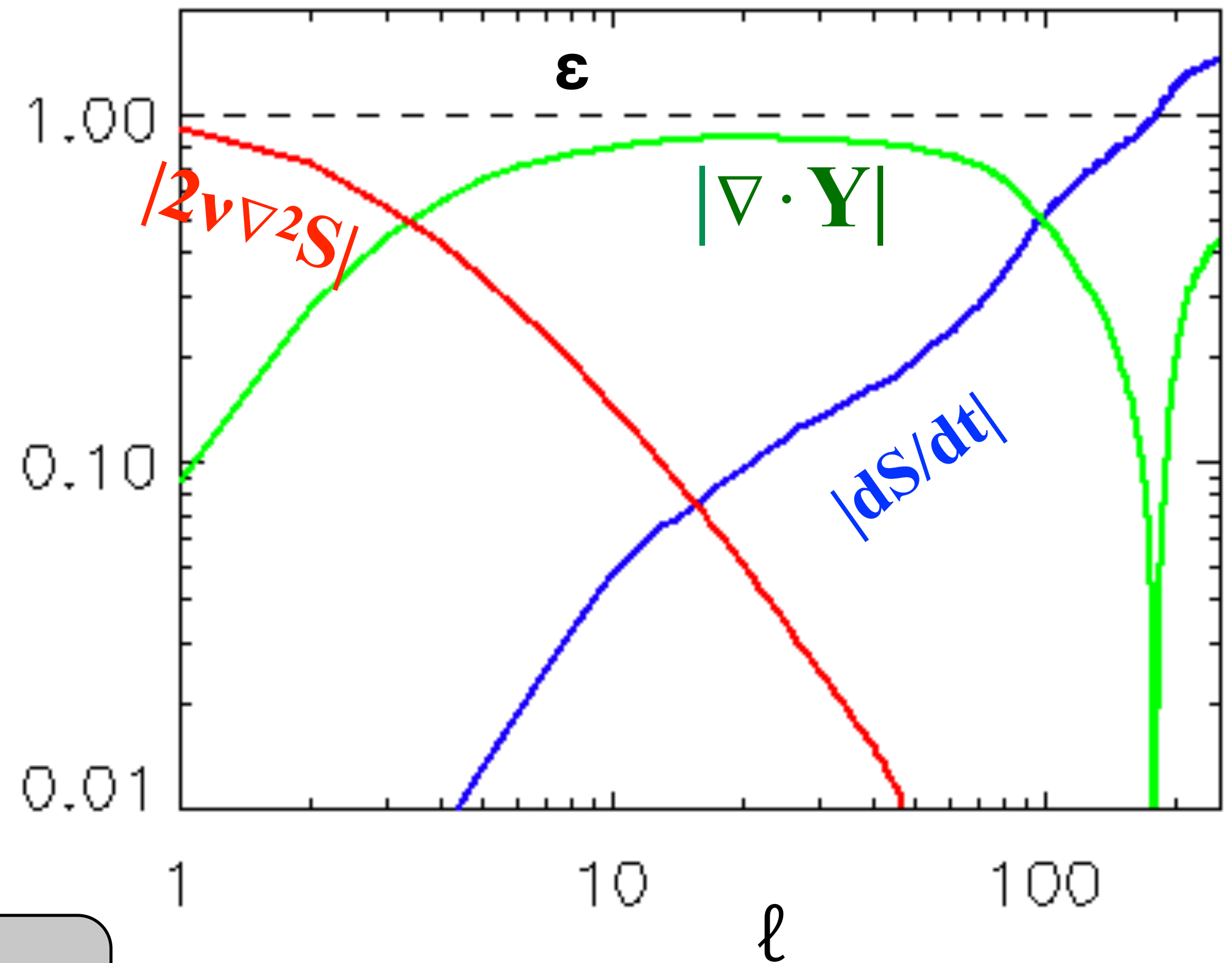
$$\frac{dS}{dt} + \nabla \cdot \mathbf{Y} - 2\nu \nabla^2 S = -4\epsilon$$



Ideally : when injection and dissipation are well separated

$$\nabla_{\ell} \cdot \mathbf{Y} = -4\epsilon$$

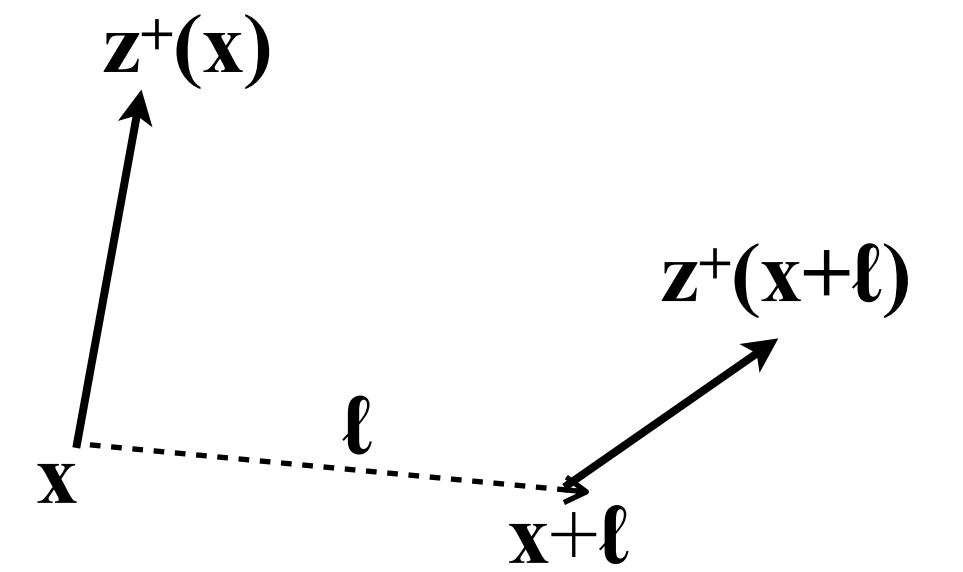
Politano Pouquet 1996 PRE



Cascade $\Leftrightarrow \nabla \cdot \mathbf{Y}$

Isotropic cascade

$$\partial_t S + \nabla_\ell \cdot \mathbf{Y} = -4\epsilon + 2\nu \nabla_\ell^2 S$$



$$\Delta \mathbf{z}^\pm(\mathbf{x}, \ell) = \mathbf{z}^\pm(\mathbf{x} + \ell) - \mathbf{z}^\pm(\mathbf{x})$$

II-order Structure Function

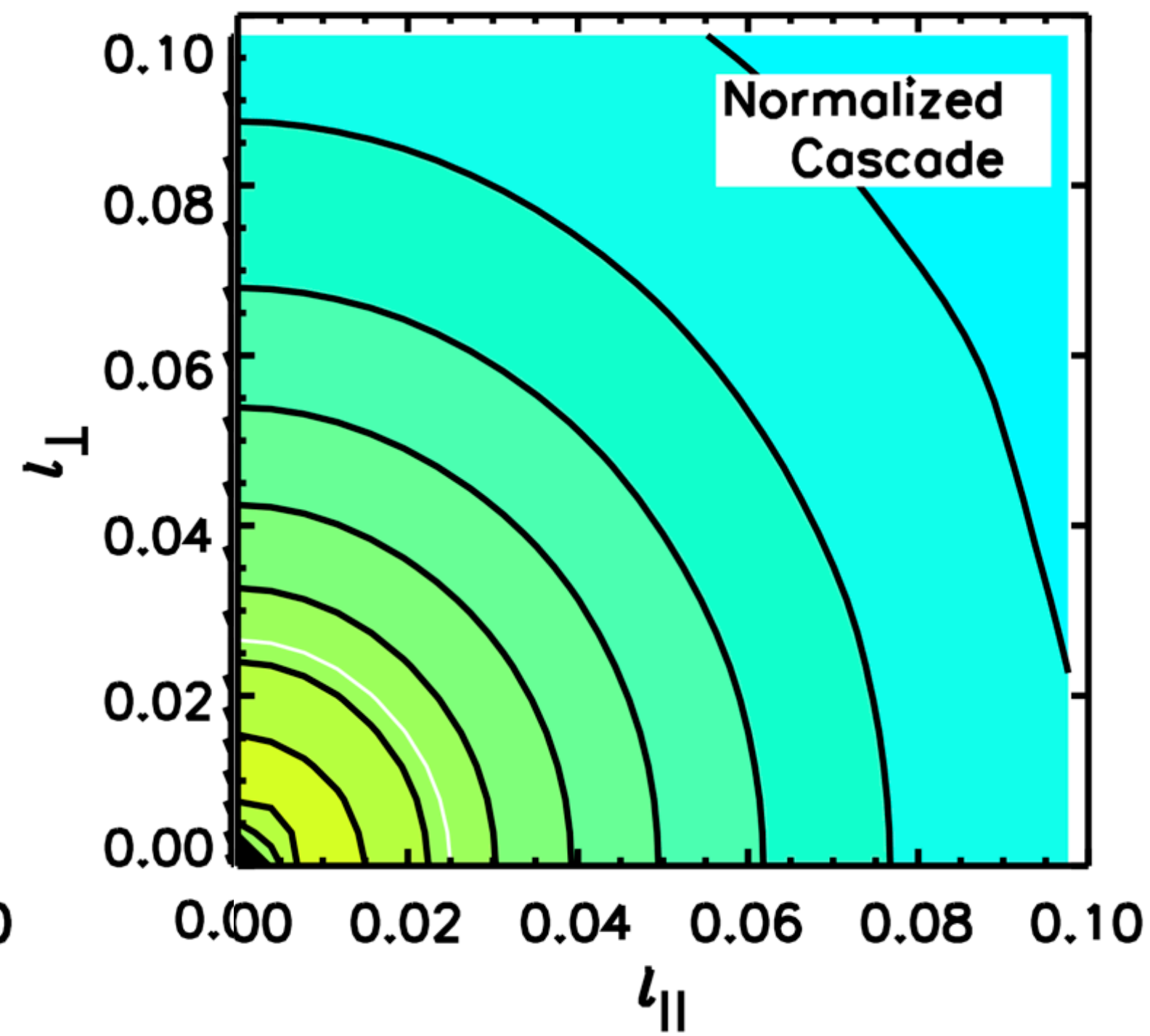
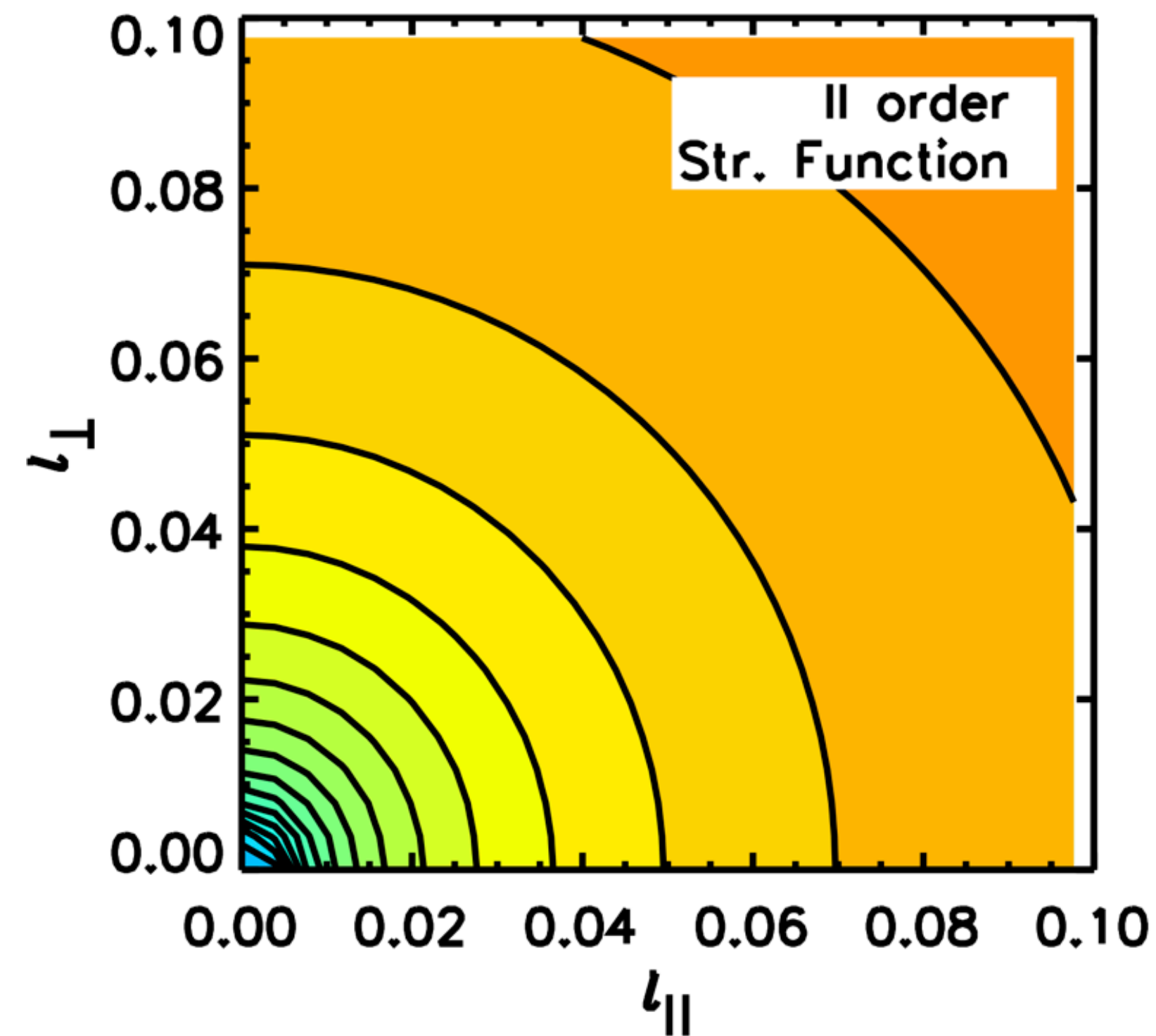
$$S = 1/2 [\langle |\Delta \mathbf{z}^-|^2 \rangle + \langle |\Delta \mathbf{z}^+|^2 \rangle]$$

Divergence of the flux

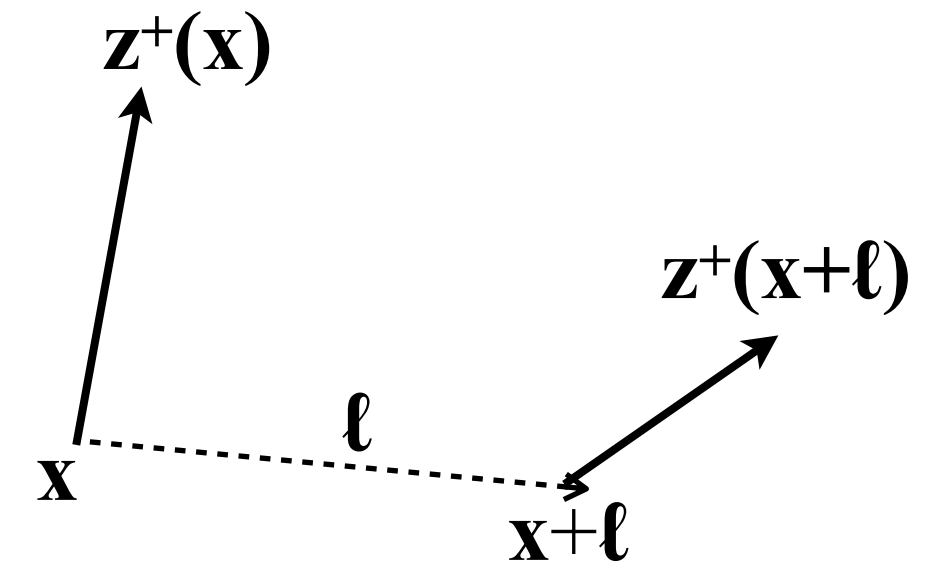
$$-\nabla \cdot \mathbf{Y} / 4\epsilon$$

III-order Structure Function

$$\mathbf{Y} = 1/2 [\langle \Delta \mathbf{z}^- |\Delta \mathbf{z}^+|^2 \rangle + \langle \Delta \mathbf{z}^+ |\Delta \mathbf{z}^-|^2 \rangle]$$



Isotropic cascade



$$\Delta \mathbf{z}^{\pm}(\mathbf{x}, \ell) = \mathbf{z}^{\pm}(\mathbf{x} + \ell) - \mathbf{z}^{\pm}(\mathbf{x})$$

$$\partial_t S + \nabla_{\ell} \cdot \mathbf{Y} = -4\epsilon + 2\nu \nabla_{\ell}^2 S$$

II-order Structure Function

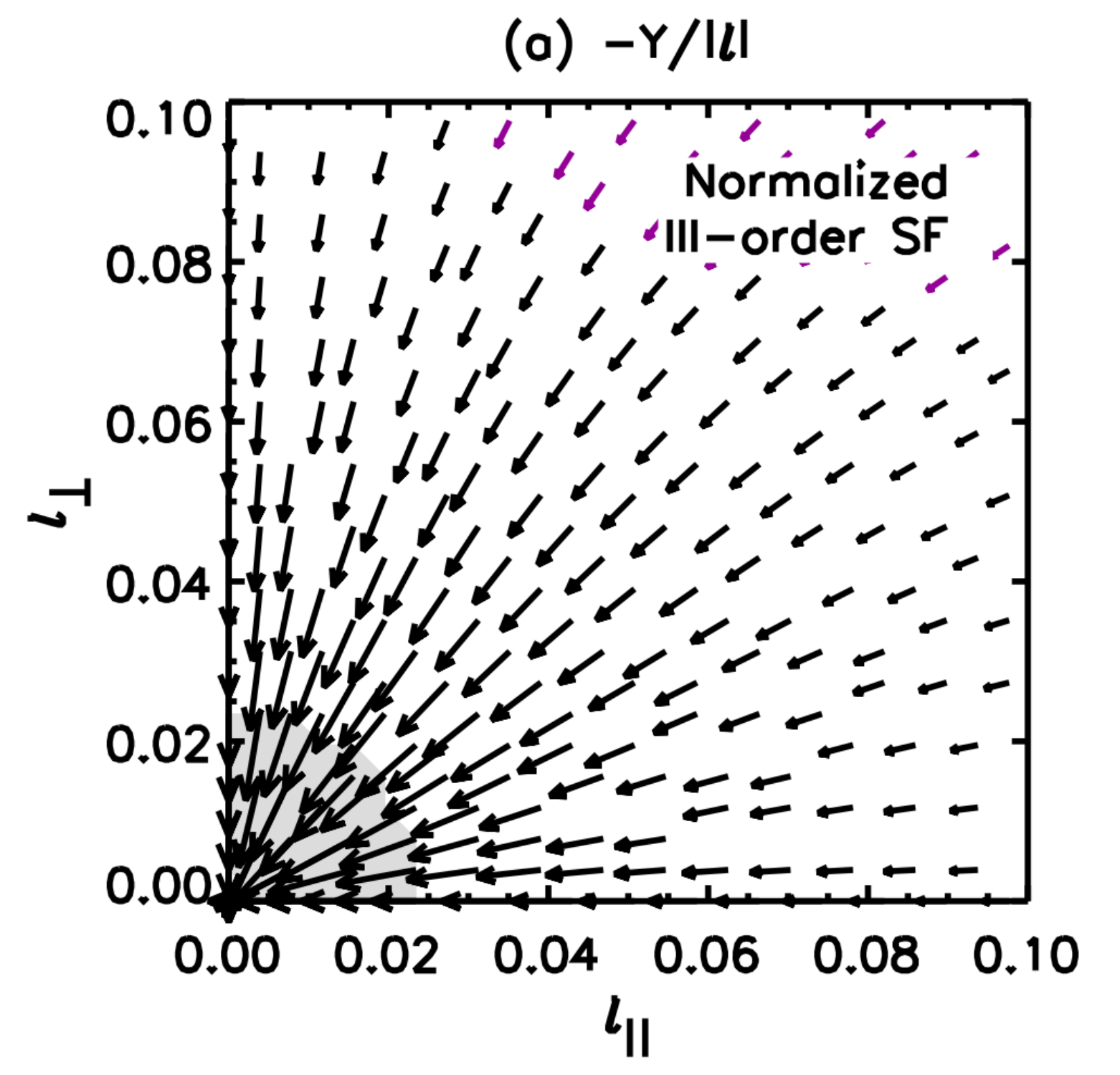
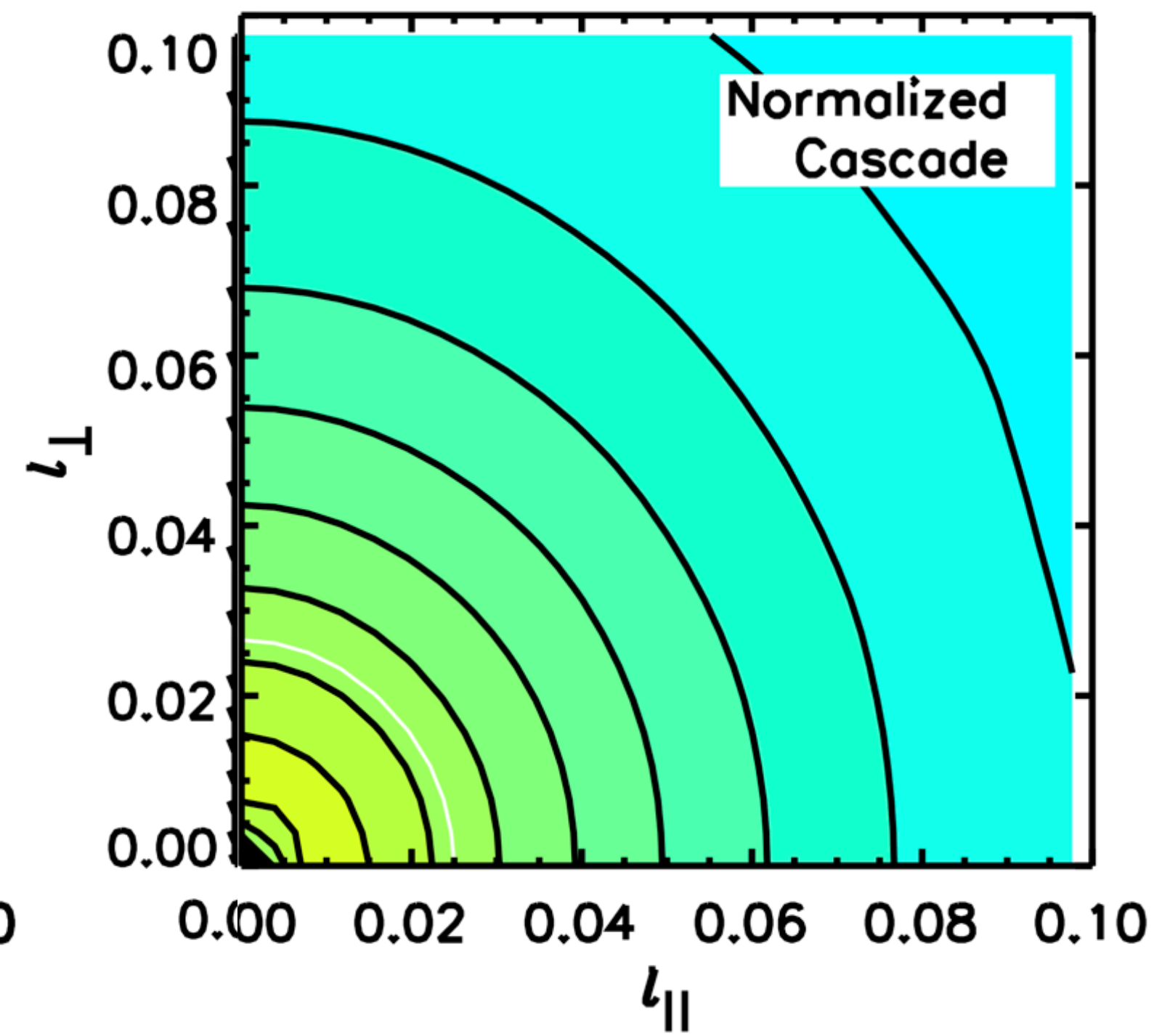
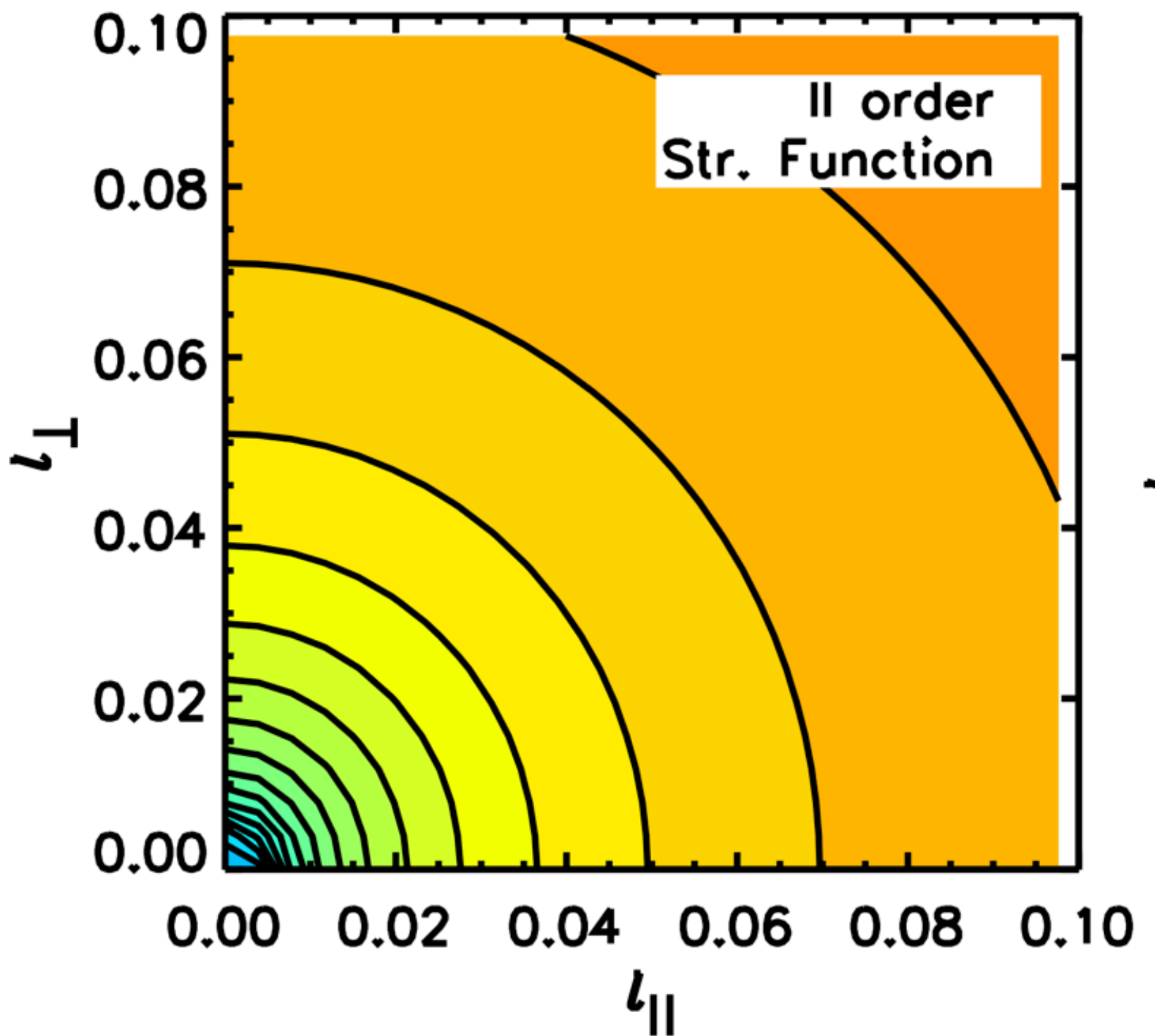
$$S = 1/2 [\langle |\Delta \mathbf{z}^{-}|^2 \rangle + \langle |\Delta \mathbf{z}^{+}|^2 \rangle]$$

Divergence of the flux

$$-\nabla \cdot \mathbf{Y} / 4\epsilon$$

III-order Structure Function

$$\mathbf{Y} = 1/2 [\langle \Delta \mathbf{z}^{-} |\Delta \mathbf{z}^{+}|^2 \rangle + \langle \Delta \mathbf{z}^{+} |\Delta \mathbf{z}^{-}|^2 \rangle]$$



Meaning of \mathbf{Y}

$$\partial_t S + \nabla_\ell \cdot \mathbf{Y} = -4\epsilon + 2\nu \nabla_\ell^2 S$$

WE DO NOT KNOW

III-order Structure Function

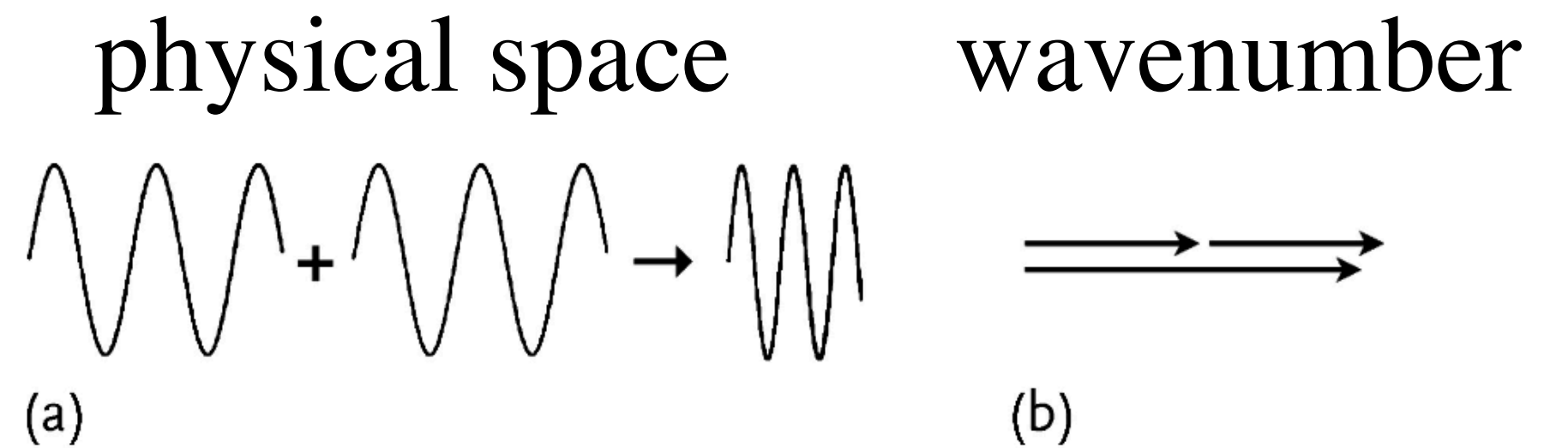
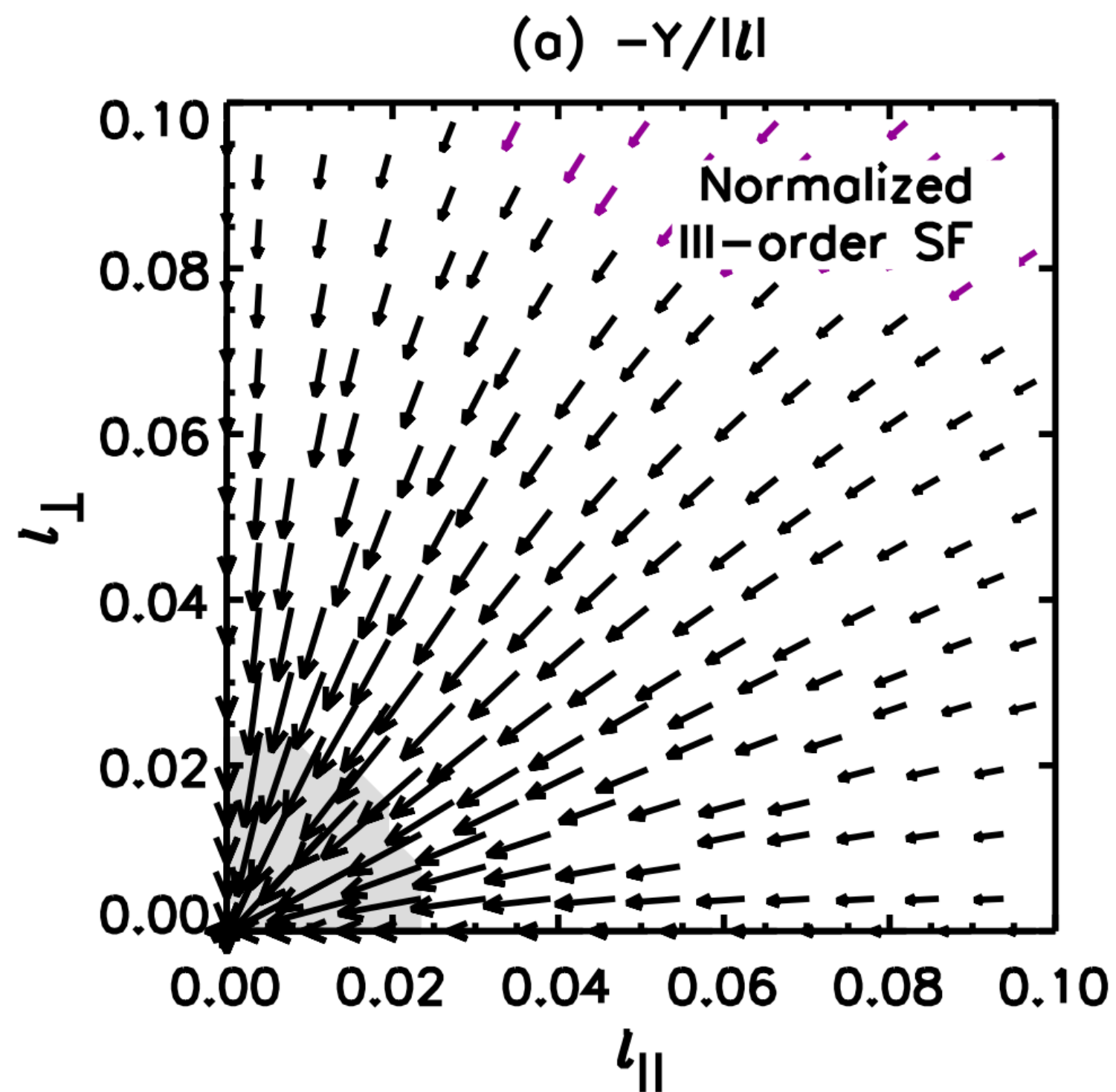
$$\mathbf{Y} = 1/2 [\langle \Delta \mathbf{z}^- | \Delta \mathbf{z}^+|^2 \rangle + \langle \Delta \mathbf{z}^+ | \Delta \mathbf{z}^-|^2 \rangle]$$

$$C(\ell) = \langle u_i(\mathbf{x} + \ell) u_i(\mathbf{x}) \rangle$$

$$SF = 2[C(0) - C(\ell)]$$

$$SF(\ell) = \langle |u_i(\mathbf{x} + \ell) - u_i(\mathbf{x})|^2 \rangle$$

$$\partial_t C - \nabla_\ell \cdot \mathbf{Y} = -2\nu \nabla_\ell^2 C,$$

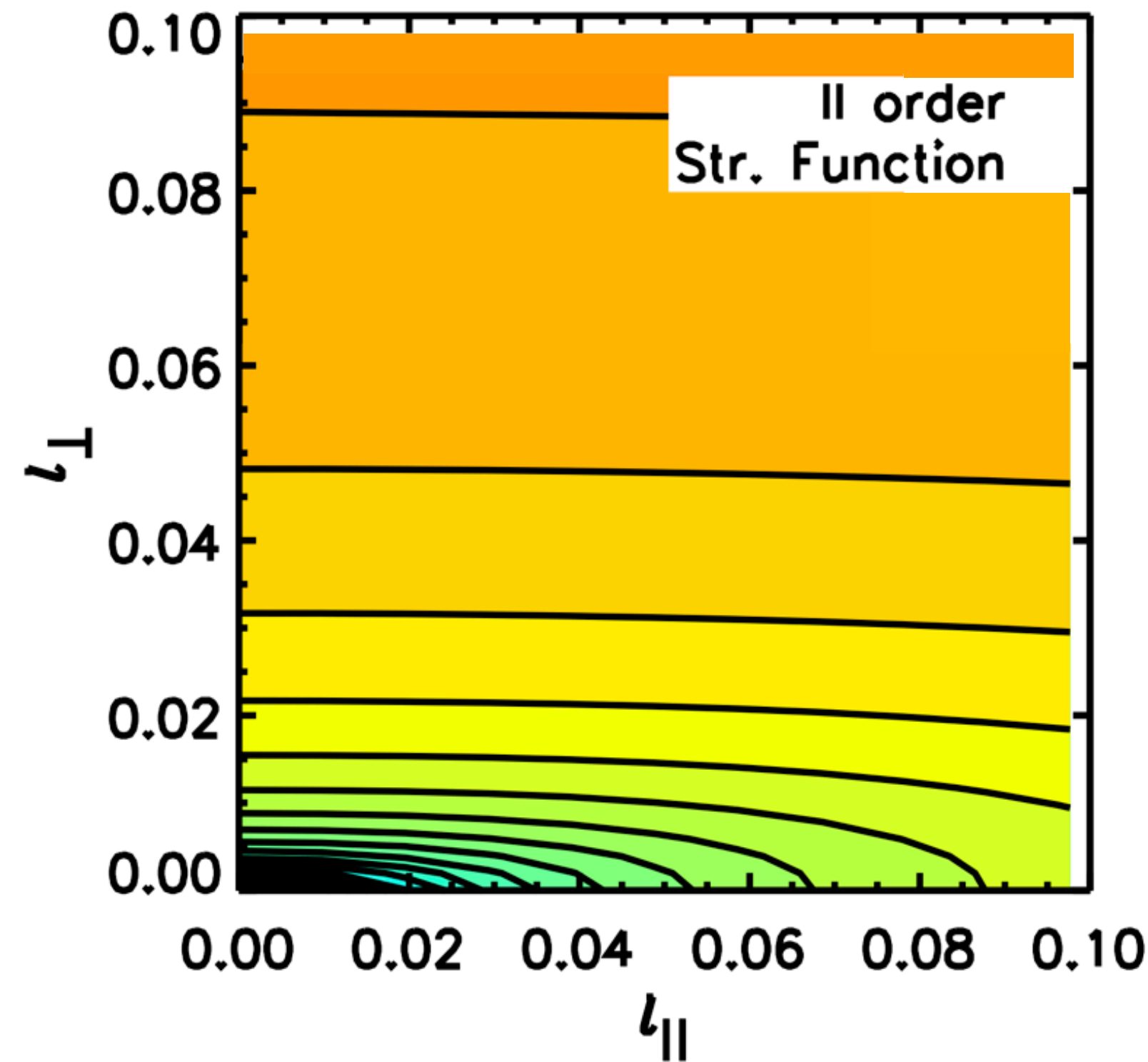


\mathbf{Y} brings negative correlations toward small scales through spherical surfaces (iso surfaces of $\nabla \cdot \mathbf{Y}$)

Anisotropic cascade

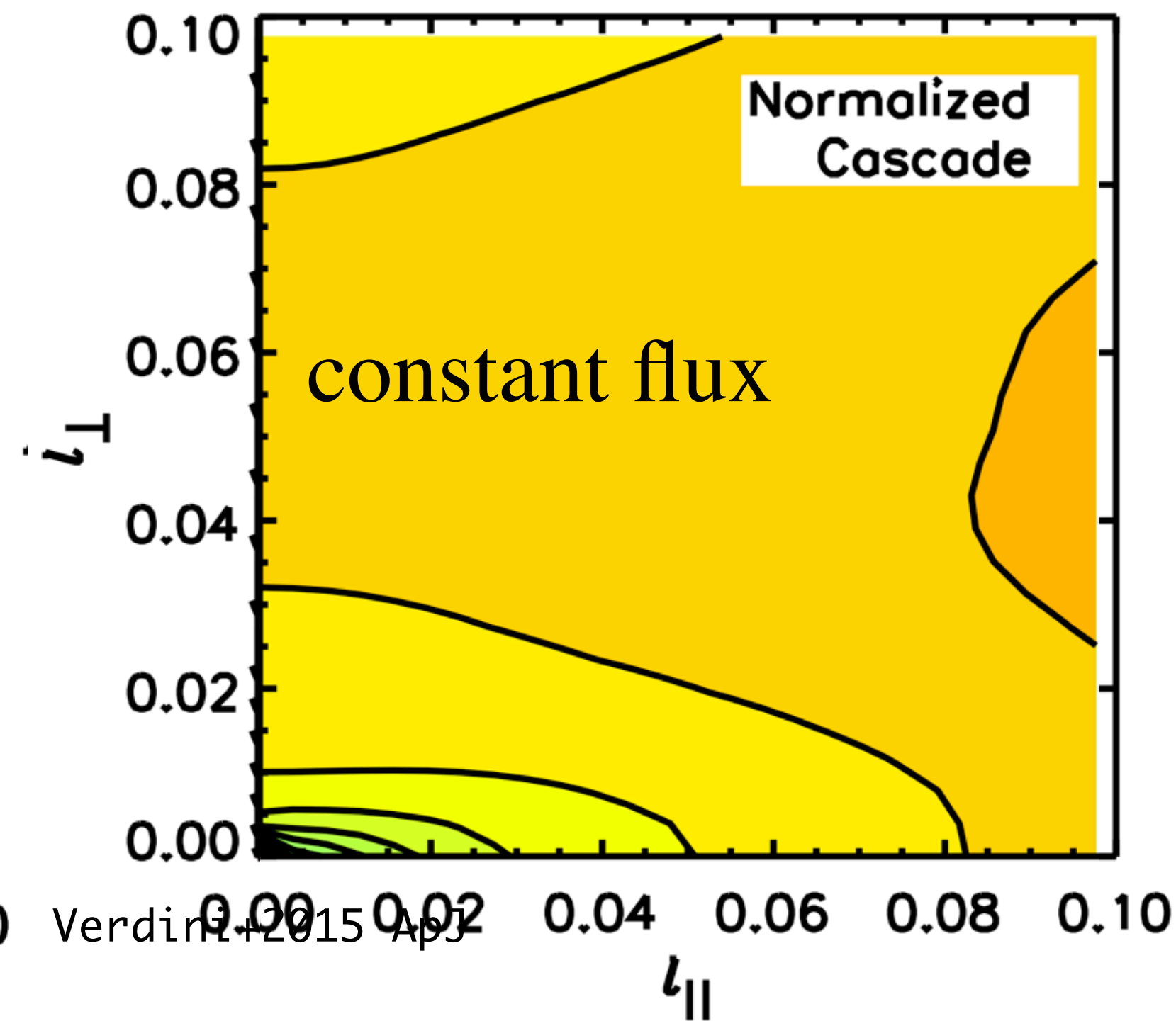
Anisotropy shows up in the difference in \parallel and \perp ranges for which $\nabla \cdot \mathbf{Y} = \text{const}$

II-order Structure Function



Divergence of the flux

$$-\nabla \cdot \mathbf{Y} / 4\varepsilon$$

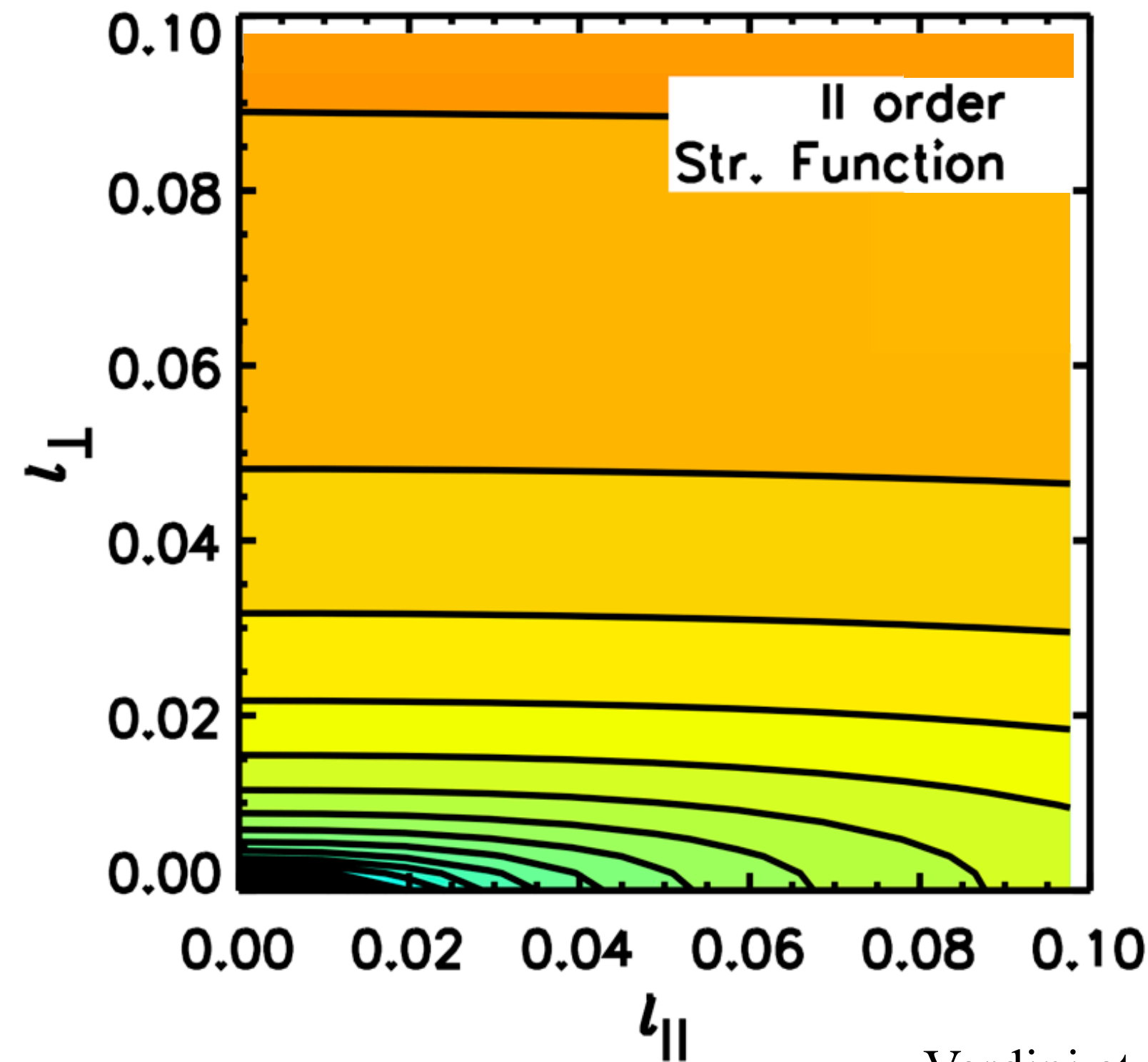


Anisotropic cascade

Anisotropy shows up in the difference in \parallel and \perp ranges for which $\nabla \cdot \mathbf{Y} = \text{const}$

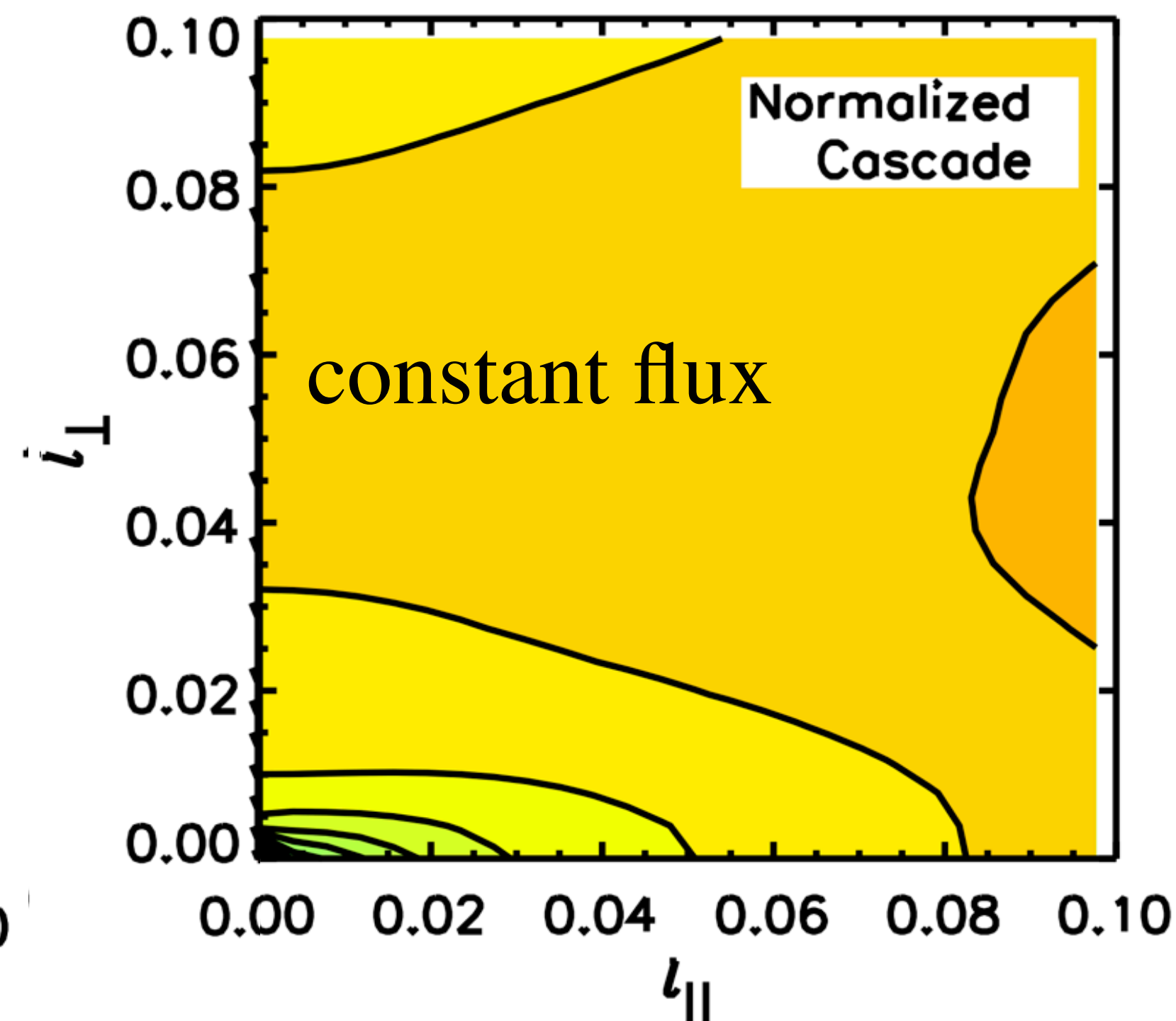
\mathbf{Y} is difficult to interpret without knowing the surfaces of constant divergence

II-order Structure Function



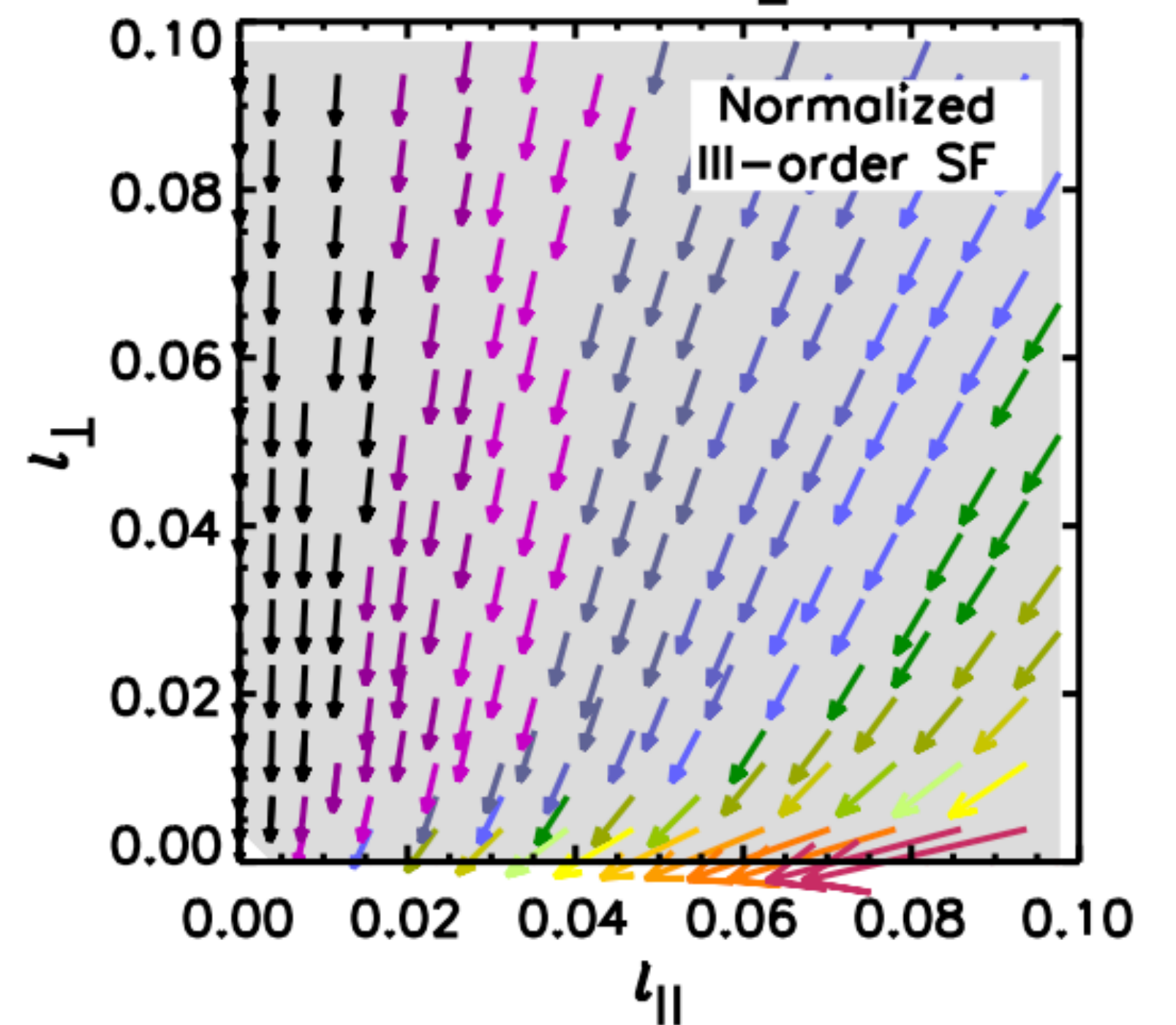
Divergence of the flux

$$-\nabla \cdot \mathbf{Y} / 4\varepsilon$$

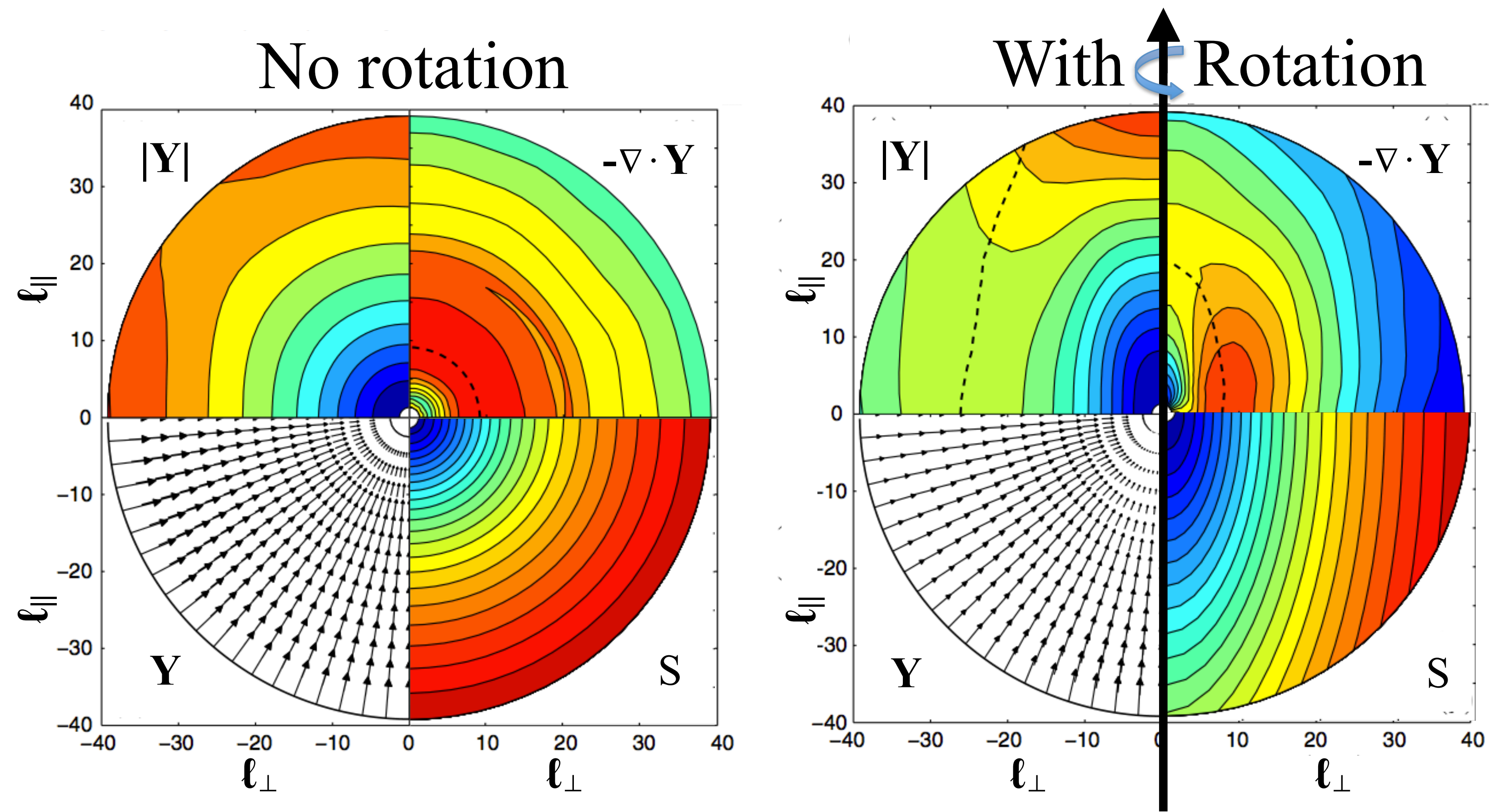


III-order Structure Function

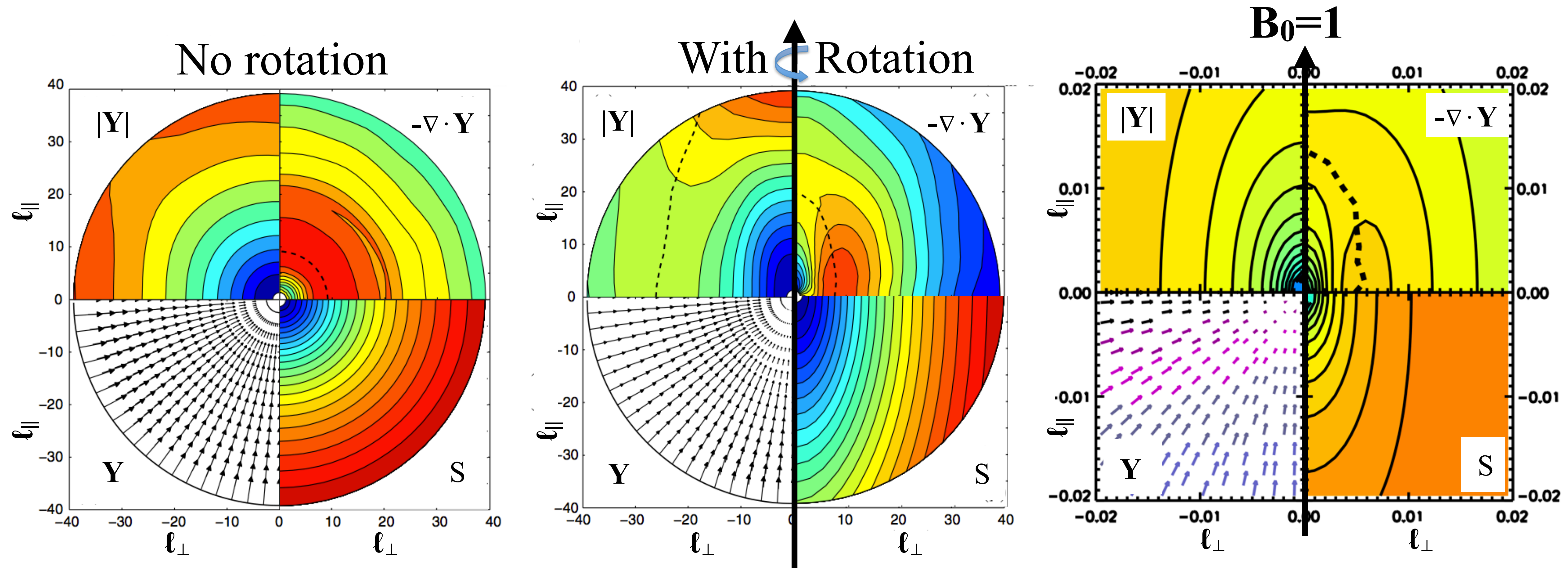
(a) $-\mathbf{Y} / l_{\perp}$



Hydrodynamics with Rotation



Hydrodynamics with Rotation

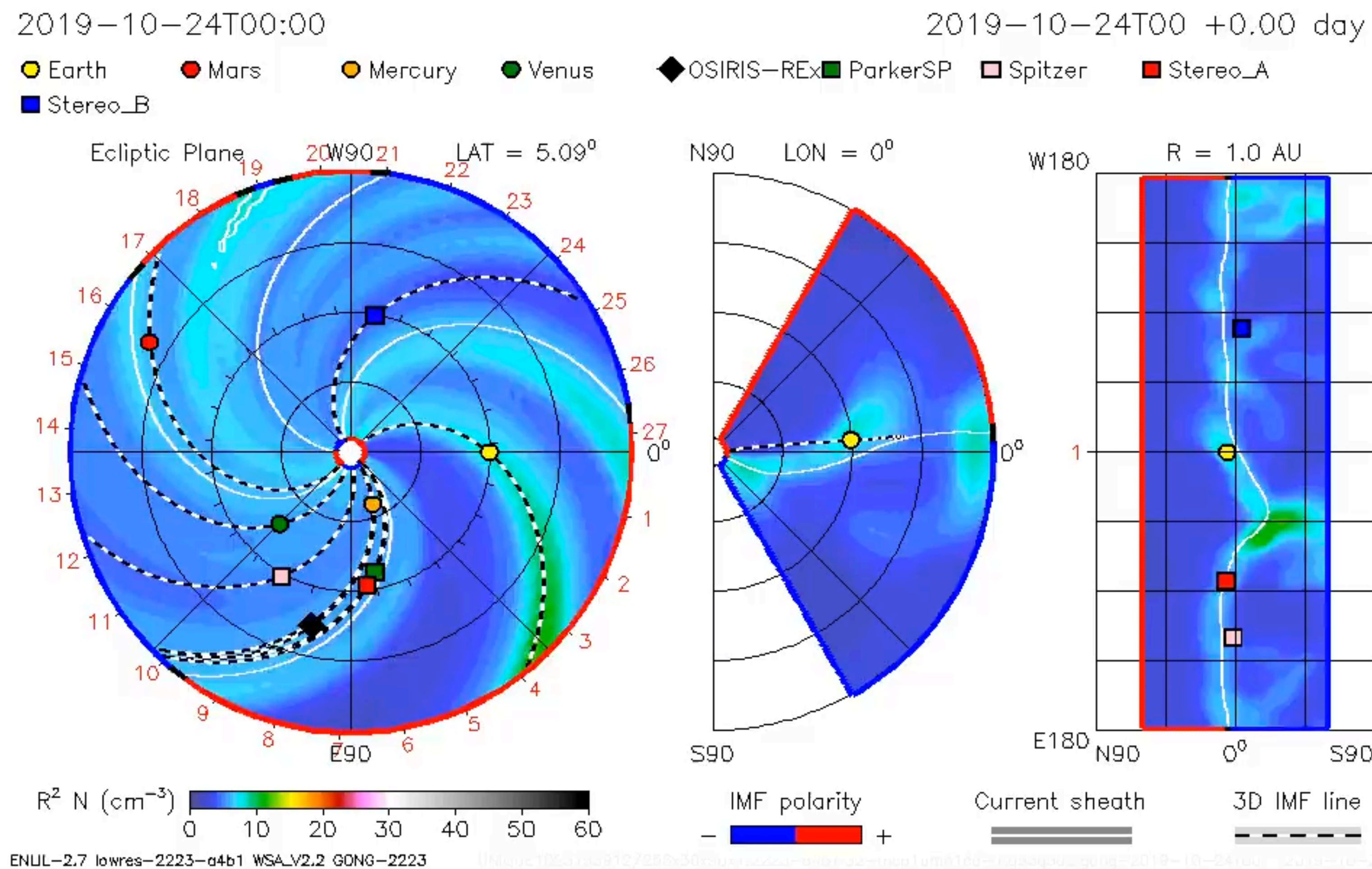


Solar Wind Measurement: Exploits Variability

Solar wind flows almost radially
(at least from 0.3 au)

Magnetic field follows a spiral
whose inclination depends on the
speed (faster, more radial)
On the ecliptic opposite sectors
cross each other due to mis-
alignment of magnetic and
rotation axis

Many samples to cover different
 θ_{BV}



Measure of anisotropy

Consider a timeseries in which B_0 has direction w.r.t. to the sampling direction $R \equiv V_{sw}$, i.e. a given θ_{BV}

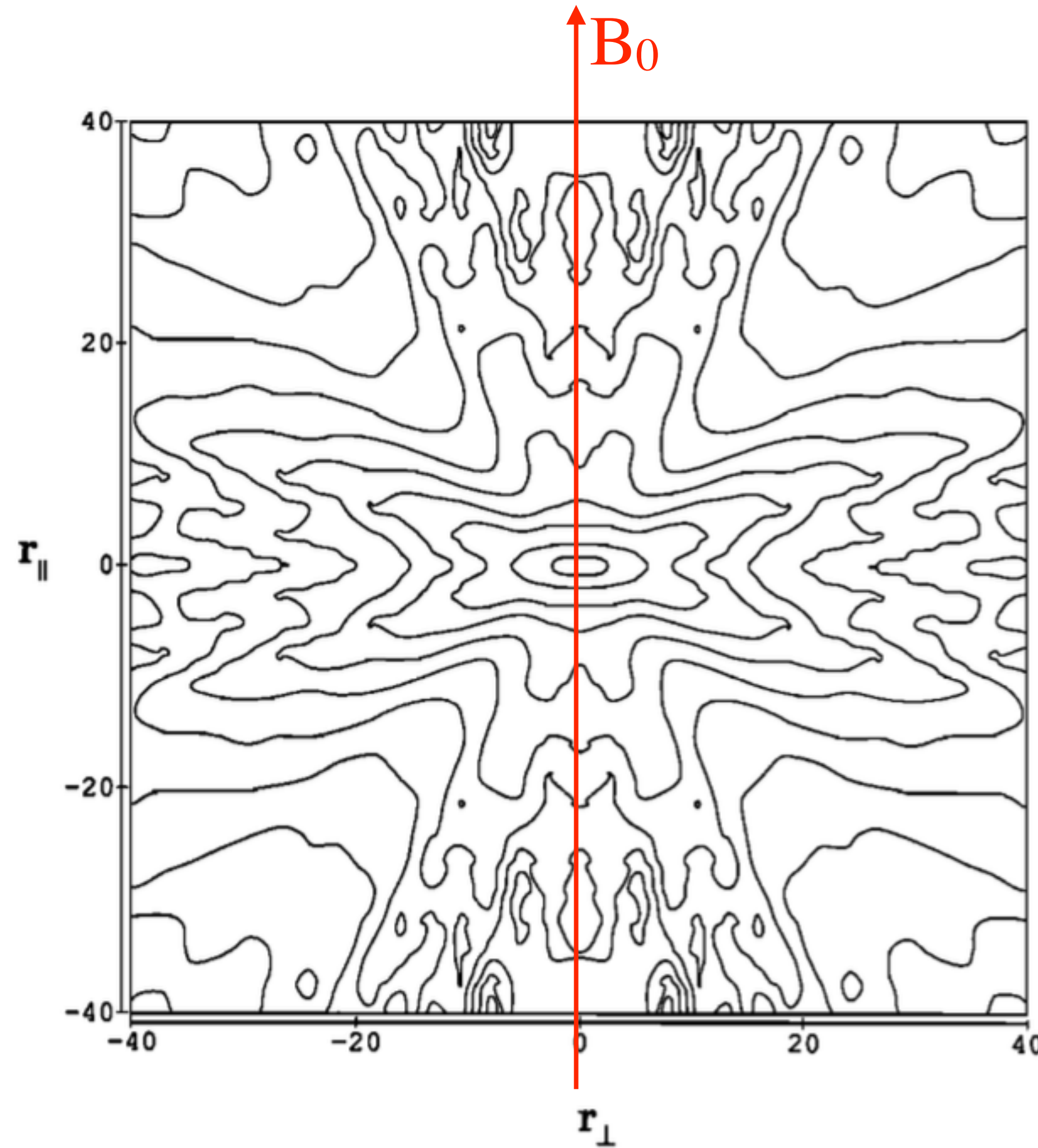
Remove Average

Compute autocorrelation

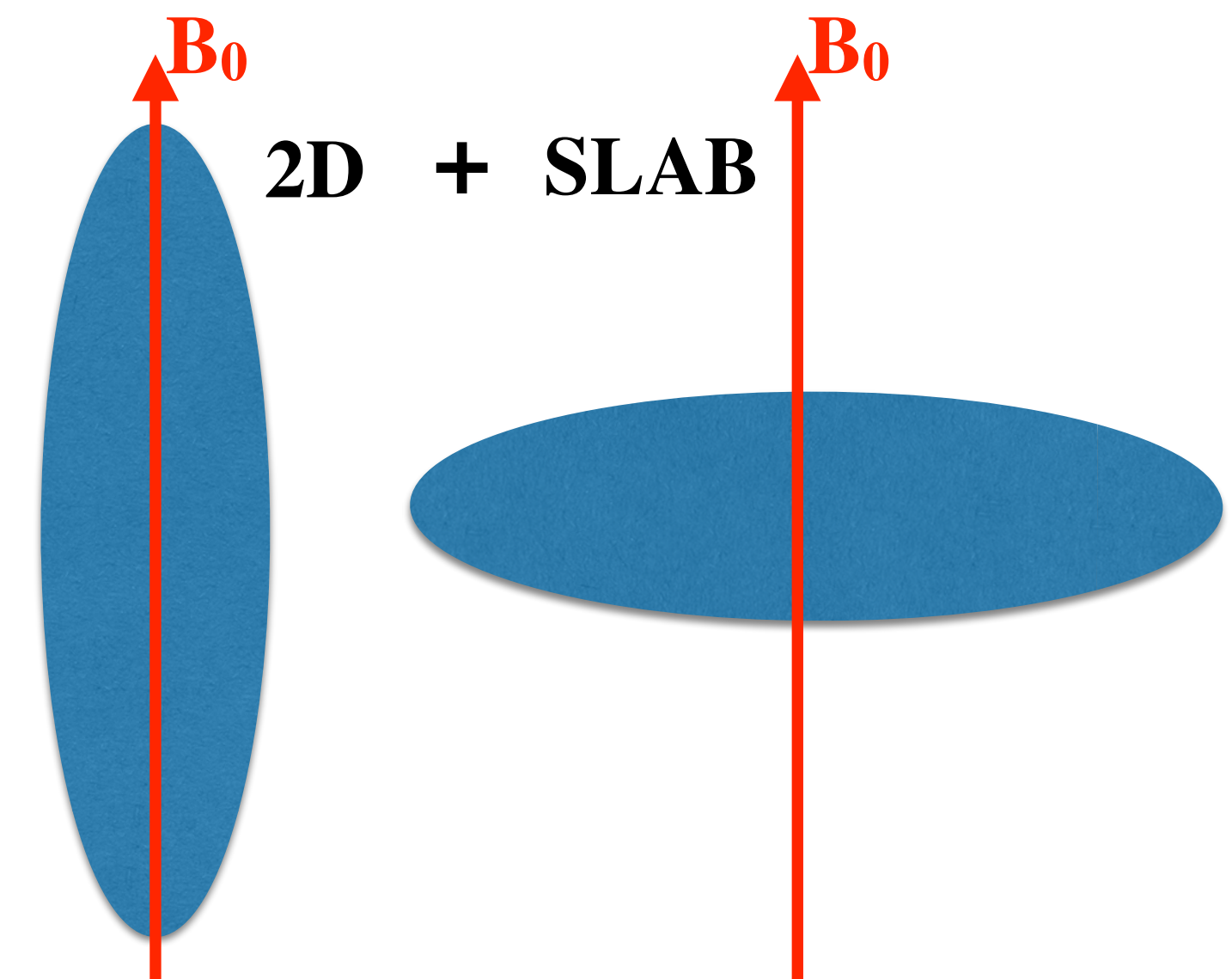
$$\ell = V_{sw} \cdot \tau$$

Normalize by $C(0) = \delta u^2$

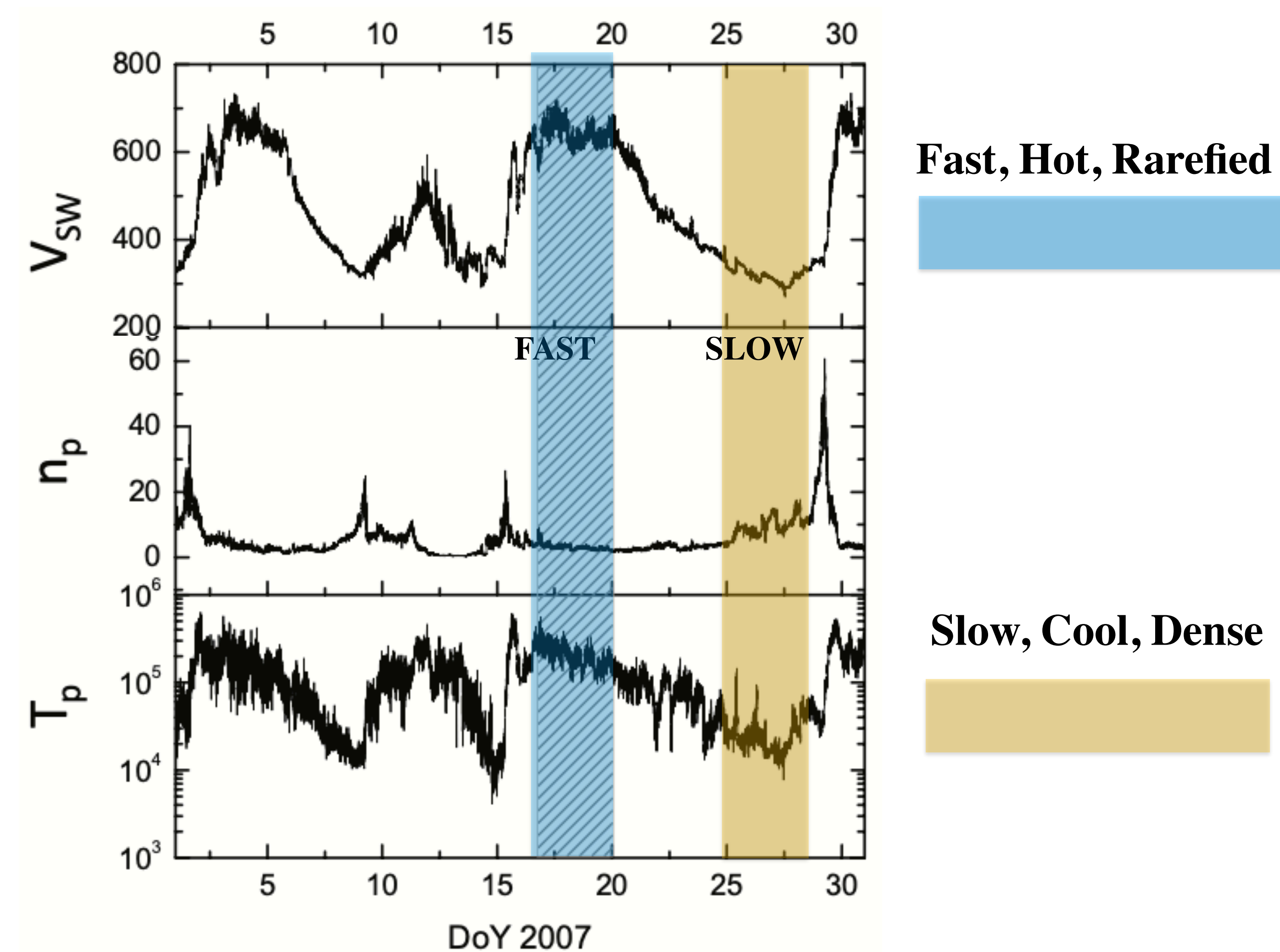
Bin AC according to θ_{BV}



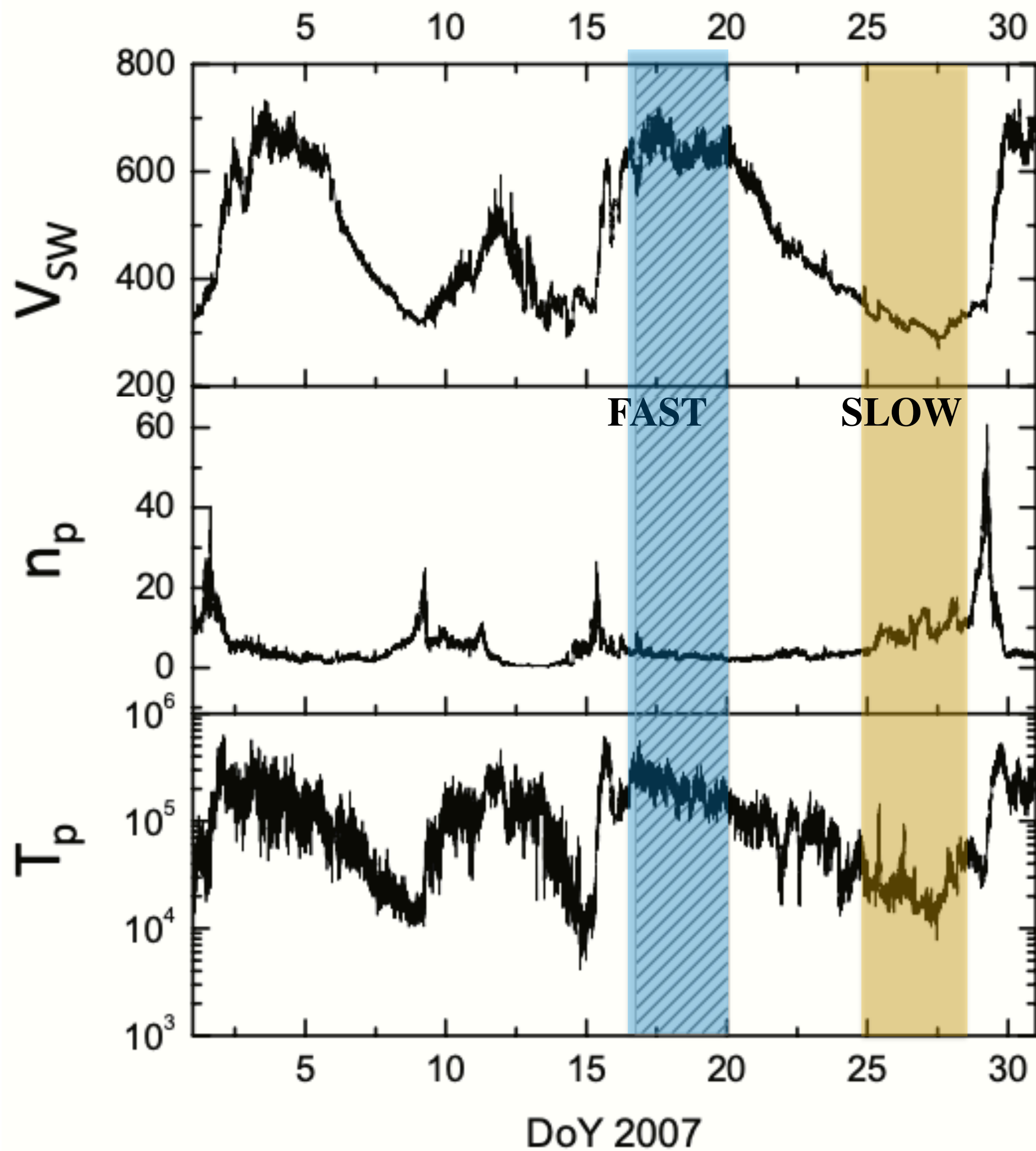
Maltese Cross



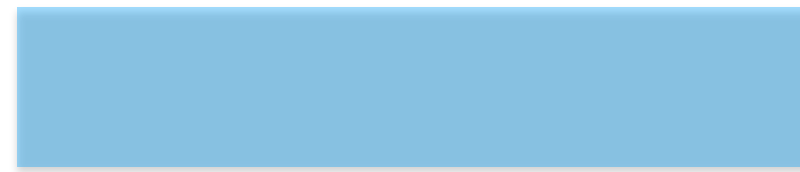
Wind Speed and Turbulence



Wind Speed and Turbulence



Fast, Hot, Rarefied

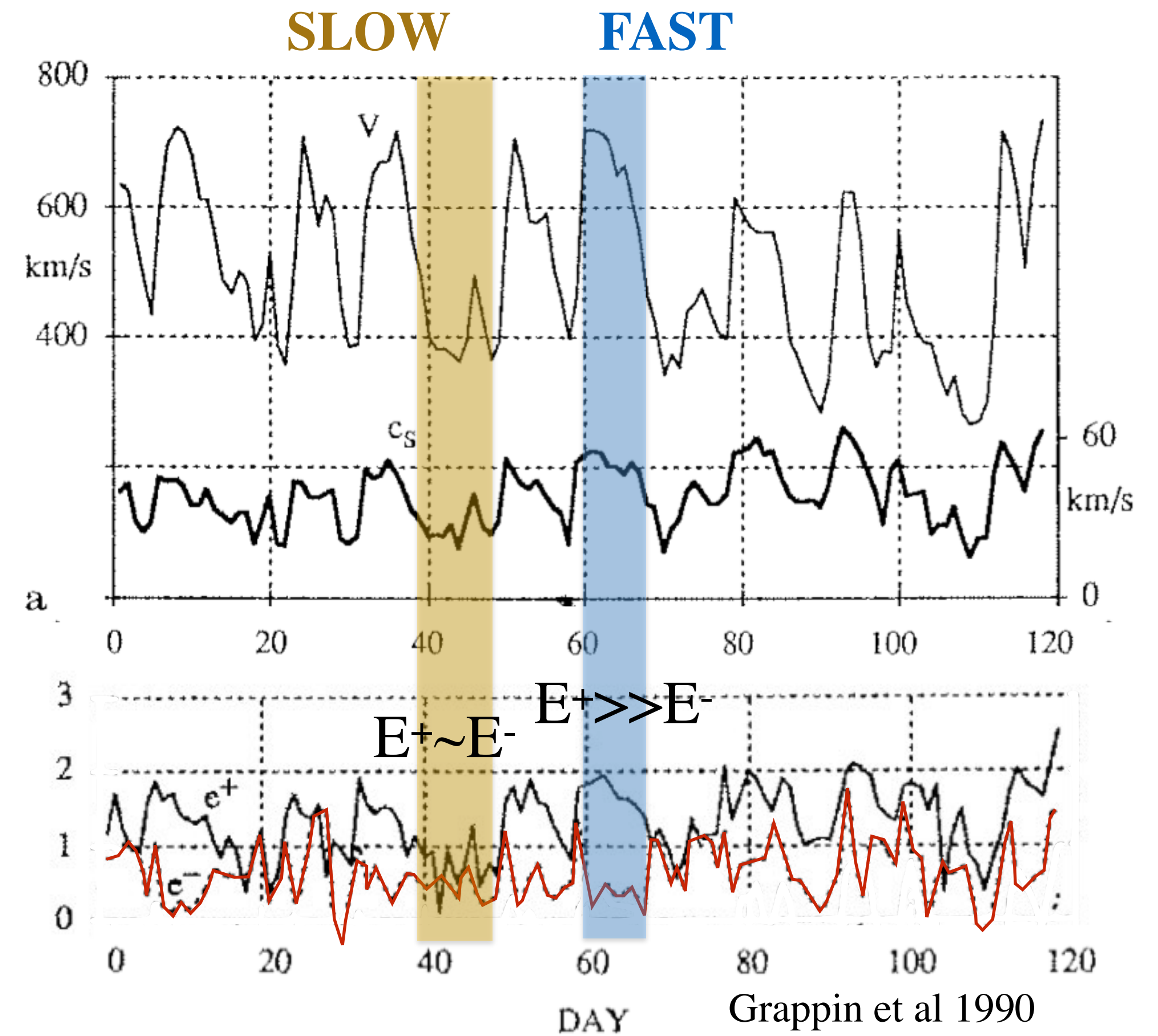


Alfvénic and almost incompressible

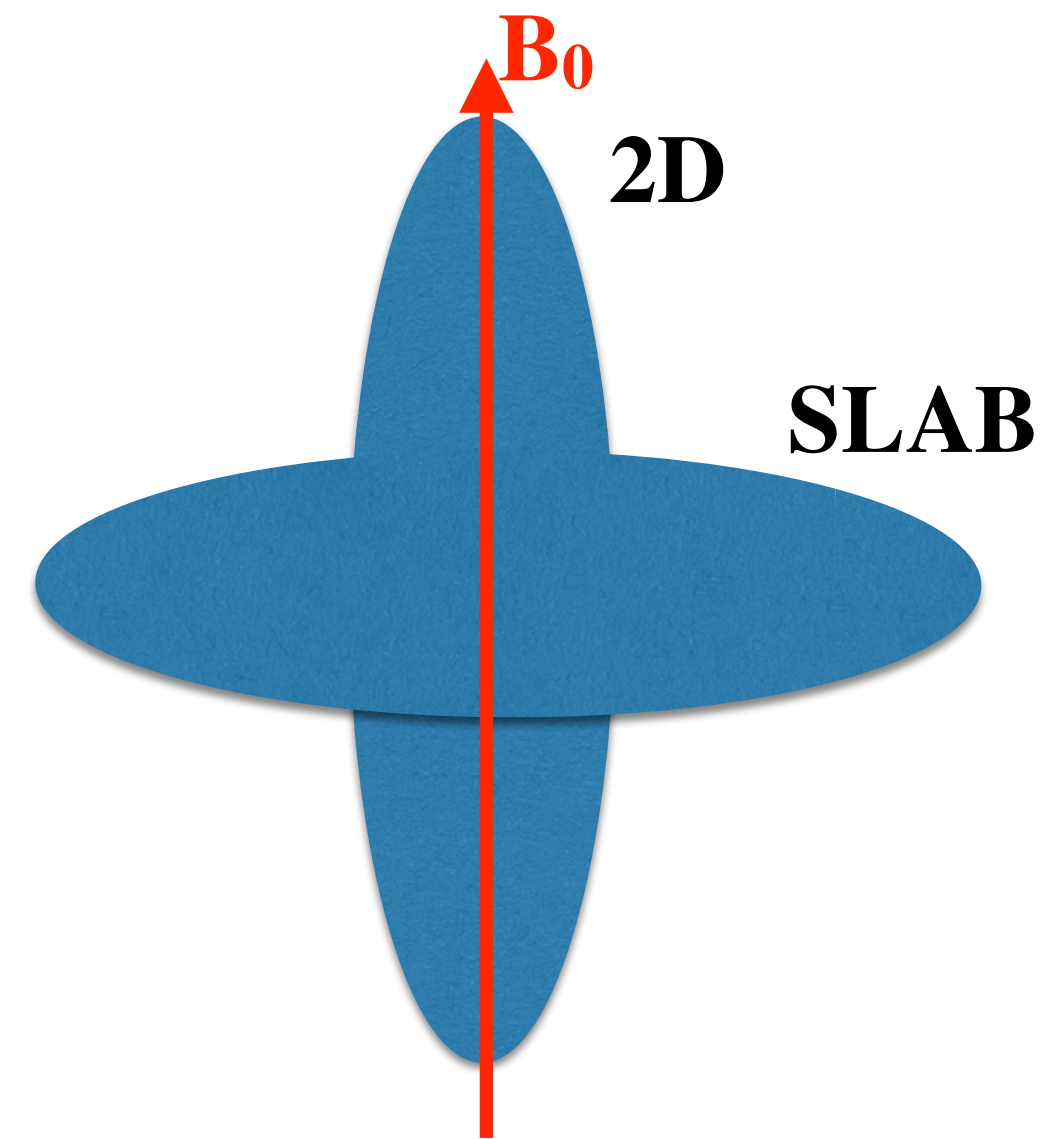
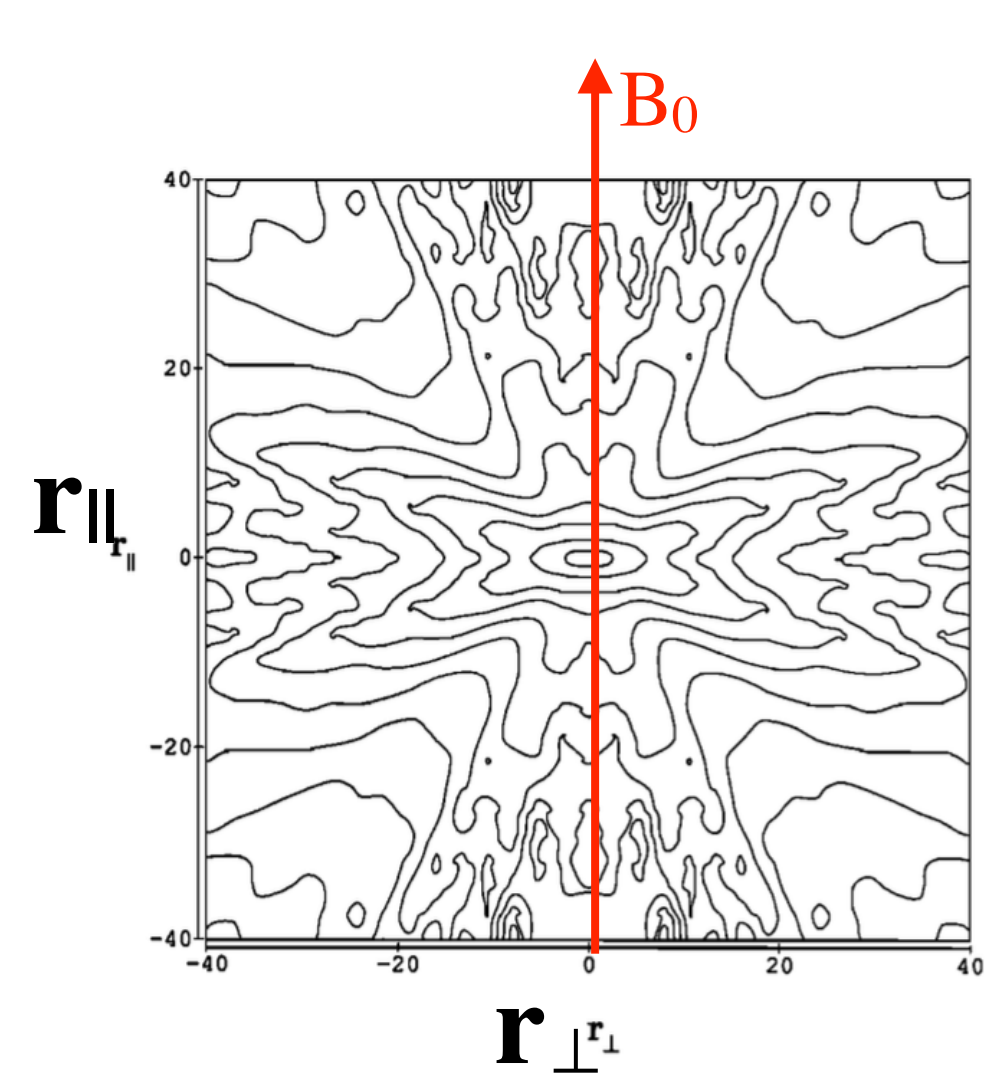
Slow, Cool, Dense



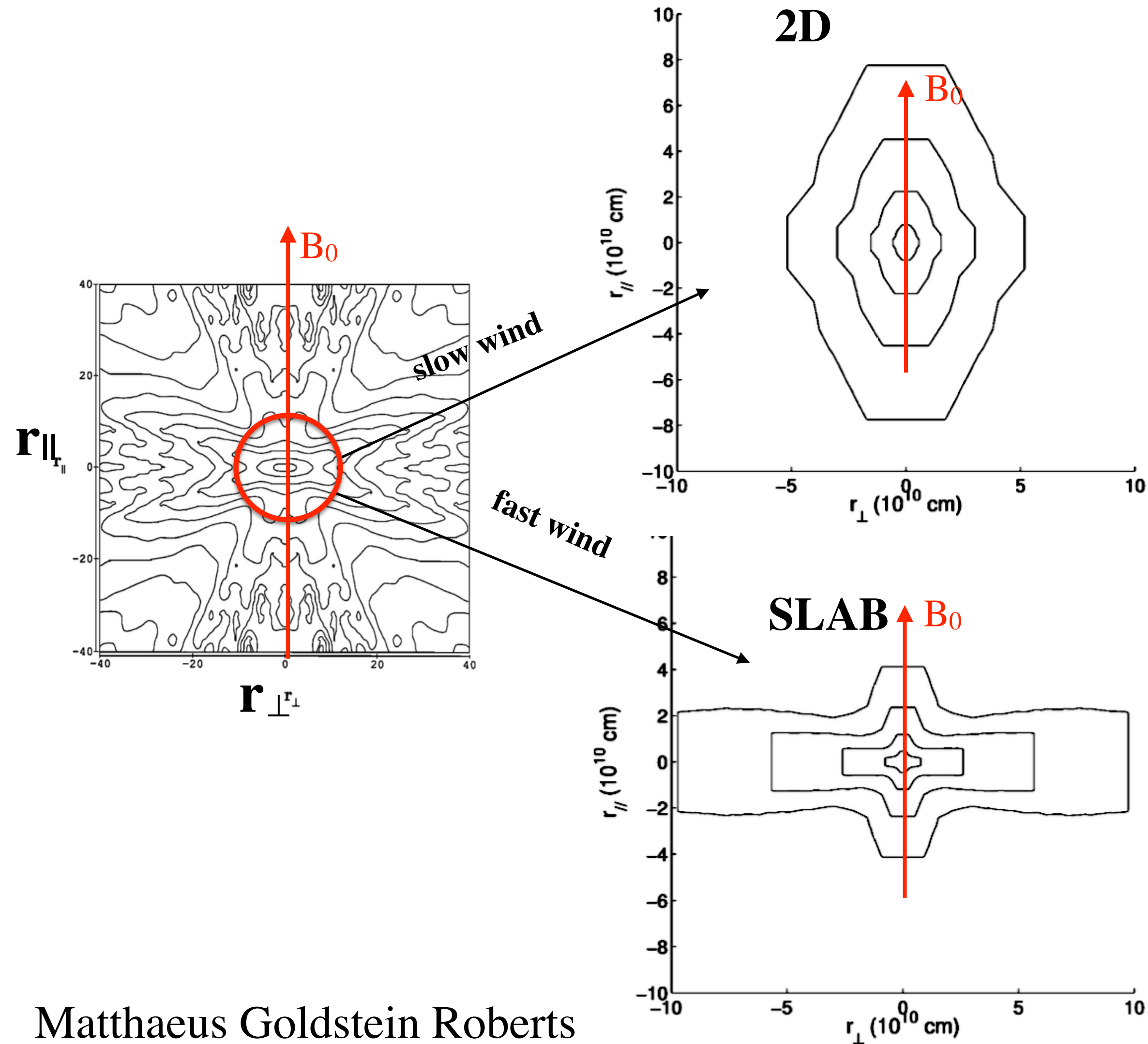
Non Alfvénic and compressible



Anisotropy at smaller scale



Anisotropy at smaller scale



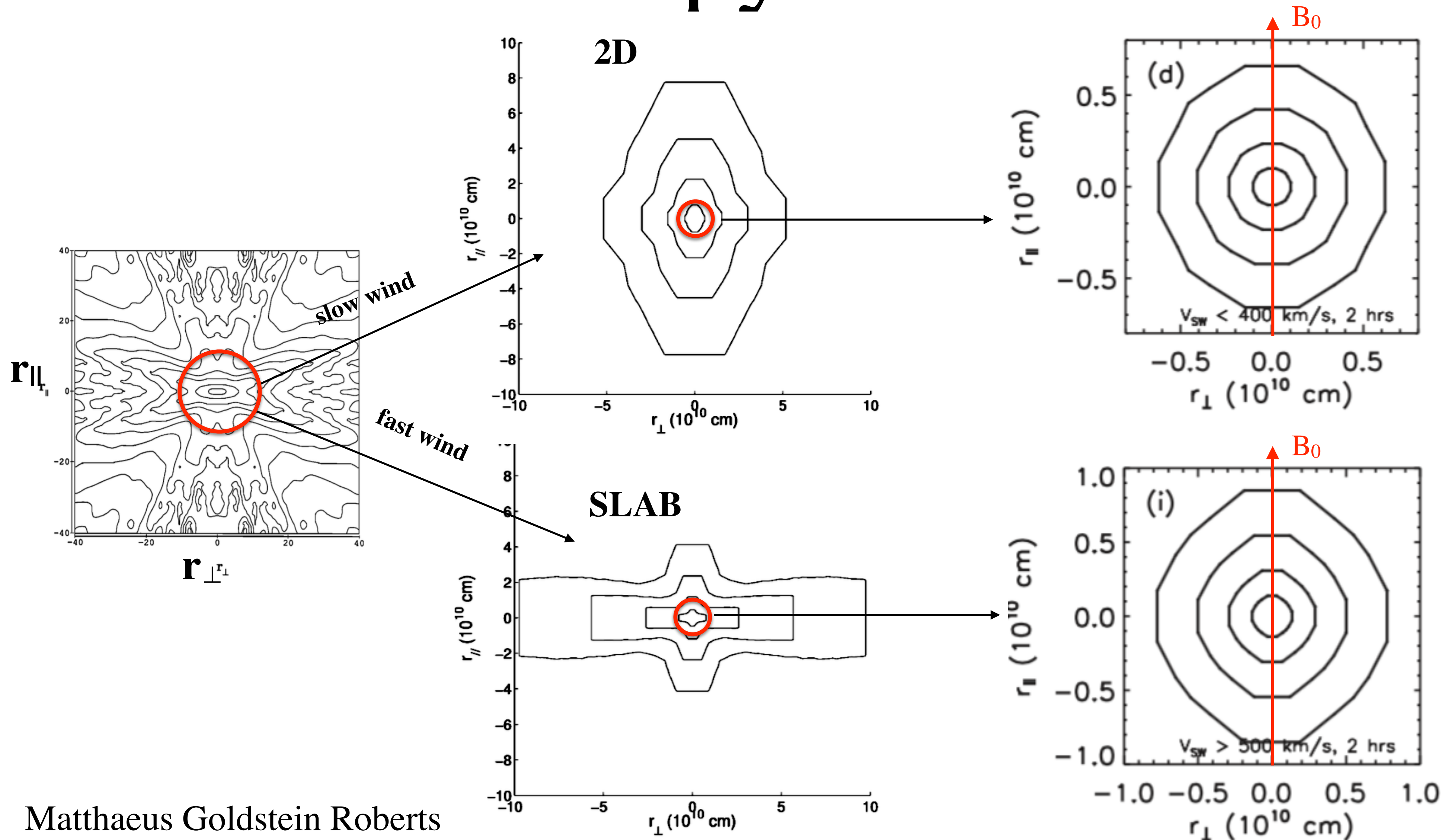
Slow is mainly 2D

Energy in k_{\perp}

Fast is mainly Slab

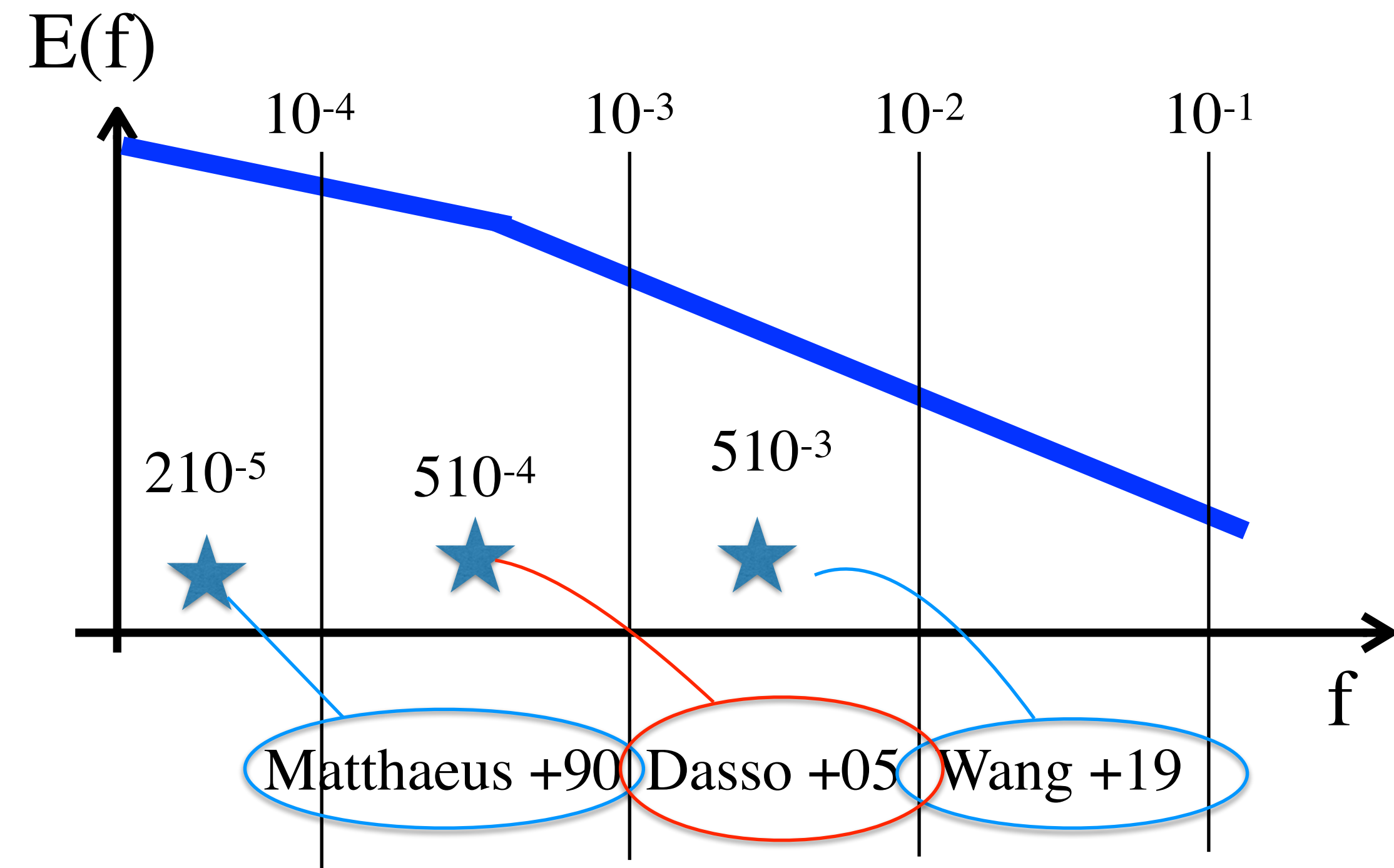
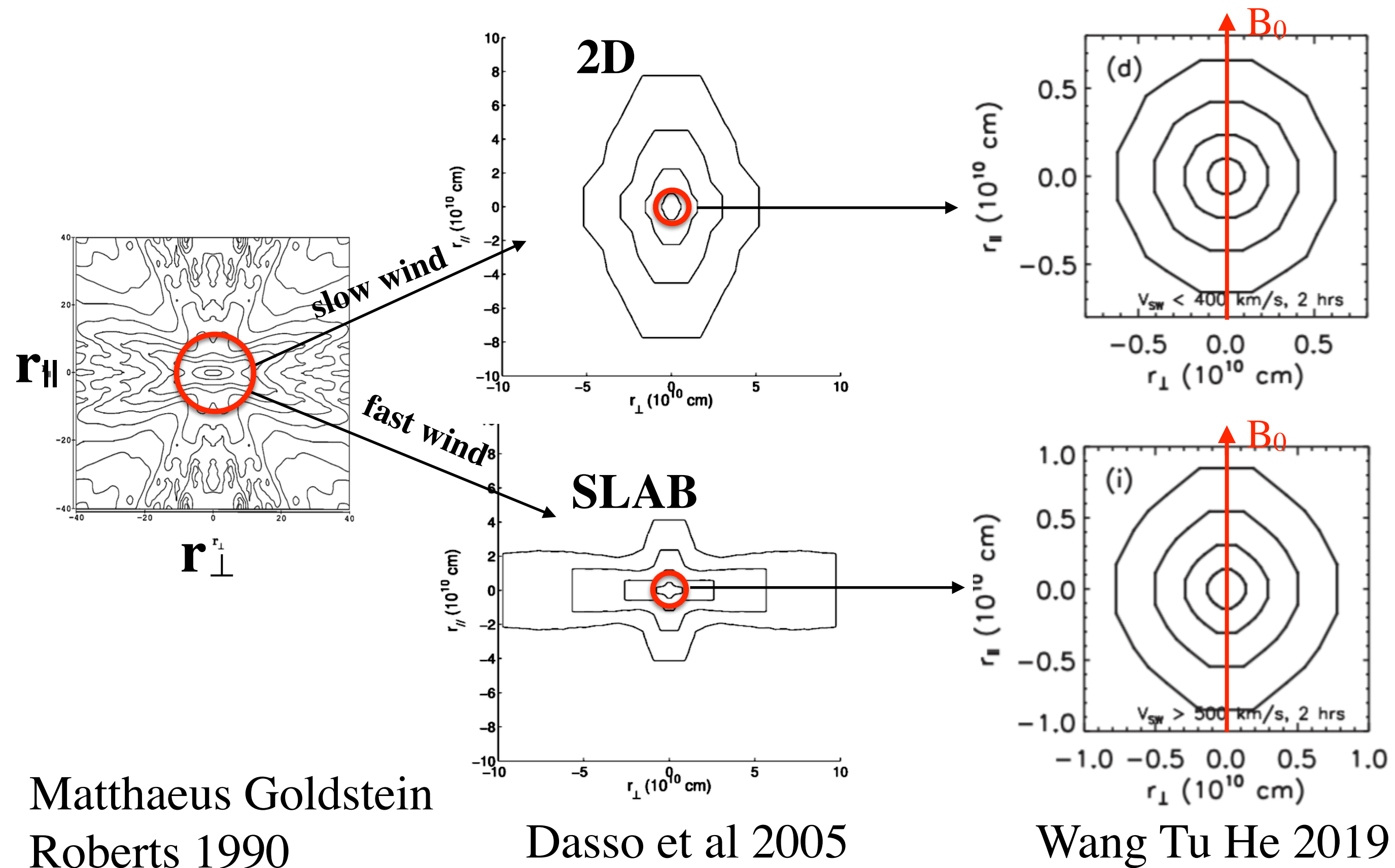
Energy in k_{\parallel}

Anisotropy at smaller scale



Recover of isotropy at small scales?

Observed Anisotropy in the SW



Measurement of Cascade

$$\nabla_{\ell} \cdot \mathbf{Y} = -4\epsilon$$

with isotropy

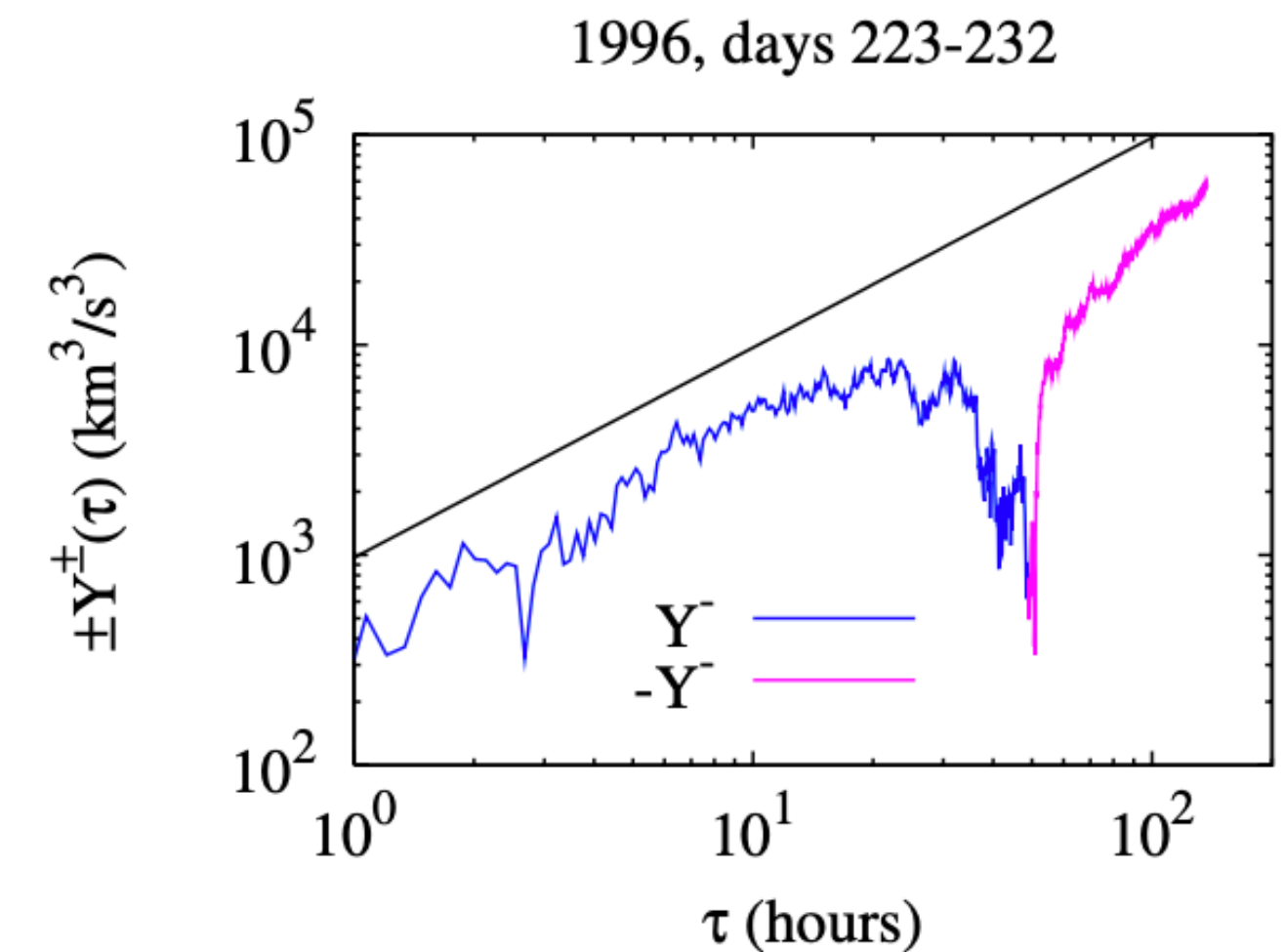
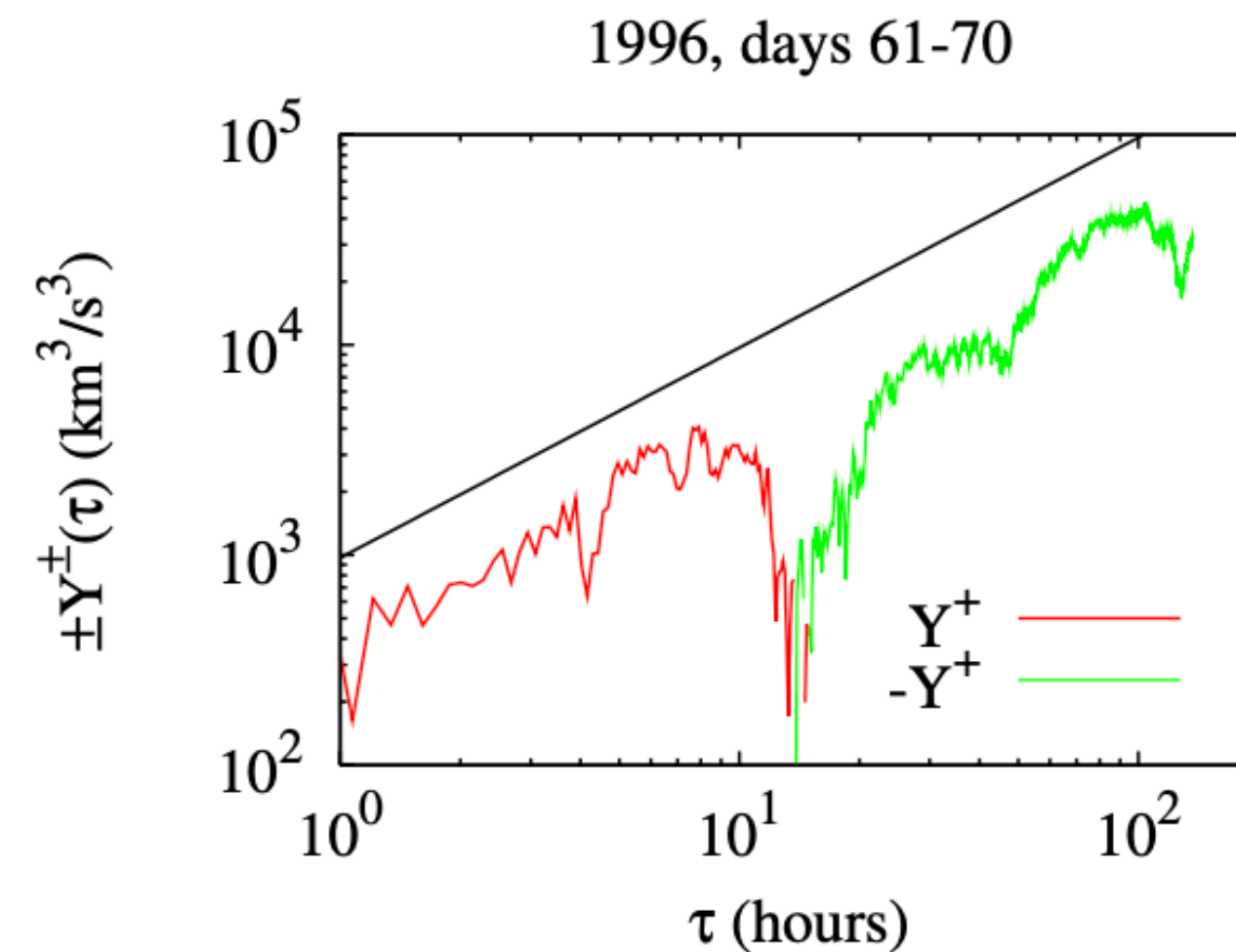
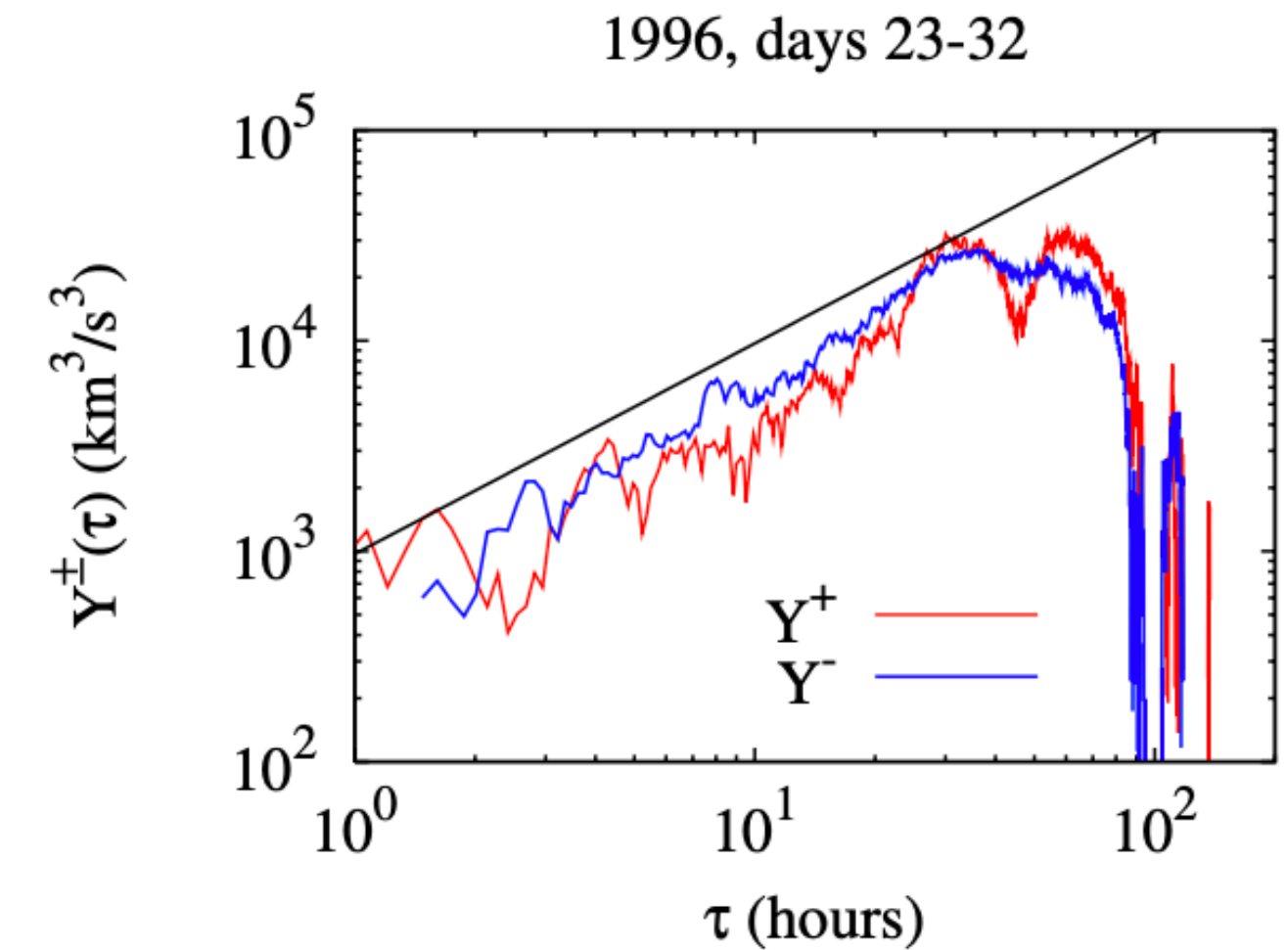
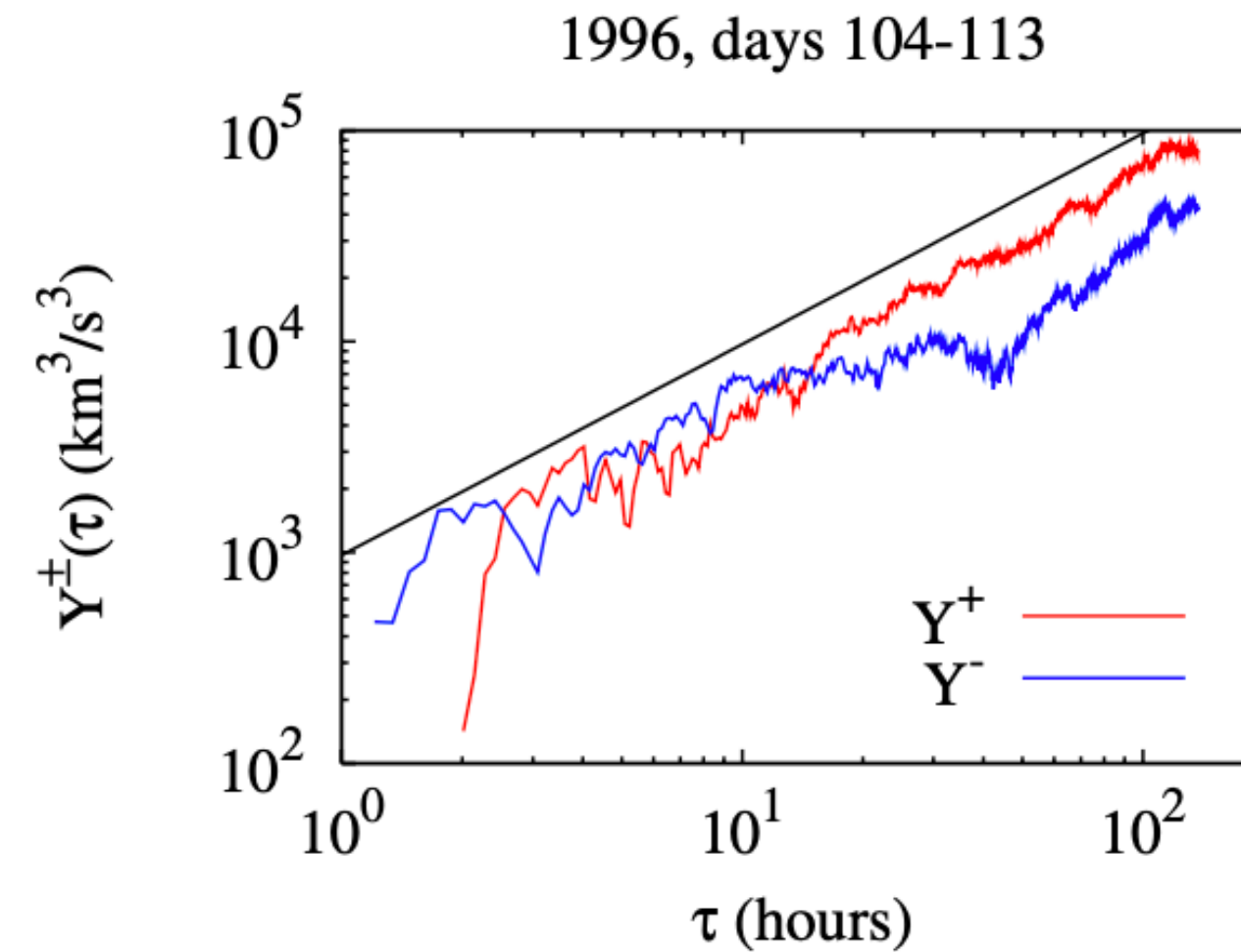
$$Y_{\ell} = -\frac{4}{3}\epsilon\ell$$

$$Y_{\ell} = \mathbf{Y} \cdot \hat{\ell}$$

Search for linear scaling of projected Yaglom vector

Change of sign

Linear scaling not found in all intervals



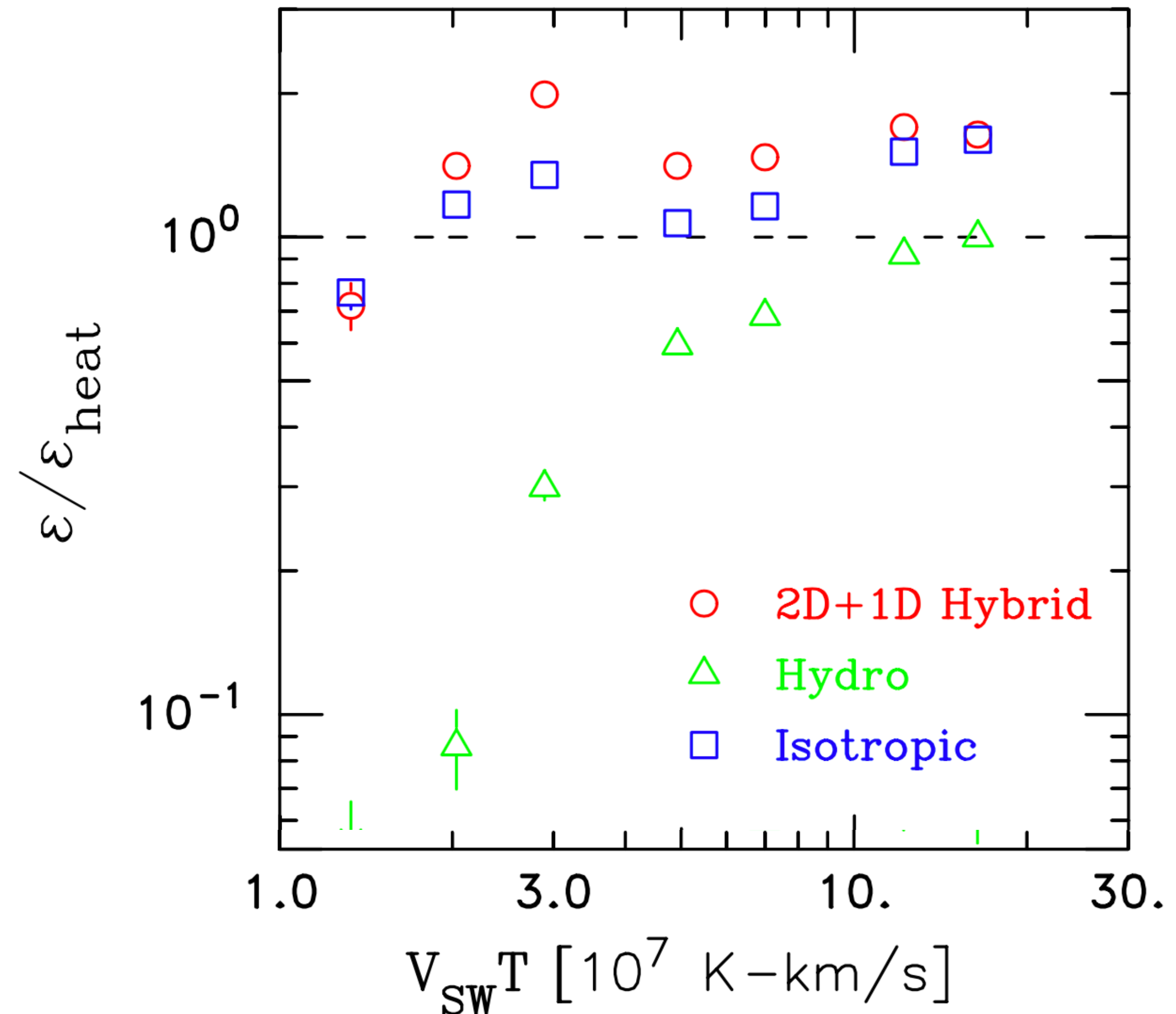
With Anisotropic Geometrical Model

$$\nabla_{\ell} \cdot \mathbf{Y} = -4\epsilon$$

$$Y_{\ell} = -\frac{4}{3}\epsilon\ell$$

with 2D + 1D
geometry

$$\epsilon^{2D} = -\frac{Y_{\ell_{\perp}}(\ell_{\perp})}{2\ell_{\perp}}, \quad \epsilon^{1D} = -\frac{Y_{\parallel}(\ell_{\parallel})}{4\ell_{\parallel}}.$$

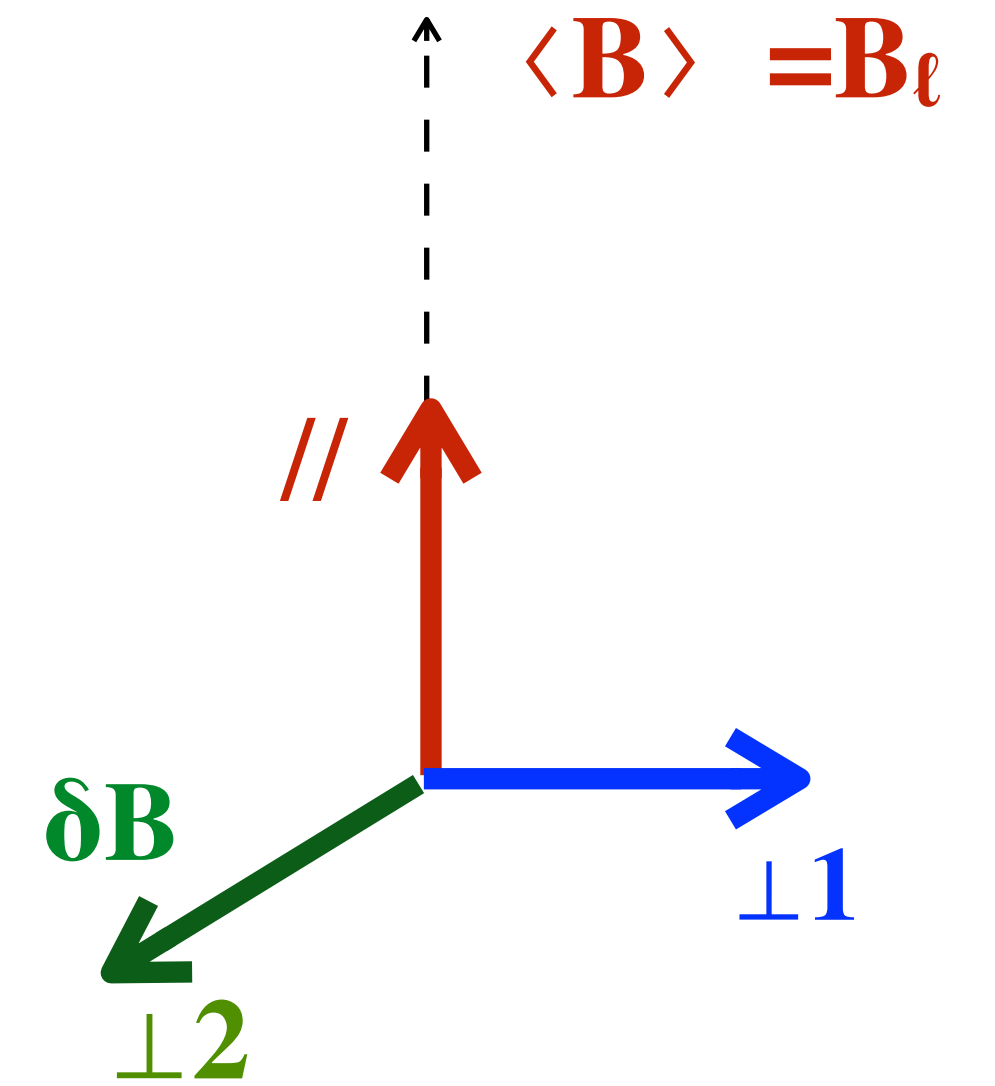
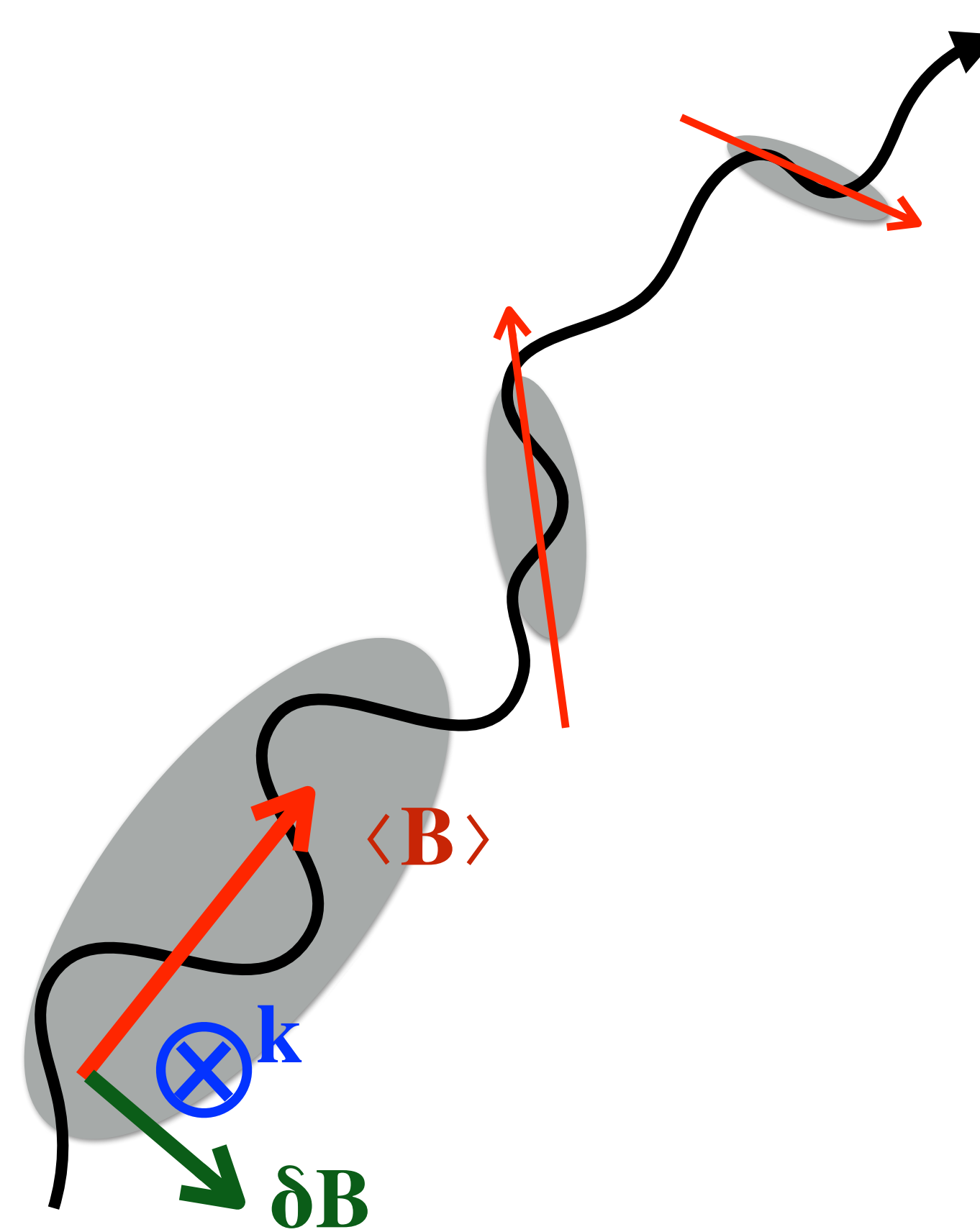
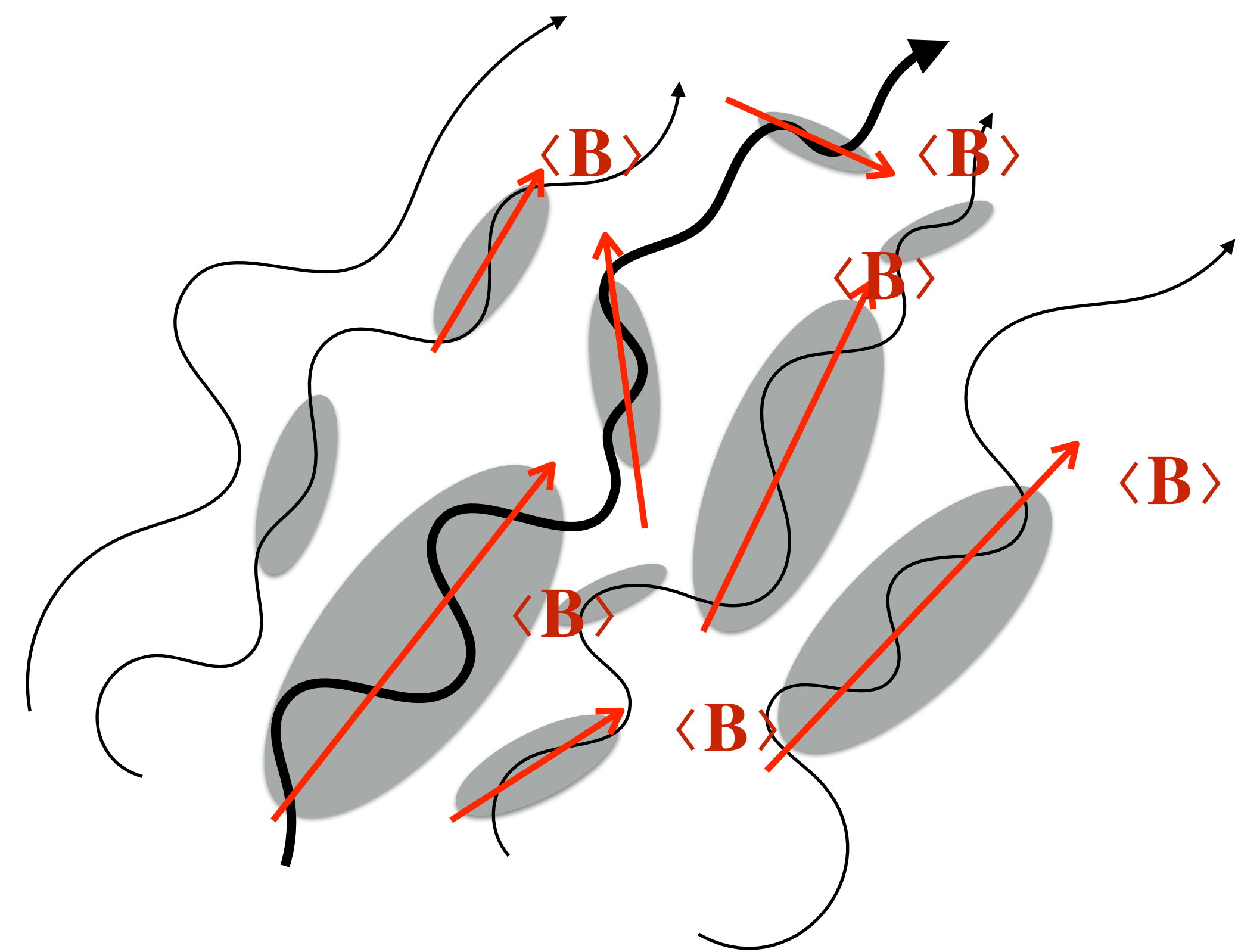


Cascade is larger in slow/cold solar wind turbulence

Scale Dependent Anisotropy

Compute anisotropy with respect to a scale dependent definition of mean magnetic field

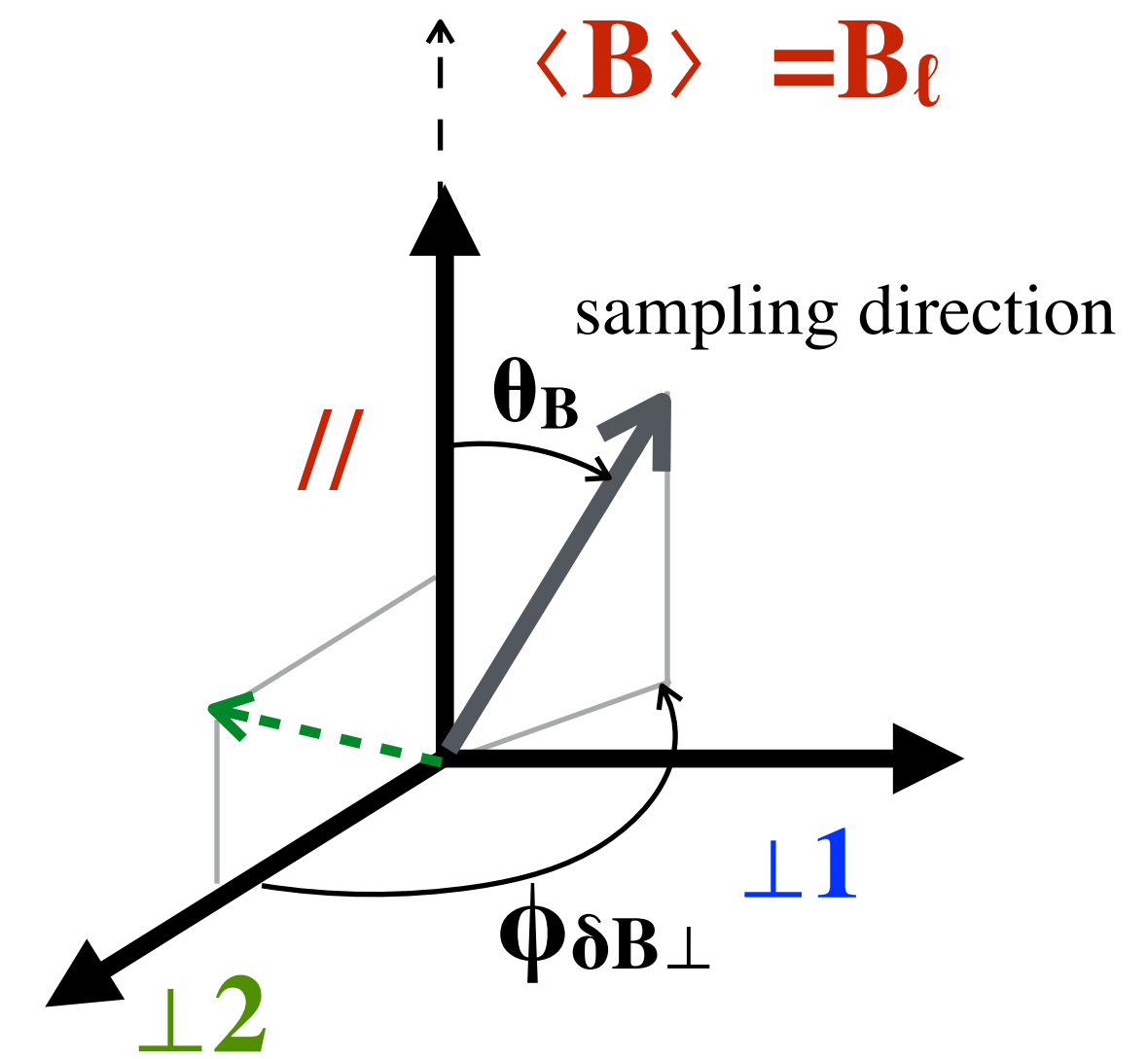
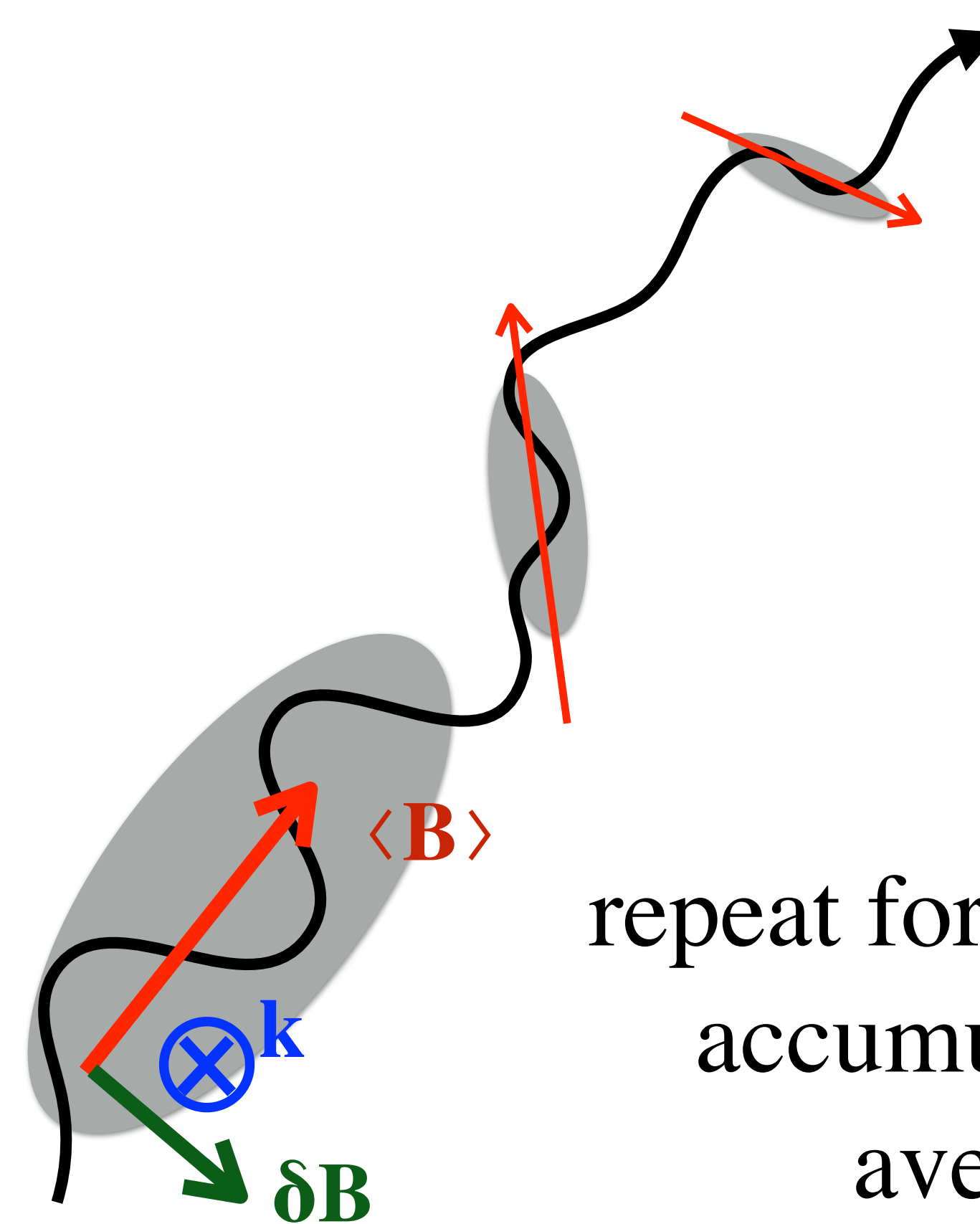
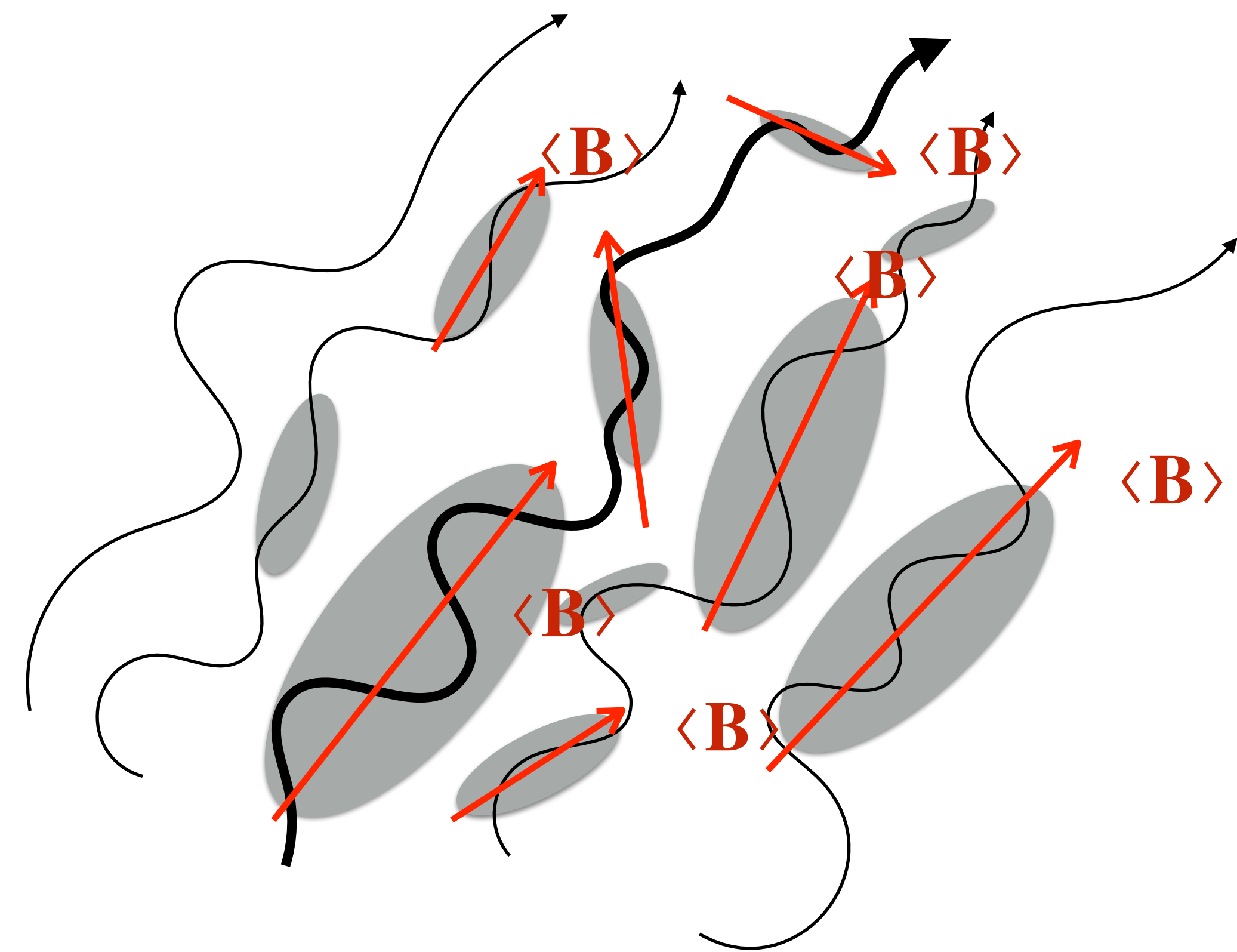
Not a Π -order quantity!



Scale Dependent Anisotropy

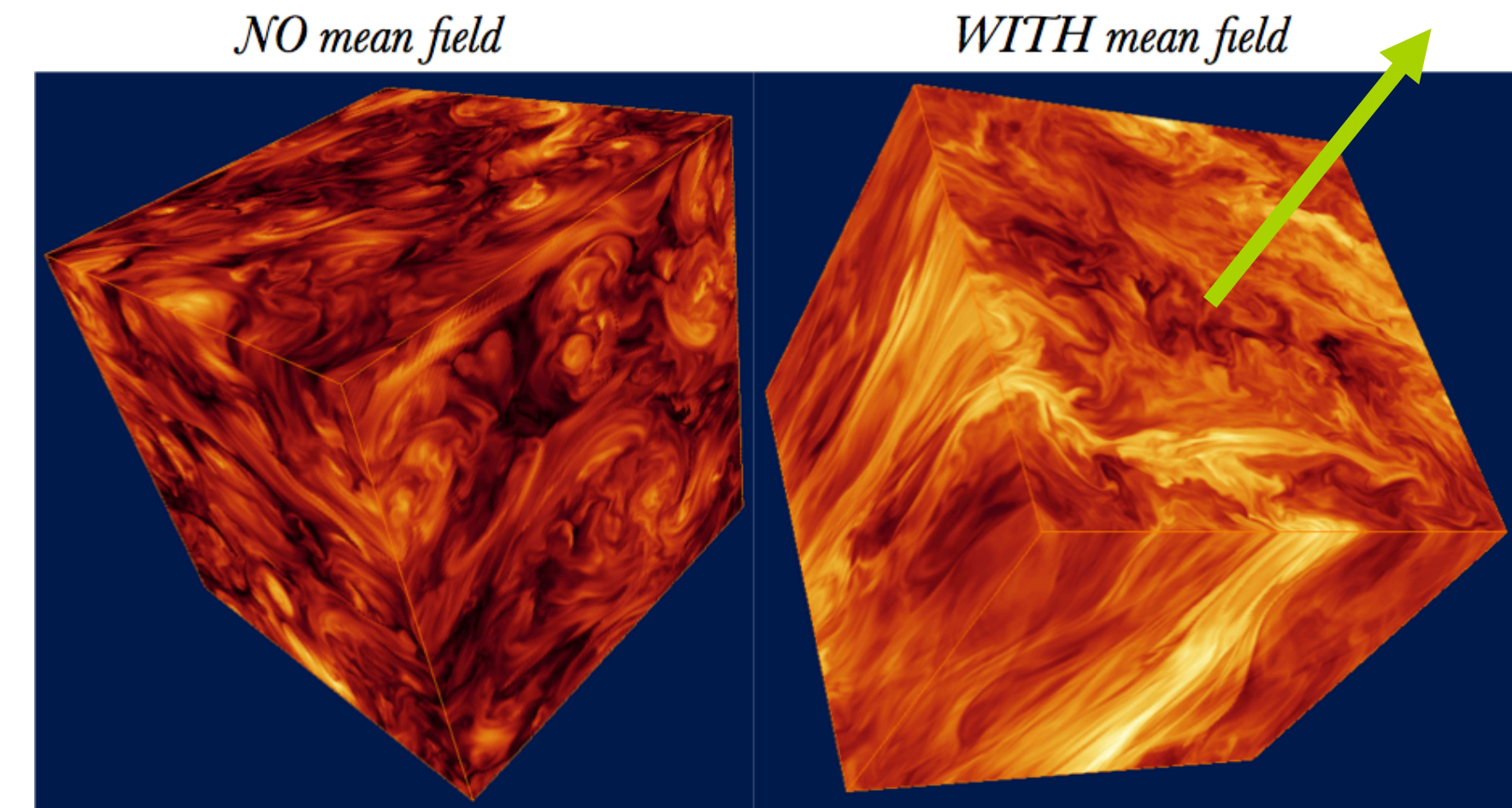
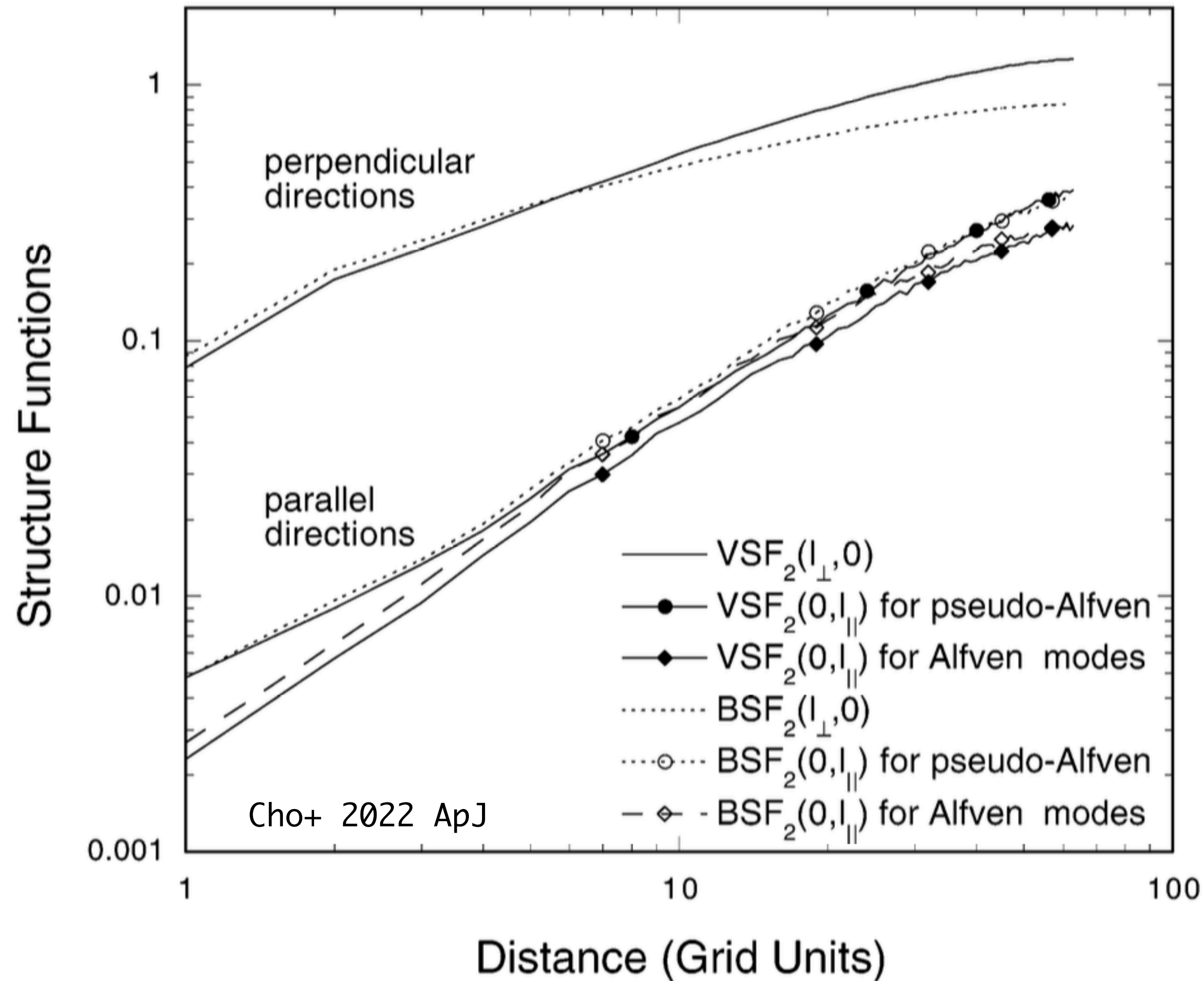
Compute anisotropy with respect to a scale dependent definition of mean magnetic field

Not a Π -order quantity!

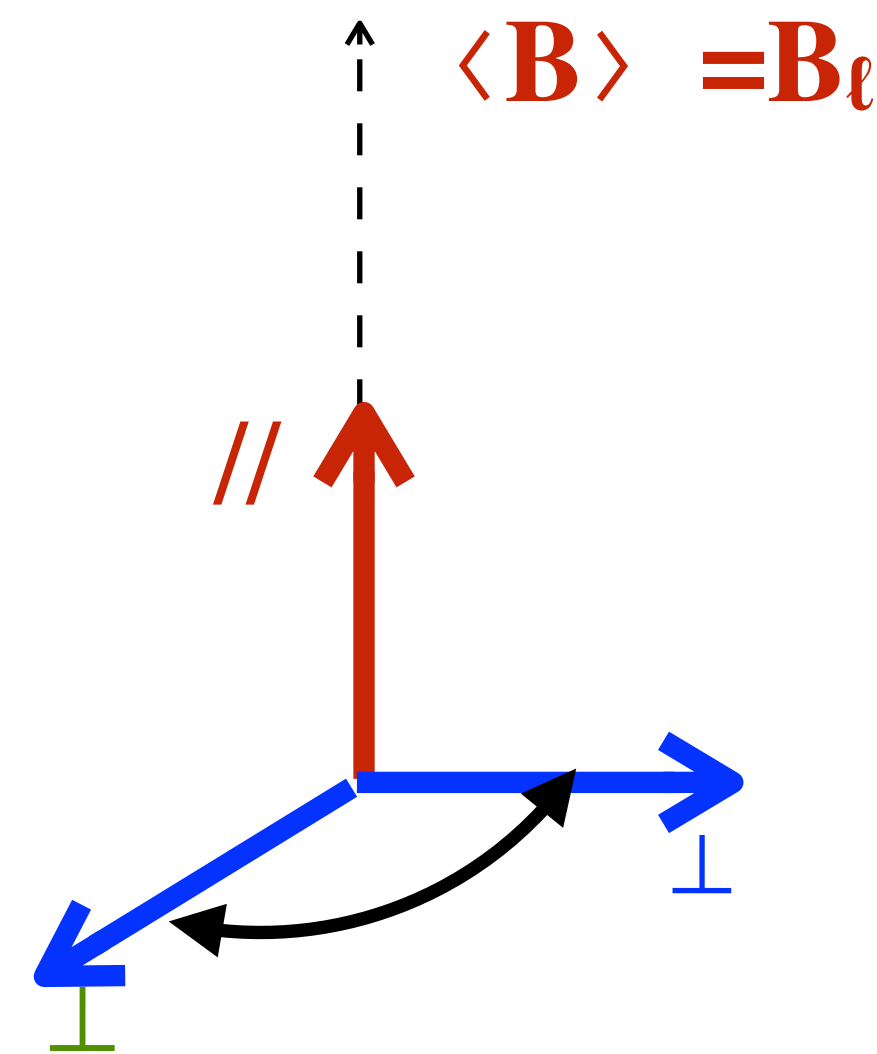


repeat for each point in space,
accumulate statistics and
average to obtain
 $SF(\ell_{\parallel}, \ell_{\perp 1}, \ell_{\perp 2})$

Scale Dependent Anisotropy



2D



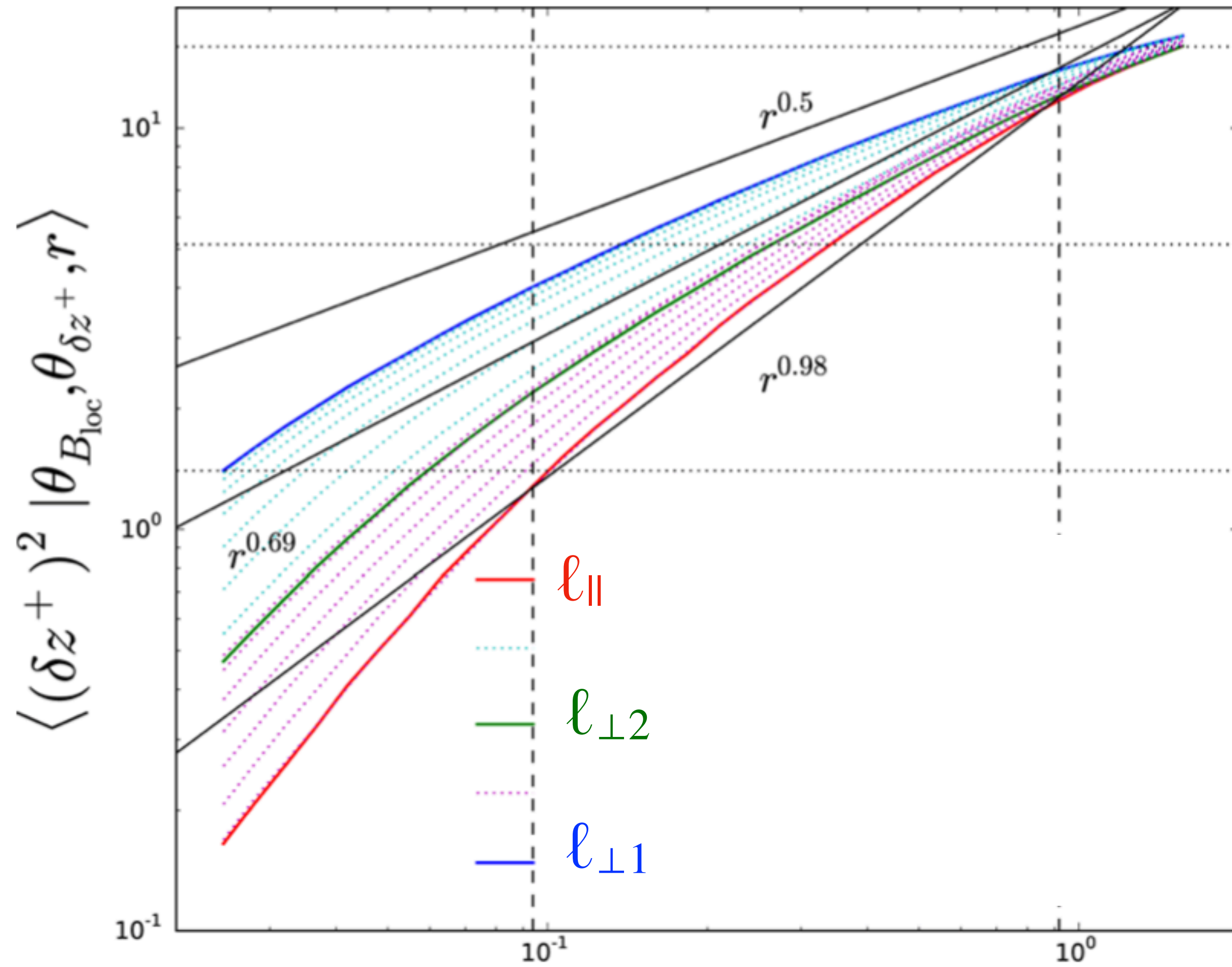
also w/o mean field
you get an
anisotropic scaling

Origin:
critical balance

Goldreich & Shridar 1995 ApJ

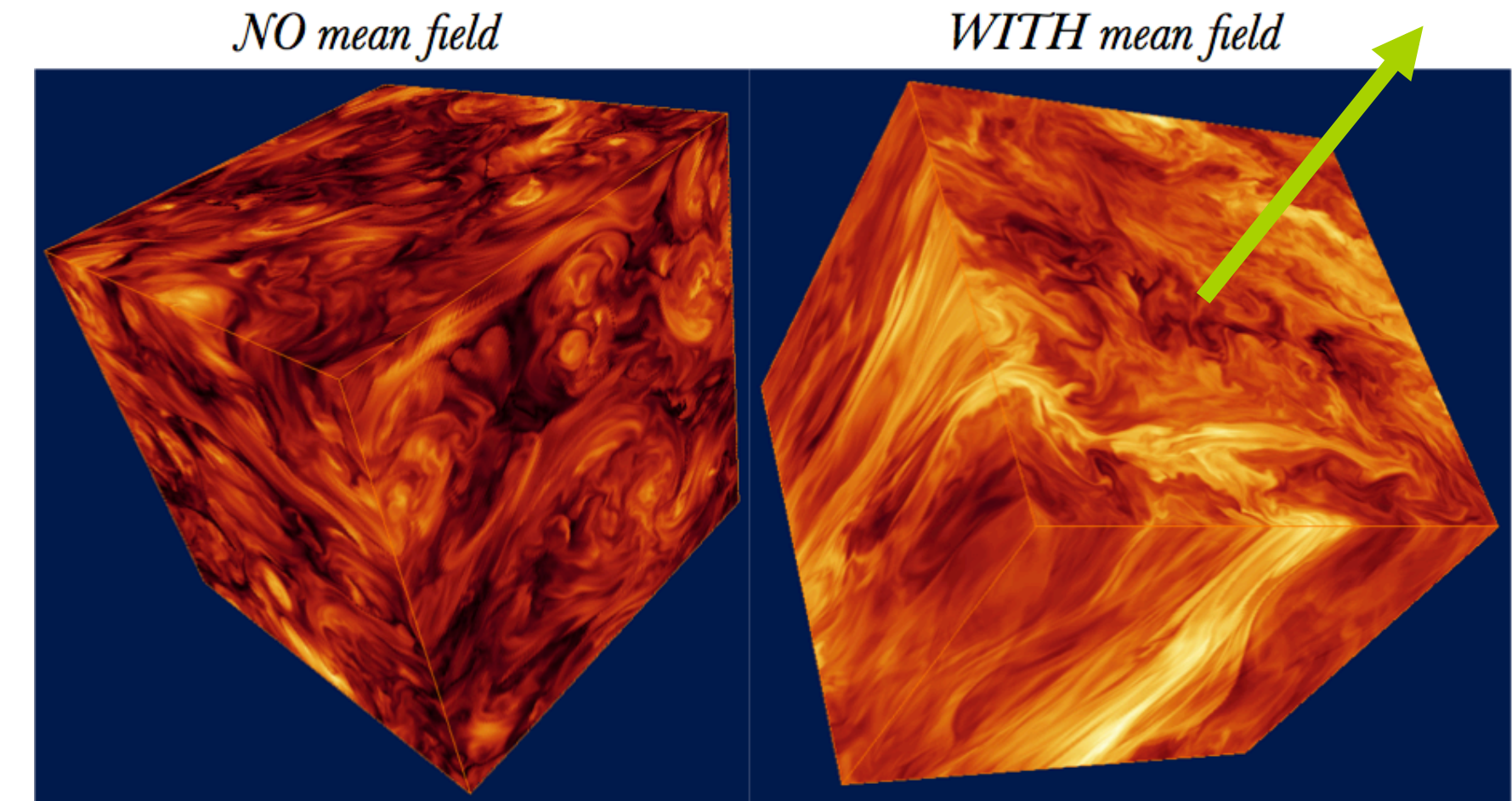
Scale Dependent Anisotropy

Scaling of $SF(\ell_{\parallel}, \ell_{\perp 1}, \ell_{\perp 2})$

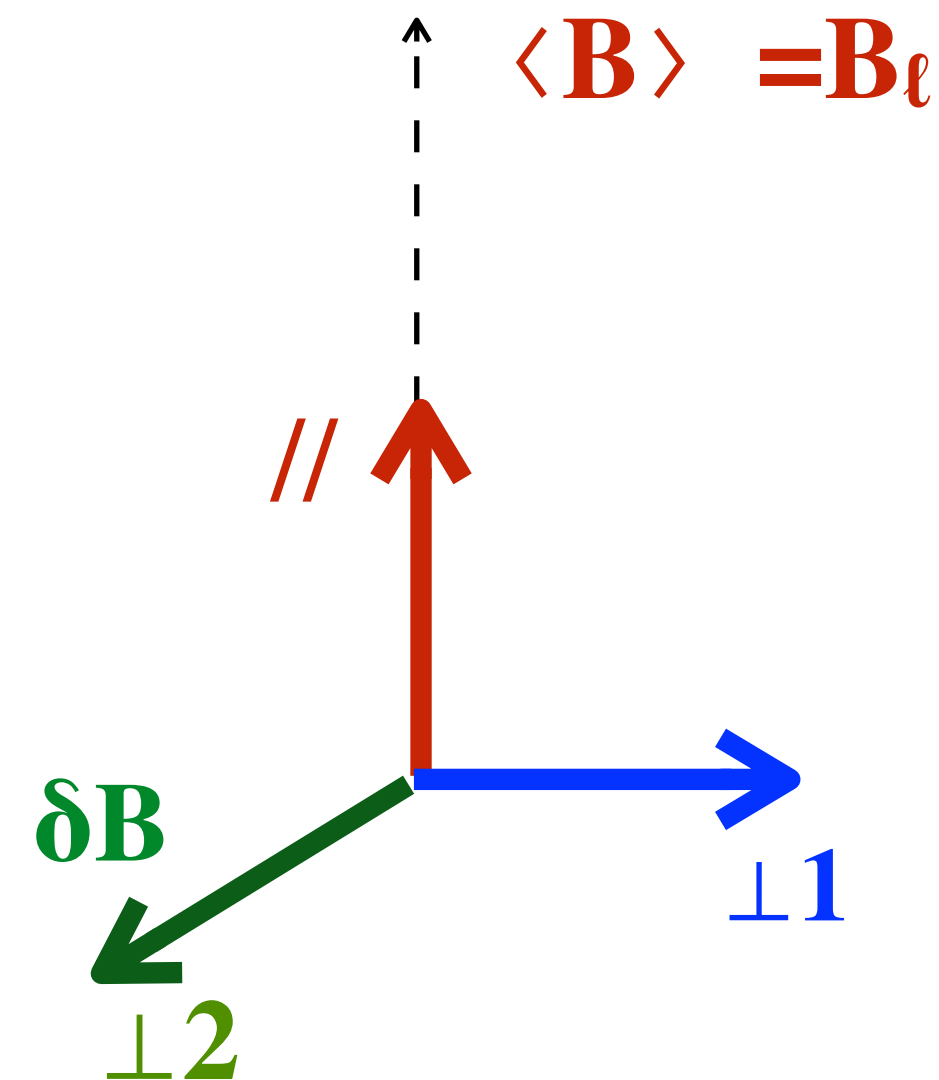


Mallet et al 2016

r



3D

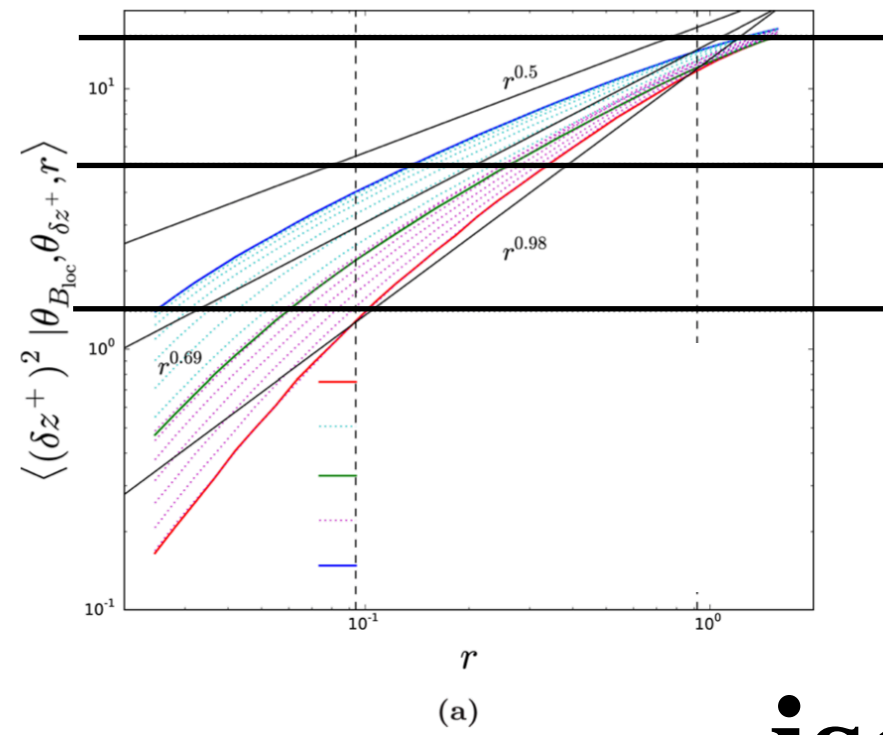


also w/o mean field
you get an
anisotropic scaling

Origin:
critical balance +
alignment

Boldyrev 2005 ApJ

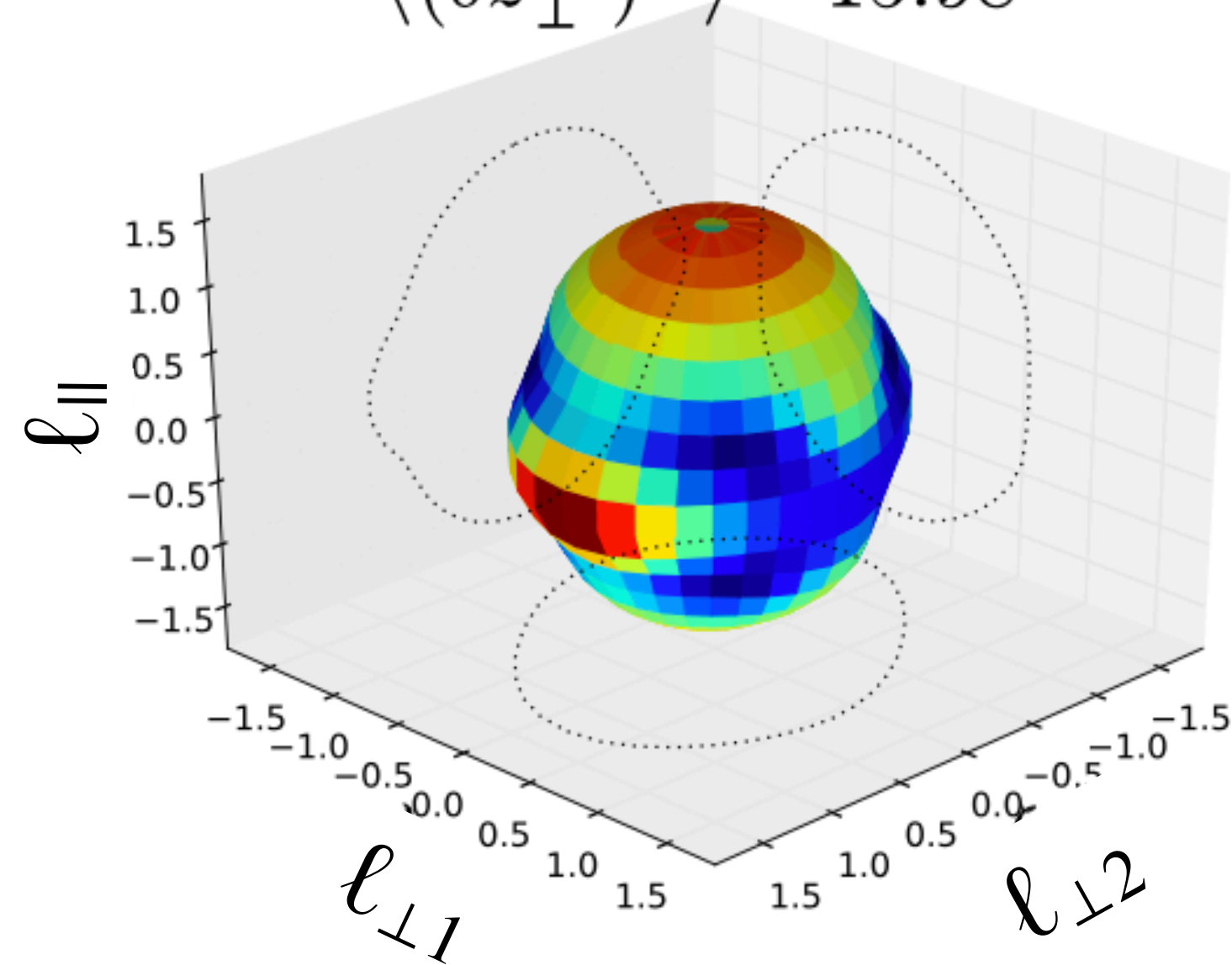
Scale Dependent Anisotropy



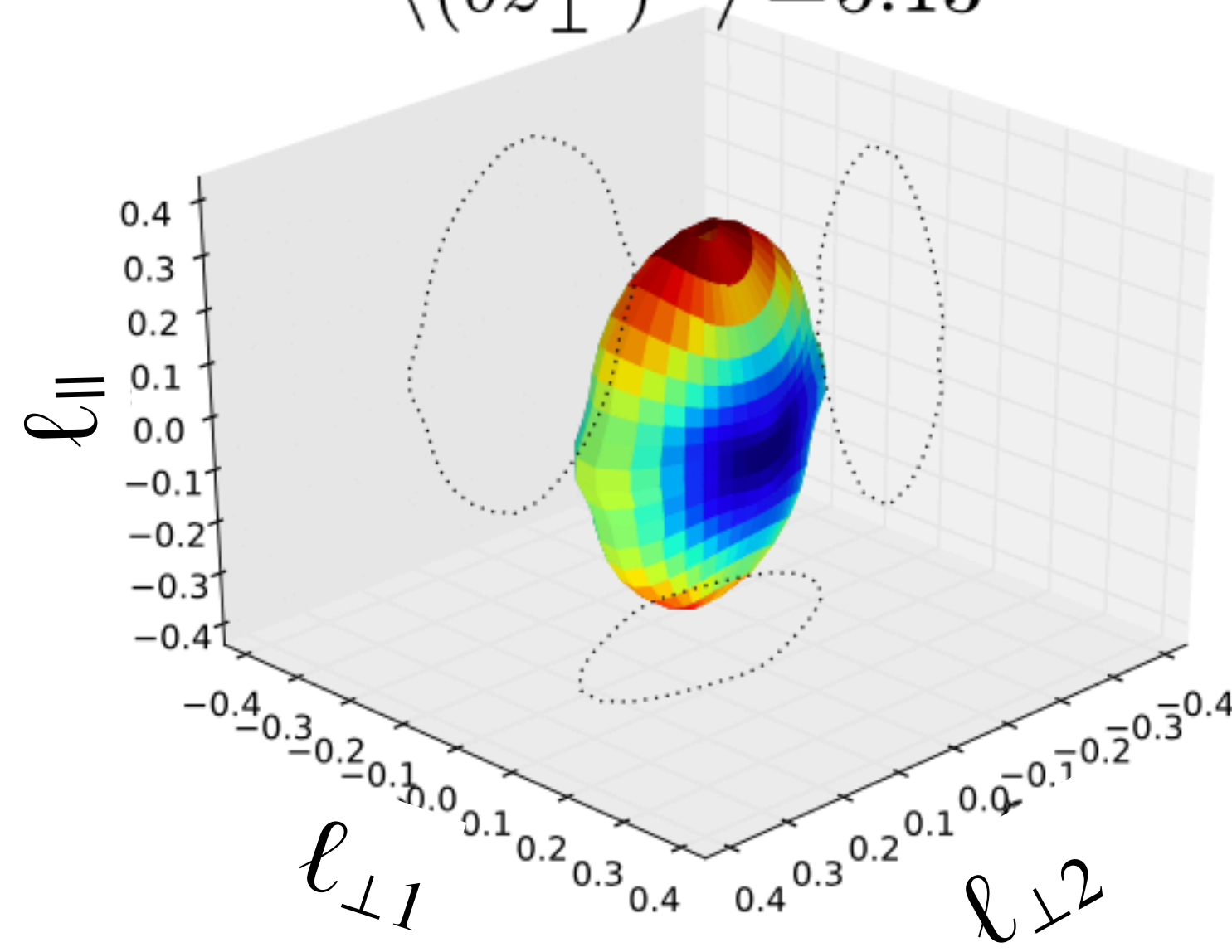
Statistical representation of turbulent eddies at different scales

isotropic

$$\langle (\delta z_{\perp}^+)^2 \rangle = 15.98$$

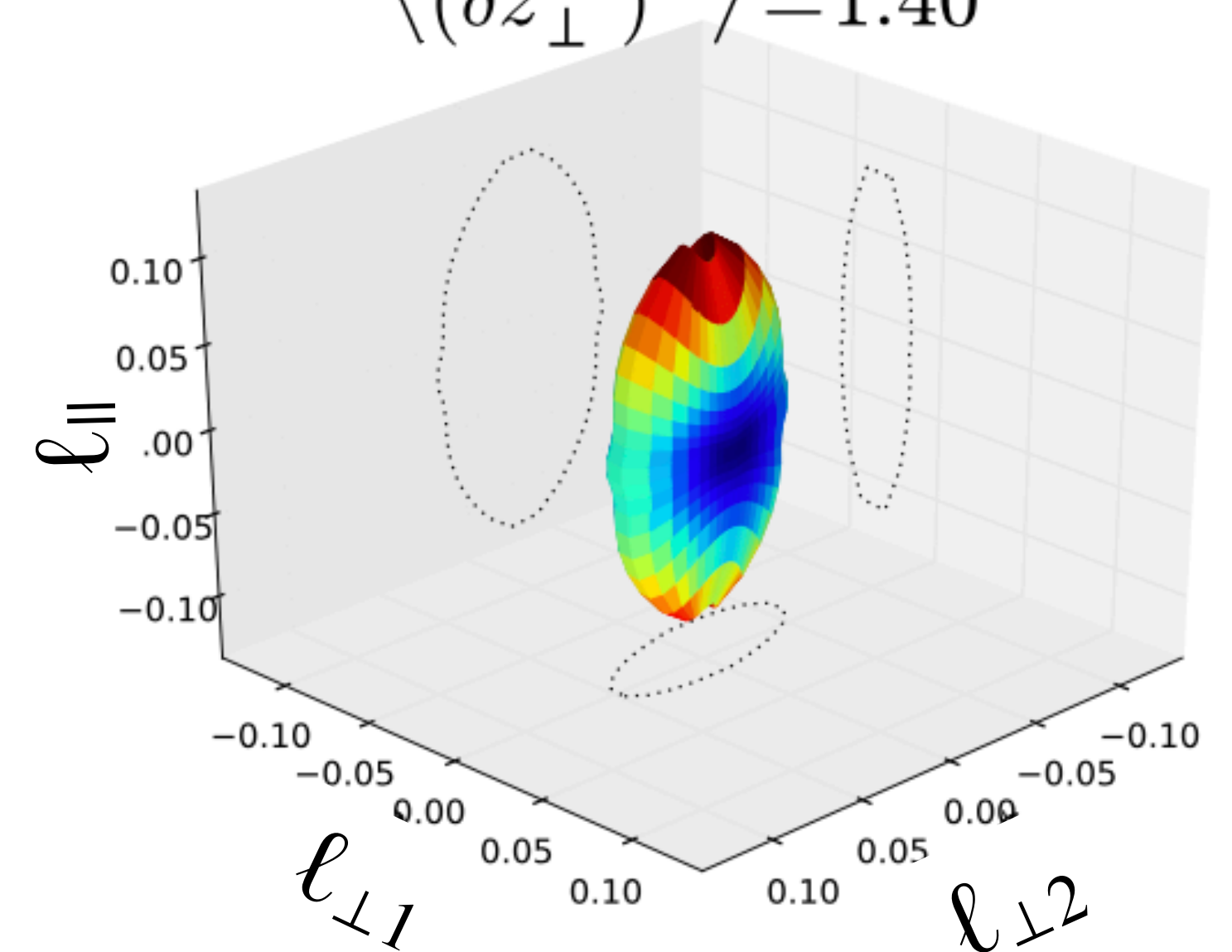


$$\langle (\delta z_{\perp}^+)^2 \rangle = 5.13$$



sheet-like
(not tube-like)

$$\langle (\delta z_{\perp}^+)^2 \rangle = 1.40$$

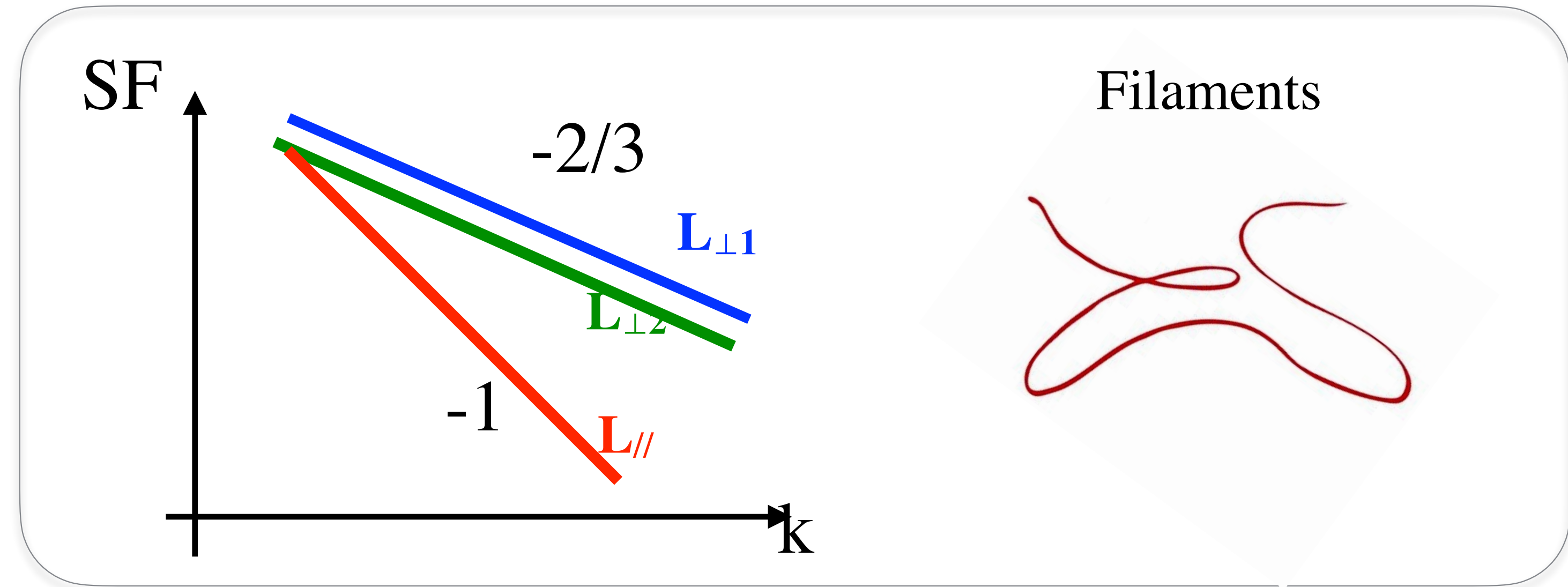


smaller energies (smaller scales)



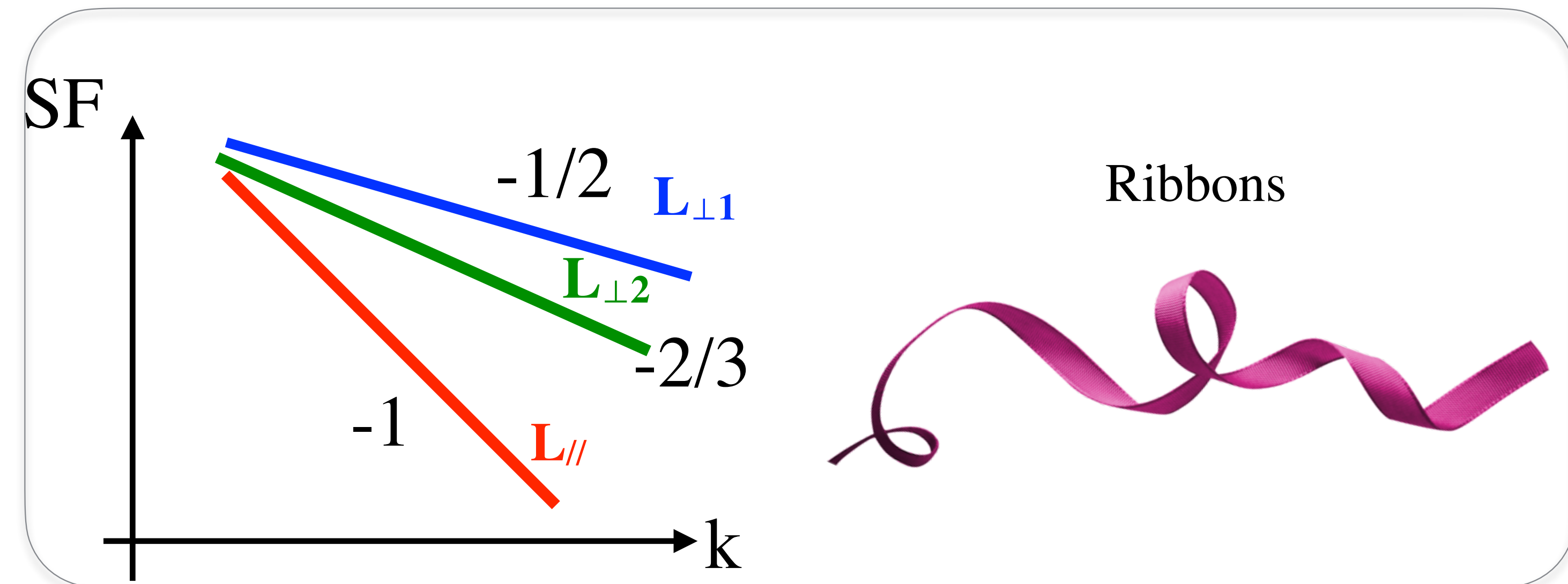
Different small-scale structures

Critical Balance
(Goldreich-Shirdar 1995)



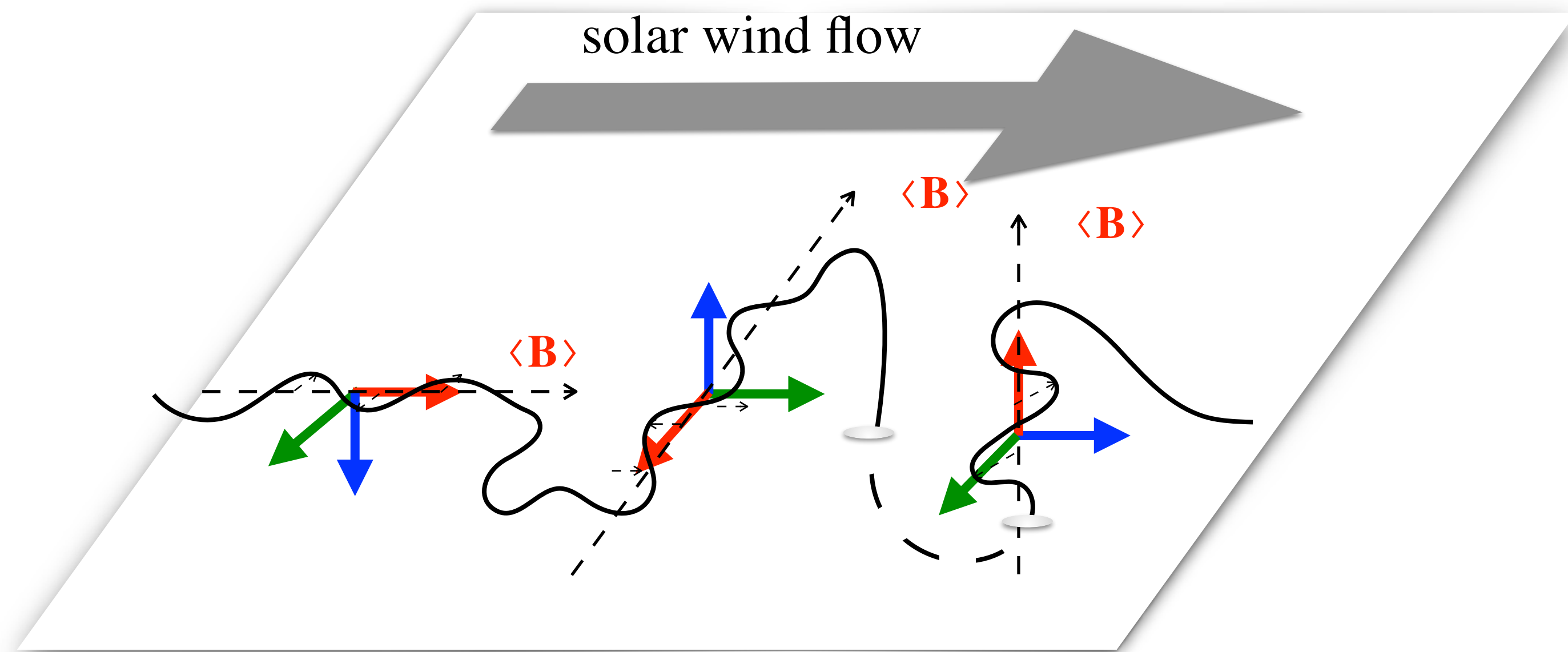
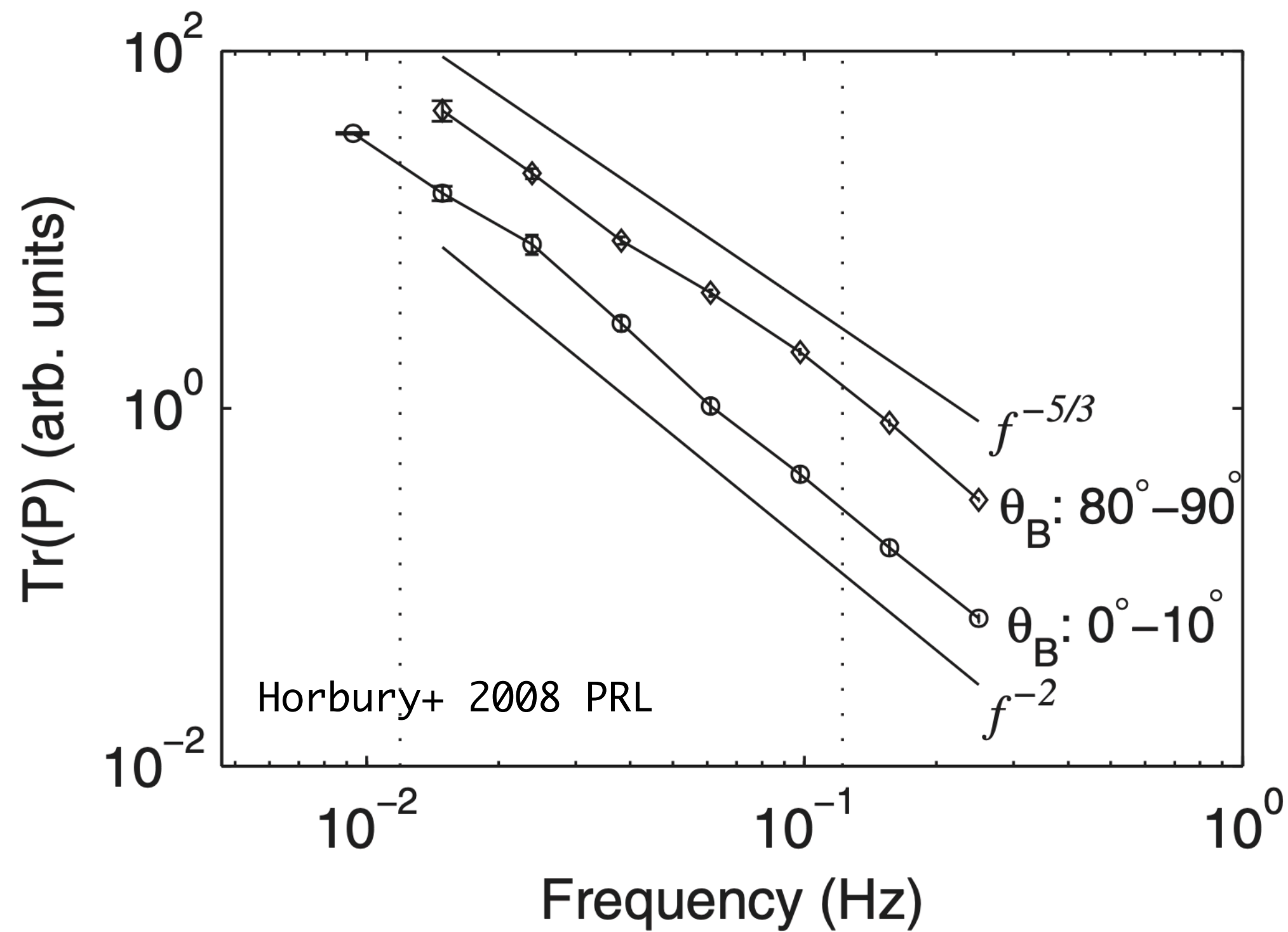
or

Critical Balance
& δV - δB alignment
(Boldyrev 2005, 2006)



Measurements in the SW

wandering of field lines allows measuring all directions



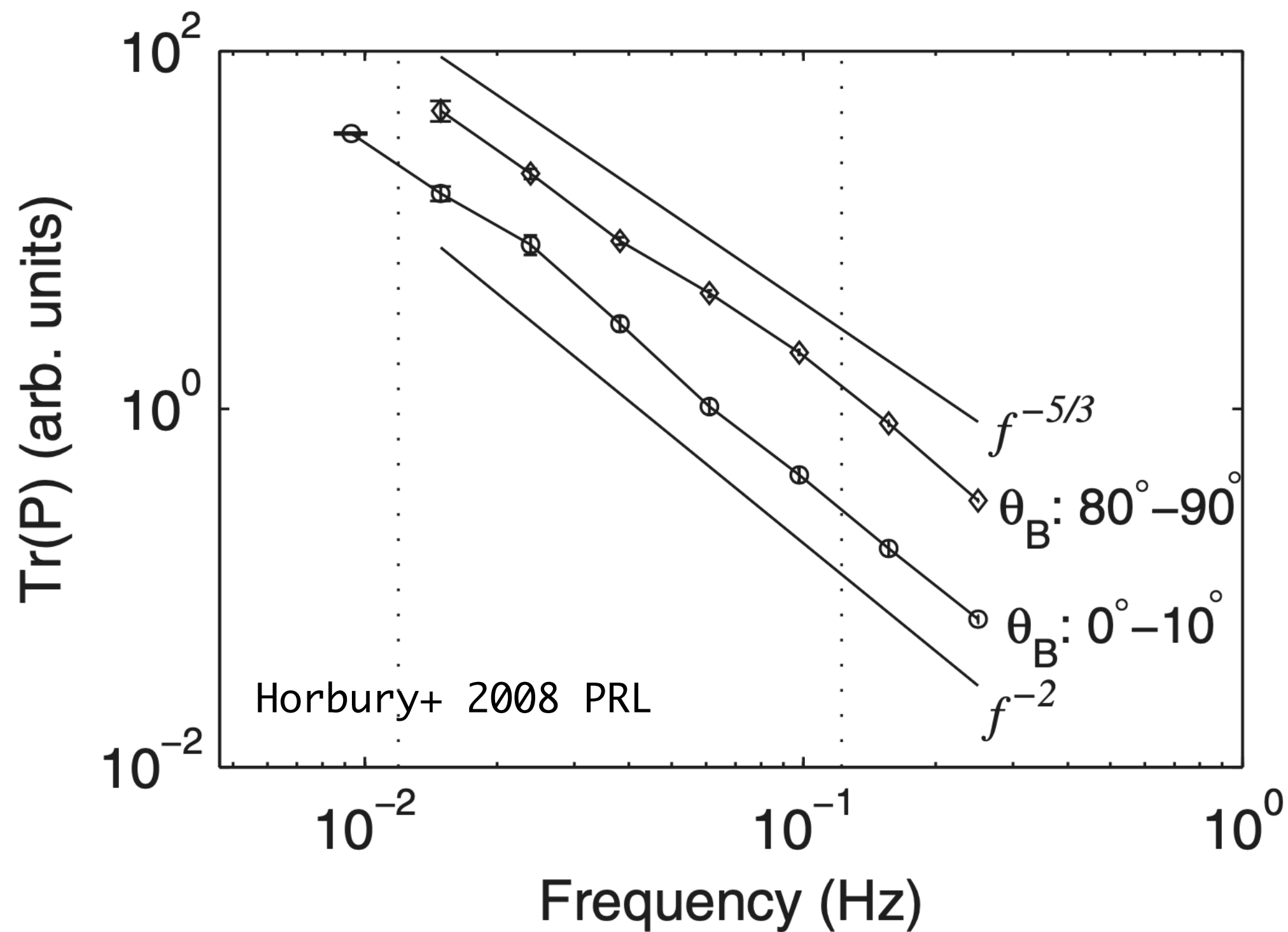
$\text{SF}(\ell_{\parallel})$

$\text{SF}(\ell_{\perp 2})$

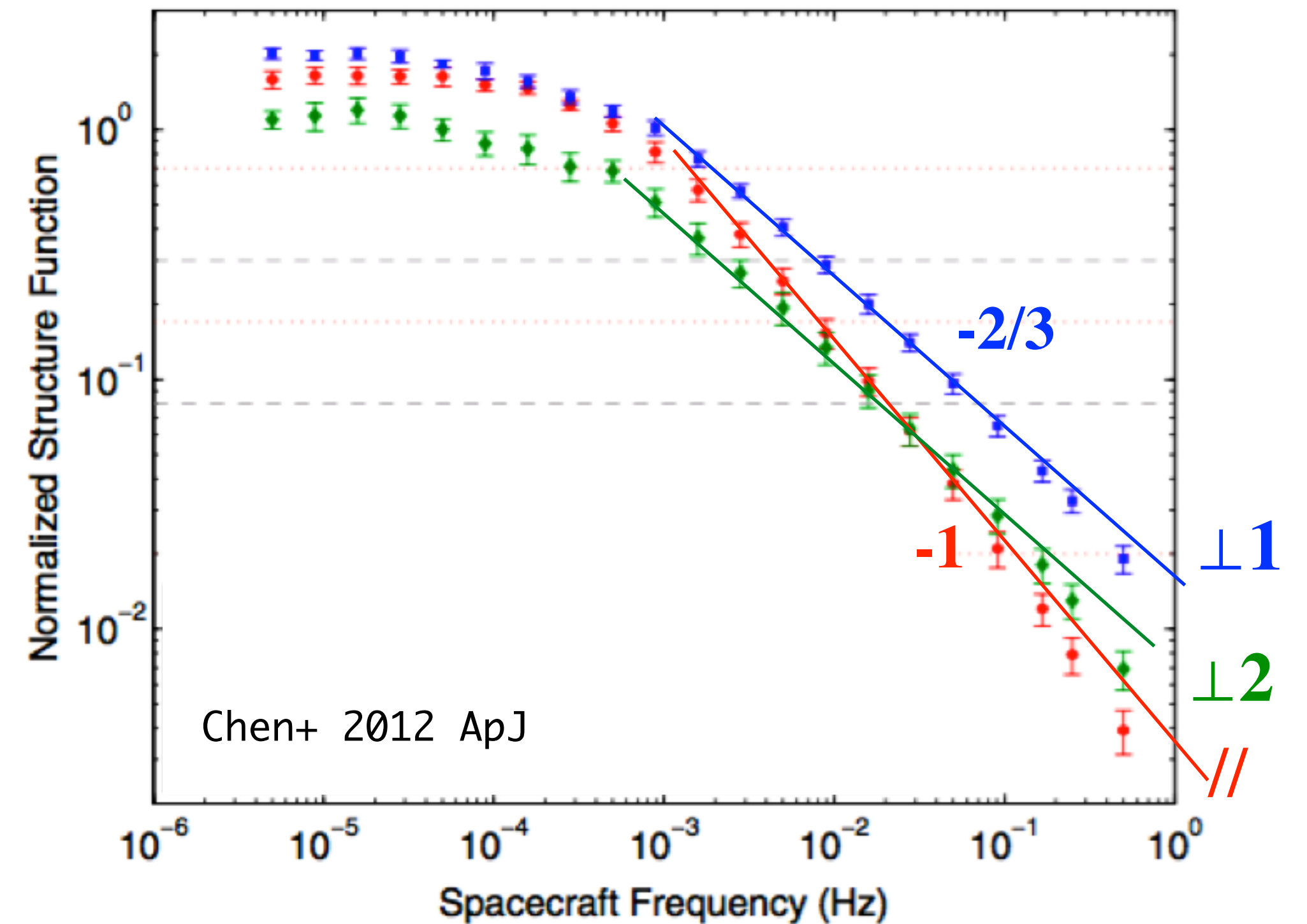
$\text{SF}(\ell_{\perp 1})$

Different Parallel and perpendicular scaling
(with power-law indexes matching theory)

Measurements in the SW



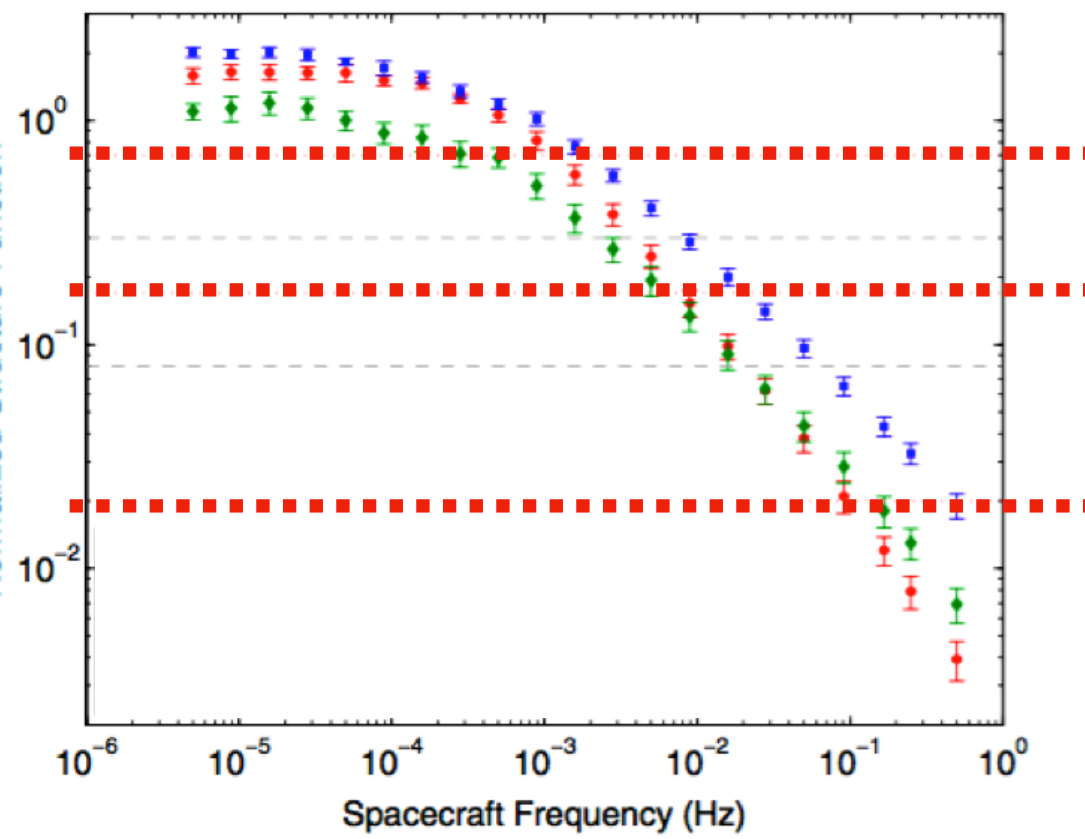
Different Parallel and perpendicular scaling
(with power-law indexes matching theory)



The two perpendicular direction have the same scaling

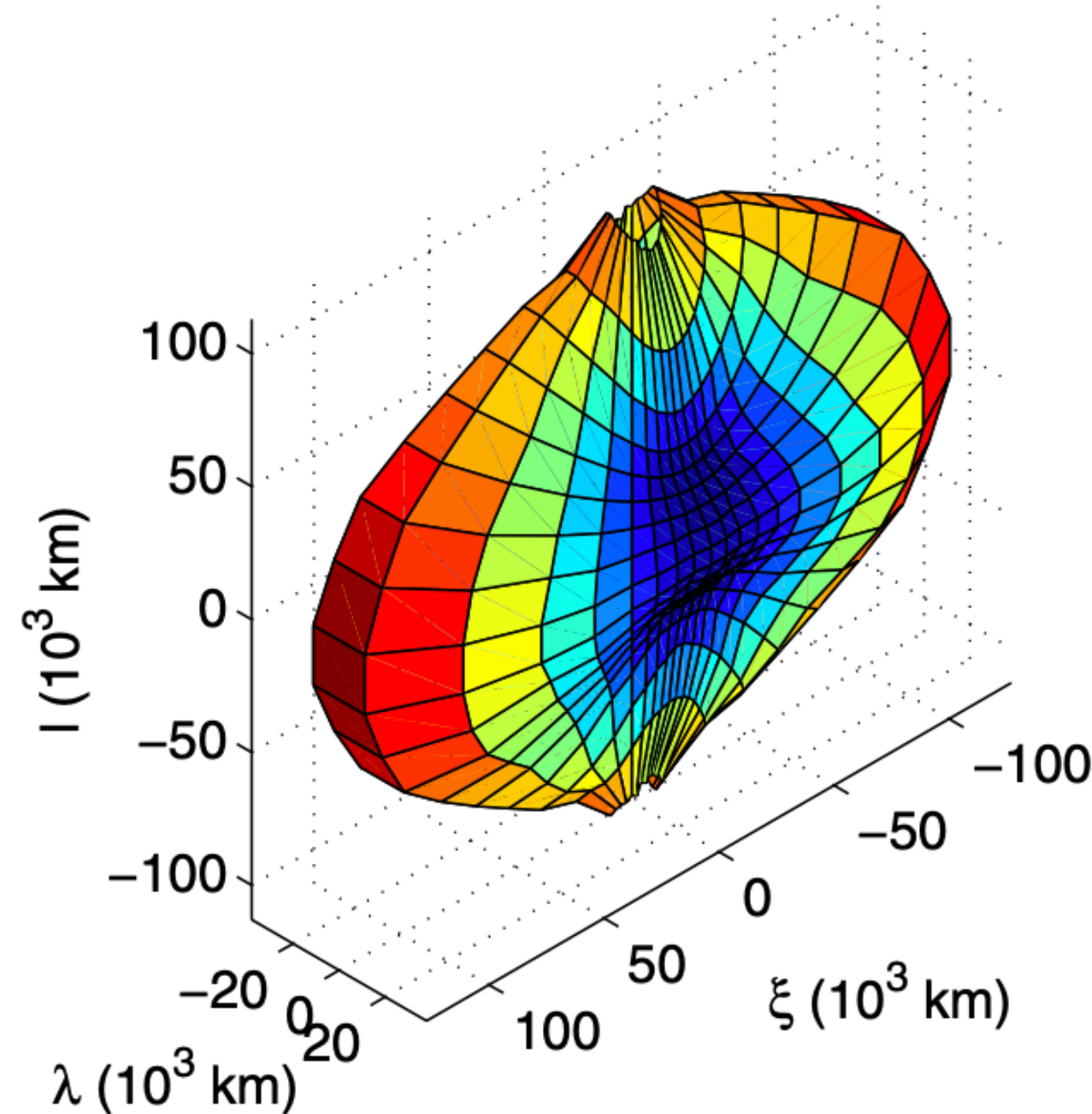
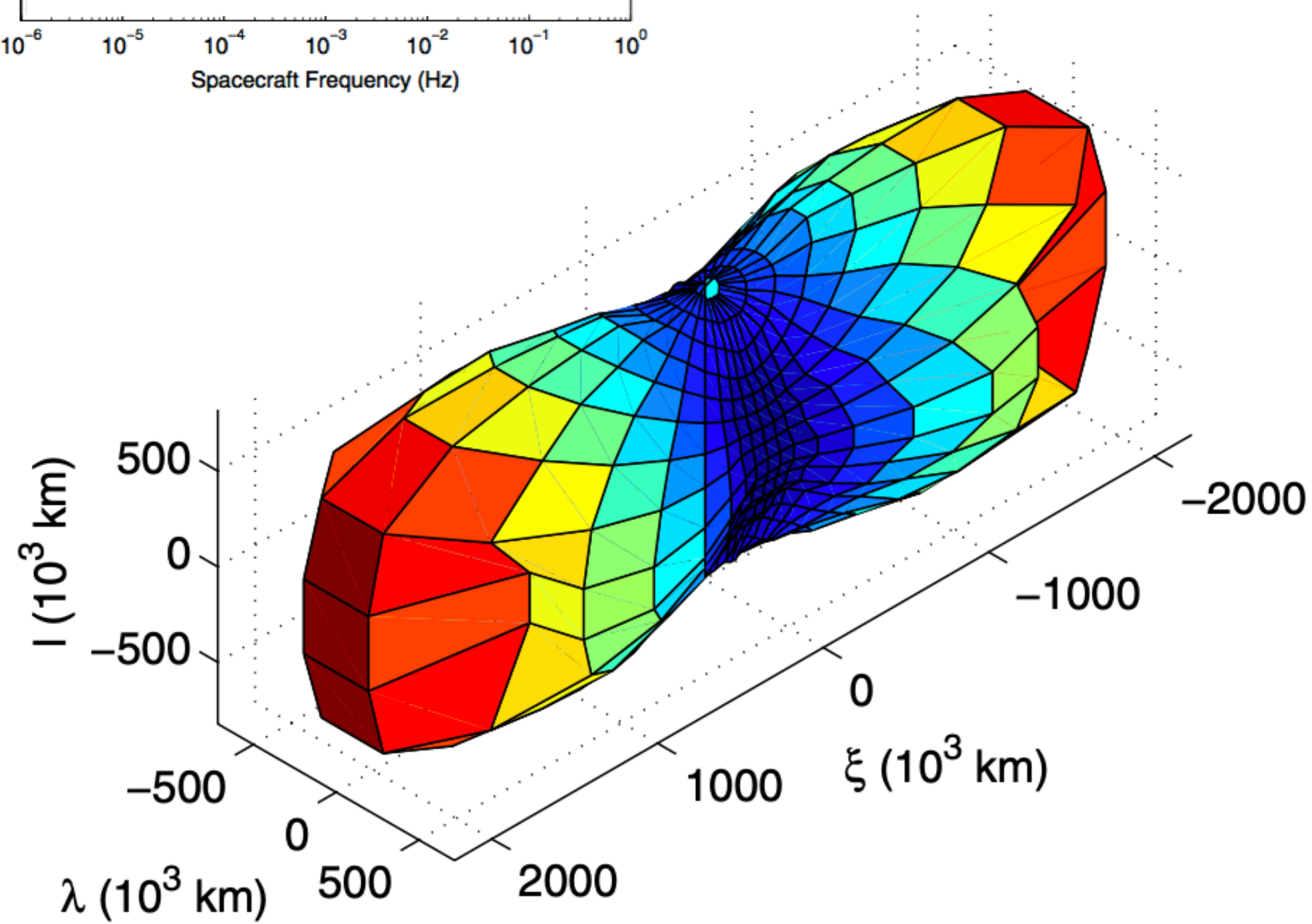
tube -like structures at small scales

Measurements in the SW



Anisotropic

Two perp direction conserve aspect ratio



smaller energies (smaller scales)

